

$$\mathbf{a} = [\mathbf{a}, \mathbf{H}] = \mathbf{aH-Ha} \text{—oh, my, gotcha!!!}$$

The ladder operator solution to the simple harmonic oscillator problem is subtle, exquisite, and rather slippery—so I thought you might appreciate a recapitulation of what I said in class You might want to go through the argument line-by-line until it clicks!

There were three steps in the argument:

1. The first step was to show that the eigenvalues of the Hamiltonian H are equal to $\hbar\omega$ times $\frac{1}{2}$ plus the eigenvalues of the number operator $N = a^\dagger a$ (which will turn out to be n , so we will end up with $(n + \frac{1}{2})\hbar\omega$). We did this by showing that the Hamiltonian H is $\hbar\omega$ times the sum of the number operator plus one half the identity operator,

$$H = (a^\dagger a + \frac{1}{2}) \hbar\omega.$$

We showed this by defining the a^\dagger and a operators, and then calculating $a^\dagger a$. Note that once we found that $H = (a^\dagger a + \frac{1}{2}) \hbar\omega$, we immediately knew that the eigenvectors of H would be the same as the eigenvalues of $a^\dagger a$ —because every vector is an eigenvector of the identity operator! We also immediately knew that the eigenvalues of H would be equal to $\hbar\omega$ times the eigenvalues of the $a^\dagger a$ operator plus $\frac{1}{2}\hbar\omega$.

2. The second step was to show that when the a^\dagger and a operators act on any eigenvector of H , we get back another eigenvector of H one step up or down the ladder of states. We showed this by calculating the three commutators:

$$\begin{aligned} [a, a^\dagger] &= +1 \\ [a, H] &= +a \\ [a^\dagger, H] &= -a^\dagger \end{aligned}$$

and considering the action of the last two commutators on any eigenvector of the Hamiltonian

$$\begin{aligned} [a, H] |\text{eigenvector of } H\rangle &= +a |\text{eigenvector of } H\rangle \\ [a^\dagger, H] |\text{eigenvector of } H\rangle &= -a^\dagger |\text{eigenvector of } H\rangle. \end{aligned}$$

By expanding the commutators, we found

$$\begin{aligned} (aH - Ha) |\text{eigenvector of } H\rangle &= +a |\text{eigenvector of } H\rangle \\ (a^\dagger H - Ha^\dagger) |\text{eigenvector of } H\rangle &= -a^\dagger |\text{eigenvector of } H\rangle. \end{aligned}$$

which allowed us to conclude that

$$\begin{aligned} a |\epsilon\rangle &= (\epsilon - 1) |\epsilon - 1\rangle \\ a^\dagger |\epsilon\rangle &= (\epsilon + 1) |\epsilon + 1\rangle. \end{aligned}$$

This showed us that the eigenvalues of H are separated by $\pm\hbar\omega$. Combining this with the $\frac{1}{2}\hbar\omega$ from step one, we then concluded that the eigenvalues of the Hamiltonian are given by

$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega$$

where n is any integer (*i.e.*, positive, negative, or zero!!!). However, in step three, we found the smallest eigenvalue of the number operator is equal to zero.

So the only other thing we did not know yet was whether the raising and lowering operators return normalized eigenvectors of the Hamiltonian, *i.e.*, are the vectors $a|\epsilon\rangle$ and $a^\dagger|\epsilon\rangle$ normalized eigenvectors of H ? We did know that they are eigenvectors of H with eigenvalues of $(\epsilon-1)\hbar\omega$ and $(\epsilon+1)\hbar\omega$, respectively, but we did not know whether they are normalized—and, in fact, they are not!

3. The third step was to calculate the normalization coefficients. To do this we started with two adjacent normalized states, $|n\rangle \equiv |E = (n + \frac{1}{2})\hbar\omega\rangle$ and $|n-1\rangle \equiv |E = ((n-1) + \frac{1}{2})\hbar\omega\rangle$ and then we calculated the expectation value of the number operator in two different ways:

(i) First, we started with the lowering operator equation

$$a|n\rangle = c_n|n-1\rangle$$

and then we calculated the adjoint of this equation

$$\langle n| a^\dagger = \langle n-1| c_n^*.$$

We combined these to evaluate the expectation value of the number operator

$$\begin{aligned} \langle n| a^\dagger a |n\rangle &= \langle n-1| c_n^* c_n |n-1\rangle = |c_n|^2 \langle n-1|n-1\rangle \\ &= |c_n|^2. \end{aligned}$$

(ii) Second, we replaced $a^\dagger a$ by $\hat{H} - \frac{1}{2}$ and recalculated the expectation value of N

$$\begin{aligned} \langle n| \hat{H} - \frac{1}{2} |n\rangle &= \langle n-1| \left[\left(n + \frac{1}{2}\right) + \frac{1}{2} \right] |n\rangle = n \langle n|n\rangle \\ &= n. \end{aligned}$$

By combining these two calculations, we found

$$\begin{aligned} |c_n|^2 = n &\Rightarrow c_n = \sqrt{n} \\ \Rightarrow a|n\rangle &= \sqrt{n}|n-1\rangle. \end{aligned}$$

Finally, to see that the lowest eigenvalue of the number operator is zero, we considered

$$a|0\rangle = \sqrt{0}|0-1\rangle = 0|-1\rangle = \vec{0}.$$

So $|0\rangle$ is the bottom rung on the ladder (lowering it we obtain the zero vector), and consequently the lowest eigenvalue of H is $\frac{1}{2}\hbar\omega$, which is the zero point energy of the oscillator.