

**Figure 41.5.**

“Orientation-entanglement relation” between a cube and the walls of a room. A  $360^\circ$  rotation of the cube entangles the threads. A  $720^\circ$  rotation might be thought to entangle them still more—but instead makes it possible completely to disentangle them.

### §41.5. SPINORS

Paint each face of a cube a different color. Then connect each corner of the cube to the corresponding corner of the room with an elastic thread (Figure 41.5). Now rotate the cube through  $2\pi = 360^\circ$ . The threads become tangled. Nothing one can do will untangle them. It is impossible for every thread to proceed on its way in a straight line. Now rotate the cube about the same axis by a further  $2\pi$ . The threads become still more tangled. However, a little work now completely straightens out the tangle (Figure 41.6). Every thread runs as it did in the beginning in a straight line from its corner of the cube to the corresponding corner of the room. More generally, rotations by  $0, \pm 4\pi, \pm 8\pi, \dots$ , leave the cube in its standard “orientation-entanglement relation” with its surroundings, whereas rotations by  $\pm 2\pi, \pm 6\pi, \pm 10\pi, \dots$ , restore to the cube only its orientation, not its orientation-entanglement relation with its surroundings. Evidently there is something about the geometry of orientation that is not fully taken into account in the usual concept of orientation; hence the concept of “orientation-entanglement relation” or (briefer term!) “version” (Latin *versor*, turn). Whether there is also a detectable difference in the physics (contact potential between a metallic object and its metallic surroundings, for example) for two inequivalent versions of an object is not known [Aharonov and Susskind (1967)].

In keeping with the distinction between the two inequivalent versions of an object, the spin matrix associated with a rotation,

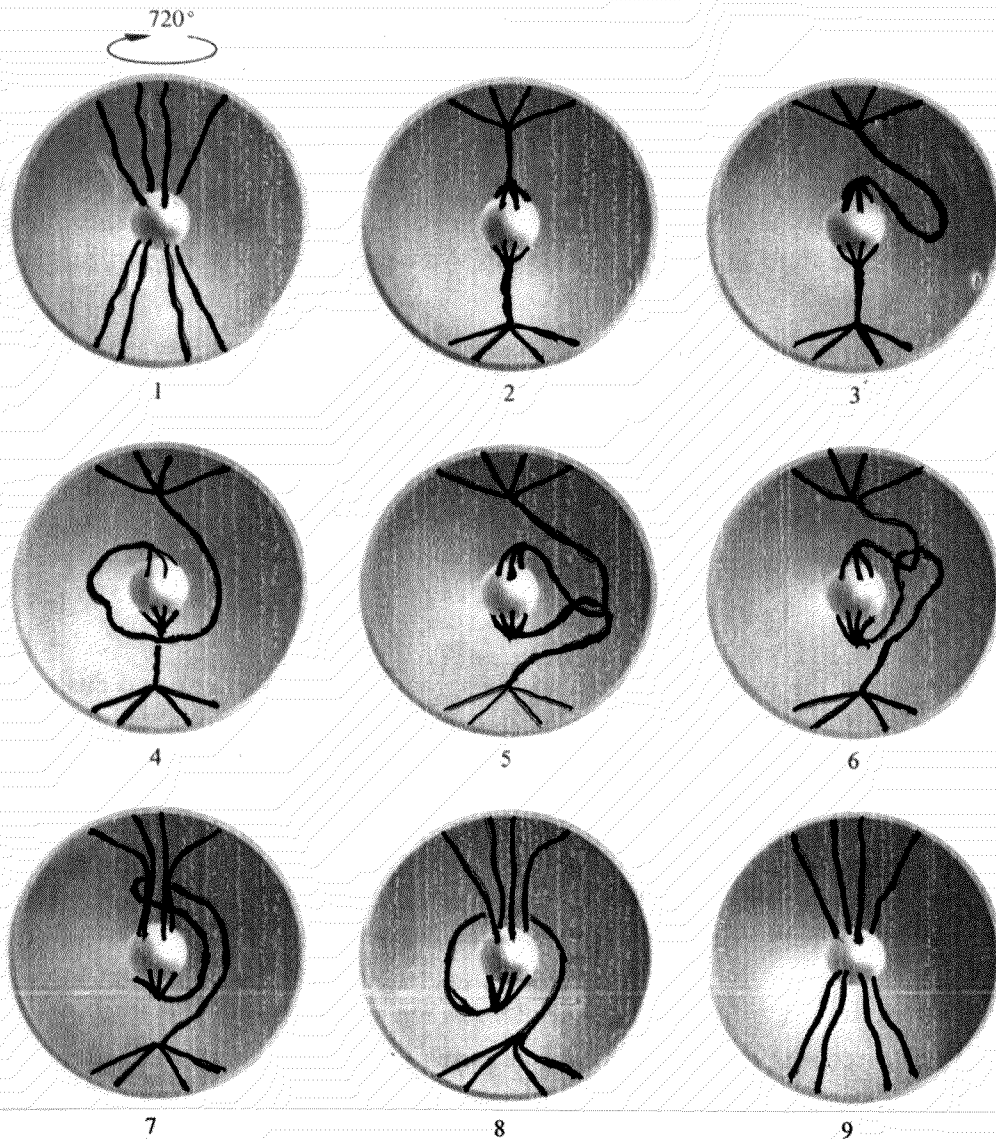
$$R = \cos(\theta/2) - i(\mathbf{n} \cdot \boldsymbol{\sigma}) \sin(\theta/2), \quad (41.48)$$

reverses sign on a rotation through an odd multiple of  $2\pi$ . This sign change never shows up in the law of transformation of a vector, as summarized in the formula

$$X \longrightarrow X' = RXR^* \quad (41.49)$$

(two factors  $R$ ; sign change in each!). The sign change does show up when one turns from a vector to a 2-component quantity that transforms according to the law

$$\xi \longrightarrow \xi' = R\xi. \quad (41.50)$$



**Figure 41.6.**

An object is connected to its surroundings by elastic threads as in Figure 41.5. (Eight are shown here; any number could be used.) Rotating the object through  $720^\circ$  and then following the procedure outlined (Edward McDonald) in frames 2–8 (with the object remaining fixed), one finds that the connecting threads are left disentangled, as in frame 9 (lower right).

Such a quantity is known as a spinor. A spinor reverses sign on a  $360^\circ$  rotation. It therefore provides a reasonable means to keep track of the difference between the two inequivalent versions of the cube. More generally, with each orientation-entanglement relation between the cube and its surroundings one can associate a different value of the spinor  $\xi$ . Moreover, there is nothing that limits the usefulness of the spinor concept to rotations. Also, for the general combination of boost and rotation, one can write

$$\xi \longrightarrow \xi' = L\xi.$$

(41.51) Lorentz transformation of a spinor