## Circles

## VS

## Hyperbolas

Geometries
Trigonometries
Calcului

## Define Euclidian angles



$$
\theta=s
$$

The angle is equal to the arc length

## Euclidian (Circle)

Geometry and Trigonometry

Distance (Pythagoras)

$$
d^{2}=x^{2}+y^{2}
$$

Fundamental Identity

$$
\cos ^{2}(\theta)+\sin ^{2}(\theta)=1
$$

Adding Angles (add arcs)

$$
\begin{aligned}
& \theta_{1}=\mathbf{s}_{1} \\
& \theta_{\mathbf{2}}=\mathbf{s}_{2} \\
& \theta_{\mathbf{3}}=\theta_{1}+\theta_{2}=\mathbf{s}_{1}+\mathbf{s}_{\mathbf{2}}=\mathbf{s}_{3}
\end{aligned}
$$



## Define Hyperbolic angles



The angle is equal to the arc length


## Hyperbolic

## Geometry and Trigonometry

Distance (Pythagoras)

$$
d^{2}=t^{2}-x^{2}
$$

Fundamental Identity

$$
\cosh ^{2}(\theta)-\sinh ^{2}(\theta)=1
$$

Adding Angles (add arcs)

$$
\theta_{1}=\mathbf{s} 1
$$

$$
\theta_{\mathbf{2}}=\mathbf{s} \mathbf{2}
$$

$$
\theta_{1}+\theta_{\mathbf{2}}=\mathbf{s}_{1}+\mathbf{s}_{2}
$$

```
Trig Identities
\(\sinh (\theta+\phi)=\) \(\sinh (\theta) \cosh (\phi)+\cosh (\theta) \sinh (\phi)\)
```


## Add Hyperbolic angles



The angle is equal to the arc length

## Euclidian Triangle Trig



## Hyperbolic Triangle Trig



Euclidian Triangles
$\cos ^{2}(\theta)=\left[1+\tan ^{2}(\theta)\right]^{-1}$
$3^{2}+4^{2}=5^{2}$


Hyperbolic Triangles
$\cosh ^{2}(\theta)=\left[1-\tanh ^{2}(\theta)\right]^{-1}$
$5^{2}-3^{2}=4^{2}$



Figure 5.3: Some hyperbolic right triangles.


Figure 5.4: More hyperbolic right triangles. The right angle is on the left!

Even More Right Triangles





Trigonometric functions on $R$ (cos: purple; sin: red; tan: blue)


Hyperbolic functions on $r$ (cosh: purple; sinh: red; tanh: blue)




## Warning!

The following article is from The Great Soviet Encyclopedia (1979).

It might be outdated or ideologically biased.



Figure 2. Graphs of trigonometric functions: (1) sine, (2) cosine, (3) tangent, (4) cotangent, (5) secant, (6) cosecant

# http://en.wikipedia.org/wiki/File:HyperbolicAnimation.gif 

## Euclidian Calculus

$$
\begin{aligned}
& \text { d } \cos (\theta) / d \theta=-\sin (\theta) \\
& \text { d } \sin (\theta) / d \theta=\cos (\theta) \\
& \text { d } \tan (\theta) / d \theta=\sec ^{2}(\theta)
\end{aligned}
$$

## Hyperbolic Calculus

> d $\cosh (\theta) / d \theta=\sinh (\theta)$
> d $\sinh (\theta) / \mathbf{d} \theta=\cosh (\theta)$
> $d \tanh (\theta) / d \theta=1-\tan ^{2}(\theta)$

## Euclidian to Hyperbolic

$$
\begin{aligned}
& \cos (\theta)=\cosh (i \theta) \\
& \sin (\theta)=i \sinh (i \theta) \\
& \tan (\theta)=i \tanh (i \theta) \\
& \cot (\theta)=-i \operatorname{coth}(i \theta) \\
& \sec (\theta)=\operatorname{sech}(i \theta) \\
& \csc (\theta)=-i \operatorname{csch}(i \theta)
\end{aligned}
$$

Hyperbolic to Euclidian

$$
\begin{aligned}
& \cosh (\theta)=\cos (i \theta) \\
& \sinh (\theta)=-i \sin (i \theta) \\
& \tanh (\theta)=-i \tan (i \theta) \\
& \operatorname{coth}(\theta)=i \tan (i \theta) \\
& \operatorname{sech}(\theta)=\sec (i \theta) \\
& \operatorname{csch}(\theta)=i \csc (i \theta)
\end{aligned}
$$

## Rapidity

The dimensionless velocities
$\beta_{1}=v_{1} / c$
$\beta_{2}=v_{2} / \mathbf{c}$

Rapidity
$\beta=\boldsymbol{\operatorname { t a n h }}(\mathbf{r})$

Rapidity adds
$r_{3}=r_{1}+r_{2}$
$\beta=\tanh \left(\mathbf{r}_{1}+\mathbf{r}_{\mathbf{2}}\right)$
$\beta=\left(\tanh \left(r_{1}\right)+\tanh \left(r_{2}\right)\right) /\left(1+\tanh \left(r_{1}\right) \tanh \left(r_{2}\right)\right)$

Einstein velocity addition

$$
\beta=\left(\beta_{1}+\beta_{2}\right) /\left(1+\beta_{1} \beta_{2}\right)
$$



Two events simultaneous in one frame are not simultaneous in any other frame

Now that we have plotted (a region of the) inertial reference frame $S^{\prime}$ moving at $v_{\text {rel }}=0.5$ relative to $S$, we can test the statement regarding the relativity of simultaneity by graphical means.


Figure 3-II. 2 Relativity of simultaneity.

For an observer in $S$, we represent simultaneous events A and B occurring at time $t$, by plotting a line parallel to the $x$ axis through points A and B . For an observer in $S^{\prime}$, similar lines plotted parallel to the $x^{\prime}$ axis through point A and B correspond with time measurements $A^{\prime}$ and $B^{\prime}$ on the $t^{\prime}$ axis. Not only are events A and B not simultaneous in $S^{\prime}$, their chronological order is reversed.

## 3-1. Relativity and swimming

The idea here is to illustrate how remarkable is the invariance of the speed of light (speed of light is the same in all free-float frames) by contrasting it with the case of a swimmer making her way through water.

Light goes through space at $3 \times 10^{8}$ meters per second, and the swimmer goes through water at 1 meter per second. "But how can there otherwise be any difference?" one at first asks oneself.

For a light flash to go down the length of a 30-meter spaceship and back again it takes

$$
\begin{aligned}
\text { time } & =(\text { distance }) /(\text { speed }) \\
& =2 \times(30 \text { metes }) /\left(3 \times 10^{8} \text { meters } / \text { second }\right) \\
& =2 \times 10^{-7} \text { second }
\end{aligned}
$$



Two events at the same place in one frame are not at the same place in any other frame

## Minkowski diagrams

Two great things
(1) Arbitrarily many reference frames
(2) Constantly reminded that space is hyperbolic

One not so great thing
(1) Hyperbolic distortion


## Time Dilation



## Length Contraction






## $X^{\prime}$ is Moving Right



Figure L-VI. $1 \boldsymbol{x}^{\prime}$ moving to the right of $\boldsymbol{x}$.

## $X$ is Moving Left



Figure L-VI. $2 \boldsymbol{x}$ moving to the left of $\boldsymbol{x}^{\prime}$.

## Lorentz transformation matrix $\Lambda$

$$
\begin{aligned}
& \boldsymbol{\Lambda}=\left[\begin{array}{cccc}
\cosh \theta & \sinh \theta & 0 & 0 \\
\sinh \theta & \cosh \theta & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1
\end{array}\right] \\
& \boldsymbol{\Lambda}=\left[\begin{array}{cccc}
\gamma & \gamma \beta & 0 & 0 \\
\gamma \beta & \gamma & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1
\end{array}\right]
\end{aligned}
$$

## Cartesian

$$
\binom{x}{y}=\left(\begin{array}{rr}
\cos \theta & \sin \theta \\
-\sin \theta & \cos \theta
\end{array}\right)\binom{x^{\prime}}{y^{\prime}}
$$

## Hyperbolic

$$
\binom{x}{c t}=\left(\begin{array}{cc}
\cosh \beta & \sinh \beta \\
\sinh \beta & \cosh \beta
\end{array}\right)\binom{x^{\prime}}{c t^{\prime}}
$$

# Making Photons using Relativistic Electrons 

Synchrotrons
Wigglers
Undulators
and
X-Ray Lasers


Figure 10: The electric field lines of a point charge.

## Compression of the Electric Field Lines



## Before $\beta=0.95$

After $\beta=0$

radial lines $\sim \mathbf{r a}^{\mathbf{- 2}}$
tangential lines $\sim \mathbf{r}^{-1}$

At rest, then acceleration a for time $\tau$, finally at fixed velocity $v$


From the Geometry
Transverse component $=\mathrm{vt} \sin \theta$ Radial component $=\mathbf{c} \tau$


## Transverse Magnetic Field




## Low Velocity Limit Non-Relativistic Limit



## Dipole Pattern $\beta=0.90$



# Relativistic with <br> Acceleration Parallel to the Velocity 



## Synchrotron Radiation $\beta=0.95$



Synchrotrons are Ultra-Relativistic with Acceleration Perpendicular to the Velocity


## Low Velocity Limit Non-Relativistic Limit



## Electron Frame



Lab Frame


## Bright and Powerful X-Rays from Relativistic Electrons



Synchrotron radiation

- $10^{10}$ brighter than the
most powerful (compact) laboratory source
- An x-ray "light bulb" in that it radiates all "colors" (wavelengths, photons energies)


Undulator radiation

- Lasers exist for the IR, visible, UV, VUV, and EUV
- Undulator radiation is quasimonochromatic and highly directional, approximating many of the desired properties of an x-ray laser


## Synchrotron Radiation in a Narrow Forward Cone

Frame moving with electron
Laboratory frame of reference


$$
\begin{gather*}
\tan \theta=\frac{\sin \theta^{\prime}}{\gamma\left(\beta+\cos \theta^{\prime}\right)}  \tag{5.1}\\
\theta \simeq \frac{1}{2 \gamma} \tag{5.2}
\end{gather*}
$$

## Relativistic Electrons Radiate in a Narrow Forward Cone

Dipole radiation


Frame of reference moving with electrons


## Three Forms of Synchrotron Radiation



## Narrow Cone Undulator Radiation, Generated by Relativistic Electrons Traversing a Periodic Magnet Structure



## An Undulator Up Close



ALS U5 undulator, beamline $7.0, \mathrm{~N}=89, \lambda_{\mathrm{u}}=50 \mathrm{~mm}$

## Undulator Radiation


$\mathrm{E}=\gamma \mathrm{mc}^{2}$
$\gamma=\frac{1}{\sqrt{1-\frac{v^{2}}{c^{2}}}}$
$N=$ \# periods

$\mathrm{e}^{-}$radiates at the Lorentz contracted wavelength:

$$
\lambda^{\prime}=\frac{\lambda_{\mathrm{u}}}{\gamma}
$$

Bandwidth:

$$
\frac{\lambda^{\prime}}{\Delta \lambda^{\prime}} \simeq \mathrm{N}
$$

## Frame of Observer



Doppler shortened wavelength on axis:

$$
\begin{aligned}
& \lambda=\lambda^{\prime} \gamma(1-\beta \cos \theta) \\
& \lambda=\frac{\lambda_{u}}{2 \gamma^{2}}\left(1+\gamma^{2} \theta^{2}\right)
\end{aligned}
$$

Accounting for transverse motion due to the periodic magnetic field:

$$
\begin{aligned}
& \lambda=\frac{\lambda_{u}}{2 \gamma^{2}}\left(1+\frac{\mathrm{K}^{2}}{2}+\gamma^{2} \theta^{2}\right) \\
& \text { where } \mathrm{K}=\mathrm{eB}_{0} \lambda_{\mathrm{u}} / 2 \pi \mathrm{mc}
\end{aligned}
$$

## Following Monochromator

For $\frac{\Delta \lambda}{\lambda} \simeq \frac{1}{N}$

$$
\theta_{\text {cen }} \simeq \frac{1}{\gamma \sqrt{N}}
$$

typically

$$
\theta_{\text {cen }} \simeq 40 \mathrm{rad}
$$

# The Synchrotrons that <br> I Have Loved 














