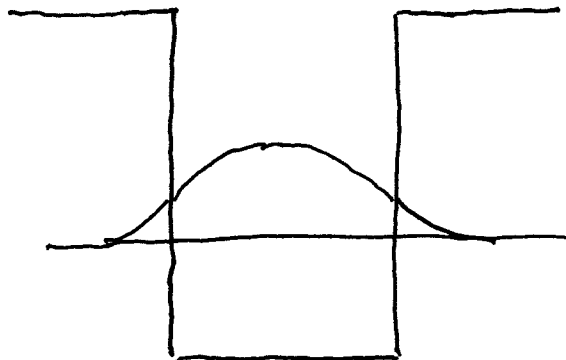
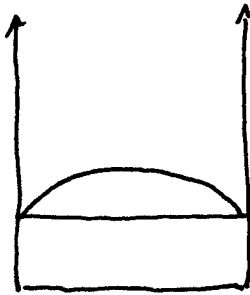


$$V(x) = \frac{1}{2} K x^2$$

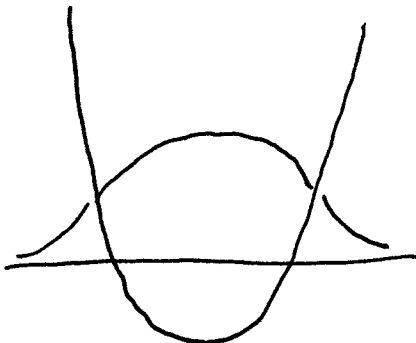
FIRST NON-TRIVIAL $V(x)$

WHAT SHOULD WE EXPECT?



$E > V$ *sin inside*

$E < V$ *exp tails outside*

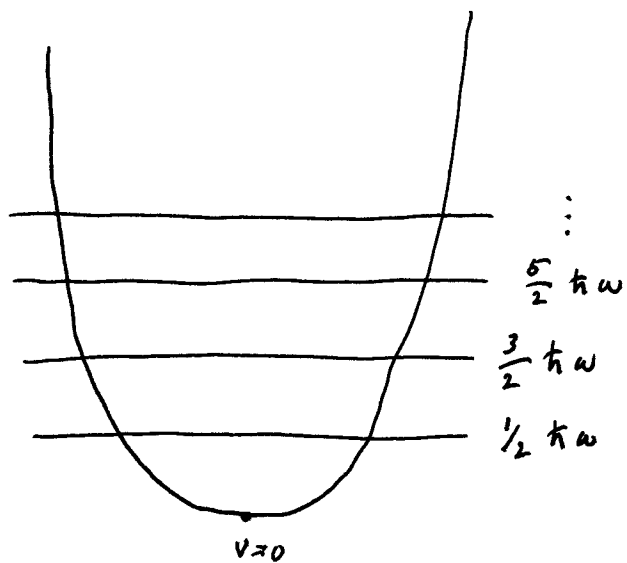


inside $E > V$ but variable

outside $E < V$ but variable

ground state \Rightarrow no nodes

WHAT WE WILL FIND



EQUALLY SPACED ENERGY LEVELS

$$\Delta E = \hbar \omega$$

GROUND STATE $E = \frac{1}{2} \hbar \omega$

STATIONARY STATES

$$\Psi_m(x) = (\text{m-TH ORDER POLYNOMIAL}) (\text{GAUSSIAN})$$

HERMITE
POLYNOMIALS

ASYMPTOTIC
DECAY

$$= A_m H_m(x) e^{-cx^2/2}$$

$$\Psi_m(x) = A_m H_m(y) e^{-y^2/2}$$

CARTOON VERSION



x



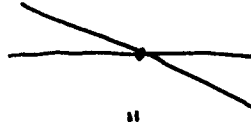
"



$\psi_0(x)$



x



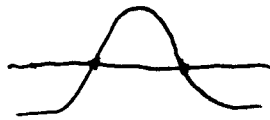
"



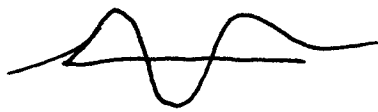
$\psi_1(x)$



x



"



$\psi_3(x)$

UBIQUITOUS

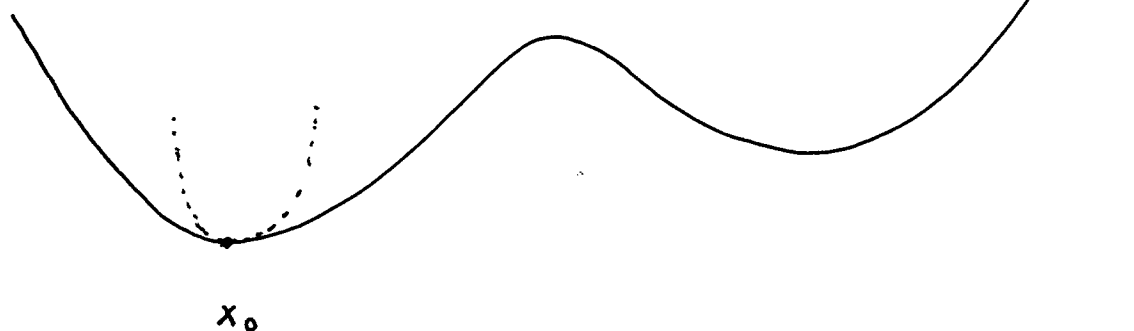
DIATOMIC MOLECULES

PHONONS

ATOMS IN SOLIDS

PHOTONS

PLASMONS



$$V(x) = V_0 + \left. \frac{\partial V}{\partial x} \right|_{x_0} (x-x_0) + \frac{1}{2!} \left. \frac{\partial^2 V}{\partial x^2} \right|_{x_0} (x-x_0)^2$$

$$V_0 = 0$$

$$x_0 = 0$$

$$V(x) = \frac{1}{2} \left. \frac{\partial^2 V}{\partial x^2} \right|_0 x^2$$

$$= \frac{1}{2} k x^2$$

LECTURE 18

JULY 23, 2007

SOLVE SHO DIER BA

$$H |\psi_m\rangle = E_m |\psi_m\rangle$$

or get: x -space, dot with $\langle x |$ bra

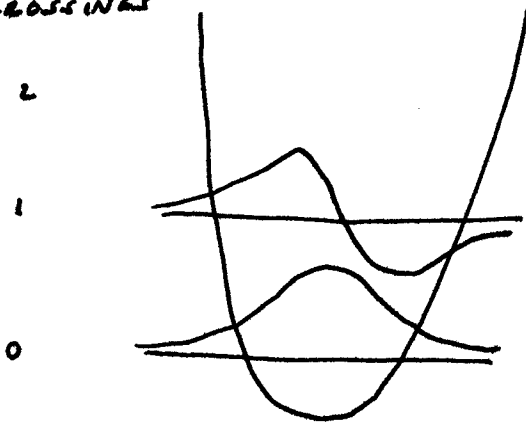
$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} kx^2 \right] \psi_m(x) = E_m \psi_m(x)$$

Find en's E_m

Find e f's $\psi_m(x)$

TELL 'EM WHAT YOU'RE GONNA TELL 'EM

ZERO
CROSSINGS



FIND ASYPTOTIC FORM OUTSIDE

= GAUSSIAN

SOLUTION = GAUSSIAN \times POLYNOMIALS

↑
HERMITE
POLYNOMIALS

TODAY : SHO-2

POWER SERIES,
ORTHOGONAL POLYNOMIALS,
DIFFERENTIAL EQUATIONS,
AND ALL THAT....

WE WANT TO SOLVE TQSE

$$H|\psi\rangle = i\hbar \frac{d}{dt}|\psi\rangle$$

SO WE FIRST SOLVE TISE

$$H|E_m\rangle = E_m|E_m\rangle$$

THEN USE

$$\begin{aligned} |\psi(t)\rangle &= \sum_m |E_m\rangle \langle E_m| e^{-iE_m t/\hbar} |\psi(0)\rangle \\ &= \sum_m |E_m\rangle \langle E_m|\psi(0)\rangle e^{-iE_m t/\hbar} \end{aligned}$$

SO, TO SOLVE TISE...

13-82 500 SHEETS FILLER 5 SQUARE
13-83 100 SHEETS FILLER 3 SQUARE
13-84 100 SHEETS FILLER 3 SQUARE
13-85 200 SHEETS EYE-EASE 5 SQUARE



Model No. 5-A

LECTURE 11: SHO II

POWER SERIES, DIFF EQ'S, AND ALL THAT

TISE

$$H | \varphi_m \rangle = E_m | \varphi_m \rangle$$

using position basis:

$$\left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} k x^2 \right] \varphi_m(x) = E_m \varphi_m(x)$$

OUR JOB: FIND φ_m 'SFIND E_m 'S

REWRITE TISE

$$\frac{d^2 \psi}{dx^2} + (\lambda - \alpha^2 x^2) \psi = 0$$

$$\lambda = \frac{2mE}{\hbar^2}$$

$$\alpha = \frac{m\omega}{\hbar}$$

FIRST STEP: FIND ASYMPTOTIC SOLUTION

FOR LARGE x , $\alpha^2 x^2 \gg \lambda$

$$\frac{d^2 \psi}{dx^2} - \alpha^2 x^2 \psi = 0$$

$$\psi = A e^{-\frac{1}{2} \alpha x^2} + B e^{+\frac{1}{2} \alpha x^2}$$

$$\frac{d\psi}{dx} = -\alpha x A e^{-\frac{1}{2} \alpha x^2} + \alpha x B e^{+\frac{1}{2} \alpha x^2}$$

$$\frac{d^2 \psi}{dx^2} = \alpha^2 x^2 A e^{-\frac{1}{2} \alpha x^2} + \alpha^2 x^2 B e^{+\frac{1}{2} \alpha x^2}$$

$$- \cancel{\alpha A e^{-\frac{1}{2} \alpha x^2}} + \cancel{\alpha B e^{+\frac{1}{2} \alpha x^2}}$$

$\alpha^2 x^2 \gg \alpha$

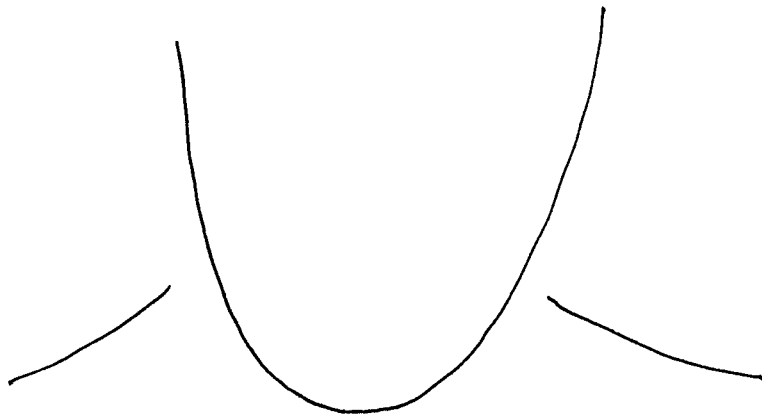
$$\psi(x) = A e^{-\frac{1}{2} \alpha x^2}$$

$$\psi(x) = B e^{+\frac{1}{2} \alpha x^2}$$



not square integ

SO, AT LARGE x , WE HAVE GAUSSIAN DECAY



LOOK FOR SOLUTIONS OF THE FORM

$$\psi(x) = e^{-\frac{1}{2}\alpha x^2} f(x)$$

$$\frac{d^2 \psi}{dx^2} + (\lambda - \alpha^2 x^2) \psi = 0$$

$$\frac{d\psi}{dx} = \frac{d}{dx} \left[e^{-\frac{1}{2}\alpha x^2} f(x) \right]$$

$$= -\alpha x e^{-\frac{1}{2}\alpha x^2} f(x) + e^{-\frac{1}{2}\alpha x^2} \frac{df}{dx}$$

$$\frac{d^2 \psi}{dx^2} = -\alpha e^{-\frac{1}{2} \alpha x^2} f(x) + \alpha^2 x^2 e^{-\frac{1}{2} \alpha x^2} f(x)$$

$$- \alpha x e^{-\frac{1}{2} \alpha x^2} \frac{df}{dx} - \alpha x e^{-\frac{1}{2} \alpha x^2} \frac{df}{dx}$$

$$+ e^{-\frac{1}{2} \alpha x^2} \frac{d^2 f}{dx^2}$$

$$\frac{d^2 \psi}{dx^2} \Rightarrow \left[\frac{d^2 f}{dx^2} - 2\alpha x \frac{df}{dx} + (\cancel{\alpha^2 x^2} - \alpha) f \right] e^{-\frac{1}{2} \alpha x^2}$$

$$+ (\lambda - \alpha^2 x^2) \psi = 0$$

$$+ (\lambda - \cancel{\alpha^2 x^2}) f e^{-\frac{1}{2} \alpha x^2} = 0$$

$$0 = \left[\frac{d^2 f}{dx^2} - 2\alpha x \frac{df}{dx} + (\lambda - \alpha) f \right] e^{-\frac{1}{2} \alpha x^2}$$

$$\xi = \sqrt{\alpha} x$$

$$f(x) \rightarrow H(\xi)$$

$$\frac{d}{dx} f(x) = \frac{d}{dx} H(\xi) = \frac{d}{d\xi} H(\xi) \frac{d\xi}{dx}$$

$$= \frac{dH(\xi)}{d\xi} \sqrt{\alpha}$$

$$\frac{d^2 A}{dx^2} = \alpha \frac{d^2 H}{d\xi^2}$$

$$\frac{d^2 f}{dx^2} - 2\alpha x \frac{df}{dx} + (\lambda - \alpha) f = 0$$

$$\frac{1}{\alpha} \left[\alpha \frac{d^2 H}{d\xi^2} - 2\alpha x \frac{dH}{d\xi} \sqrt{\alpha} + (\lambda - \alpha) H \right] = 0$$

$$\xi = \sqrt{\alpha} x$$

$$\frac{d^2 H}{d\xi^2} - 2\xi \frac{dH}{d\xi} + \left(\frac{\lambda}{\alpha} - 1 \right) H = 0$$

HERMITE'S EQN

SOLVE USING POWER SERIES...

$$H(\xi) = a_0 + a_1 \xi + a_2 \xi^2 + a_3 \xi^3 + \dots$$

$$\frac{dH}{d\xi} = 0 + a_1 + 2a_2 \xi + 3a_3 \xi^2 + \dots$$

$$\frac{d^2H}{d\xi^2} = 0 + 0 + 1 \cdot 2 a_2 + 2 \cdot 3 a_3 \xi$$

KNOW

$$\frac{d^2H}{d\xi^2} - 2\xi \frac{dH}{d\xi} + \left(\frac{\lambda}{\alpha} - 1\right) H = 0 \quad \text{FOR ALL } \xi!$$

\Rightarrow each coefficient of ξ must vanish

$$\xi^0 \quad 1 \cdot 2 a_2 + \left(\frac{\lambda}{\alpha} - 1\right) a_0 = 0$$

$$\xi^1 \quad 2 \cdot 3 a_3 + \left(\frac{\lambda}{\alpha} - 1 - 2\right) a_1 = 0$$

$$\xi^2 \quad 3 \cdot 4 a_4 + \left(\frac{\lambda}{\alpha} - 1 - 2 \cdot 2\right) a_2 = 0$$

$$\xi^3 \quad 4 \cdot 5 a_5 + \left(\frac{\lambda}{\alpha} - 1 - 2 \cdot 3\right) a_3 = 0$$

\Rightarrow COEFF OF $\{^m$

$$(m+1)(m+2) a_{m+2} + \left(\frac{\lambda}{2} - 1 - 2m\right) a_m = 0$$

$$a_{m+2} = \frac{-\left(\frac{\lambda}{2} - 2m - 1\right)}{(m+2)(m+1)} a_m$$

RECURSION RELATION

$a_0 \Rightarrow$ all even COEFF'S

$a_1 \Rightarrow$ all odd COEFF'S

FOR EACH m :

$$\text{WHEN } \left(+\frac{\lambda}{2} - 2m - 1\right) = 0$$

all higher terms vanish!

SPECIAL VALUES OF $\lambda \Rightarrow$ FINITE POLYNOMIALS

$$\lambda = \frac{2mE}{\hbar^2}$$

\Rightarrow EIGENENERGIES!

TWO CLASSES OF SOLUTIONS

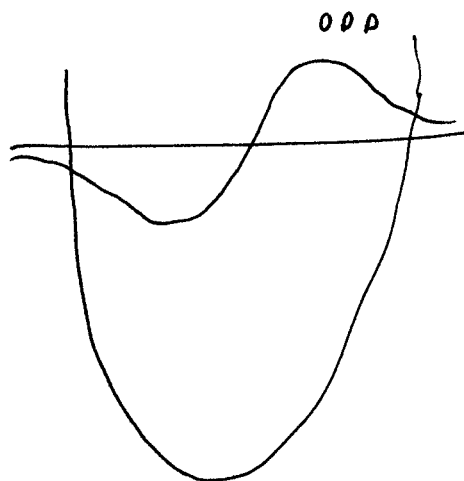
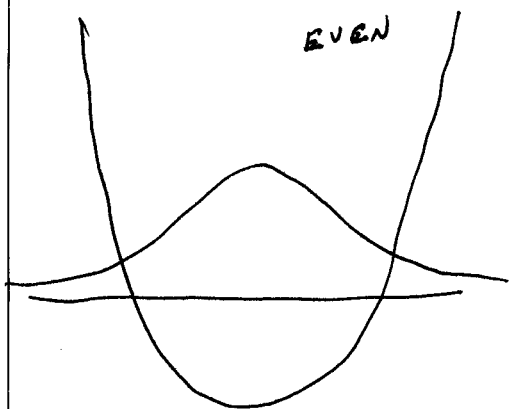
EVEN

$$H(\xi) = a_0 + a_2 \xi^2 + a_4 \xi^4 + \dots$$

ODD

$$H(\xi) = a_1 \xi + a_3 \xi^3 + a_5 \xi^5 + \dots$$

ODD AND EVEN COME FROM SYMMETRY OF
THE PROBLEM



N. B., FOR AN ARBITRARY ENERGY, NO STATIONARY

STATE! SERIES ^{DOES} NOT TERMINATE!

$$a_{m+2} = \frac{- \left[\frac{\lambda}{\alpha} - 2m - 1 \right]}{(m+2)(m+1)} a_m$$

CAN CHOOSE λ TO TERMINATE EVEN OR ODD SERIES,

BUT NOT BOTH.

LARGE m

$$a_{m+2} = \frac{- \left[\cancel{\frac{\lambda}{\alpha}} - 2m - \cancel{1} \right]}{(m+\cancel{2})(m+\cancel{1})} a_m$$

$$a_{m+2} = + \left(\frac{2}{m} \right) a_m$$

COMPARE WITH POWER SERIES FOR A

GAUSSIAN

$$e^{\xi^2} = \sum_{m=0}^{\infty} \frac{(\xi^2)^m}{m!} = 1 + \xi^2 + \frac{1}{2!} \xi^4 + \frac{1}{3!} \xi^6$$

$$+ \frac{\xi^m}{\left(\frac{m}{2}\right)!} + \frac{\xi^{m+2}}{\left(\frac{m+2}{2}\right)!}$$

$$\frac{b_{m+2}}{b_m} = \left(\frac{2}{m}\right) \Rightarrow \text{SAME!}$$

CALCULATE THE EIGENENERGIES

$$\left[\frac{\lambda}{\alpha} - 2m - 1 \right] = 0$$

$$\frac{\lambda}{\alpha} = 2m + 1$$

$$\lambda = (2m + 1) \alpha$$

$$\frac{2mE}{\hbar^2} = (2m + 1) \sqrt{\frac{mK}{\hbar^2}}$$

$$E = (2m + 1) \frac{\hbar^2}{2m} \sqrt{\frac{mK}{\hbar^2}}$$

$$= (m + \frac{1}{2}) \hbar \sqrt{\frac{K}{m}}$$

$$= (m + \frac{1}{2}) \hbar \omega$$

$$\begin{aligned}
H_0(\xi) &= 1 \\
H_1(\xi) &= 2\xi \\
H_2(\xi) &= 4\xi^2 - 2 \\
H_3(\xi) &= 8\xi^3 - 12\xi \\
H_4(\xi) &= 16\xi^4 - 48\xi^2 + 12 \\
H_5(\xi) &= 32\xi^5 - 160\xi^3 + 120\xi \\
H_6(\xi) &= 64\xi^6 - 480\xi^4 + 720\xi^2 - 120 \\
H_7(\xi) &= 128\xi^7 - 1344\xi^5 + 3360\xi^3 - 1680\xi \\
H_8(\xi) &= 256\xi^8 - 3584\xi^6 + 13440\xi^4 - 13440\xi^2 + 1680 \\
H_9(\xi) &= 512\xi^9 - 9216\xi^7 + 48384\xi^5 - 80640\xi^3 + 30240\xi \\
H_{10}(\xi) &= 1024\xi^{10} - 23040\xi^8 + 161280\xi^6 - 403200\xi^4 + 302400\xi^2 \\
&\quad - 30240.
\end{aligned}
\tag{11-23}$$

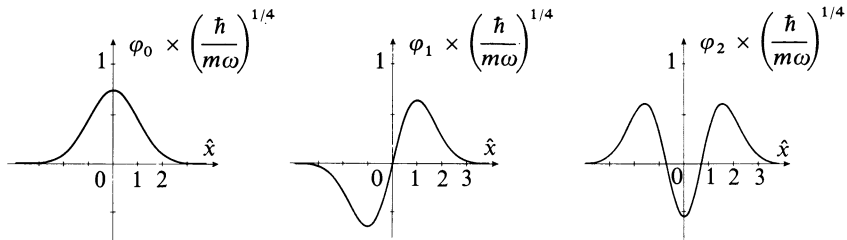


FIGURE 4

Wave functions associated with the first three levels of a harmonic oscillator.

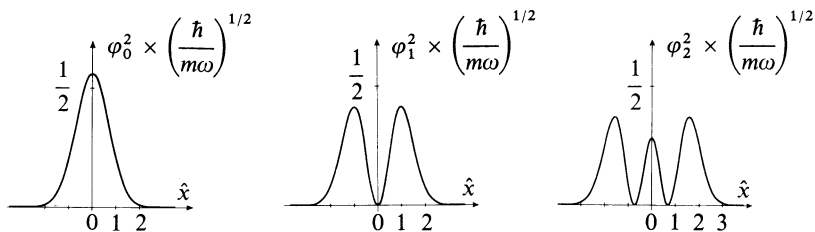


FIGURE 5

Probability densities associated with the first three levels of a harmonic oscillator.

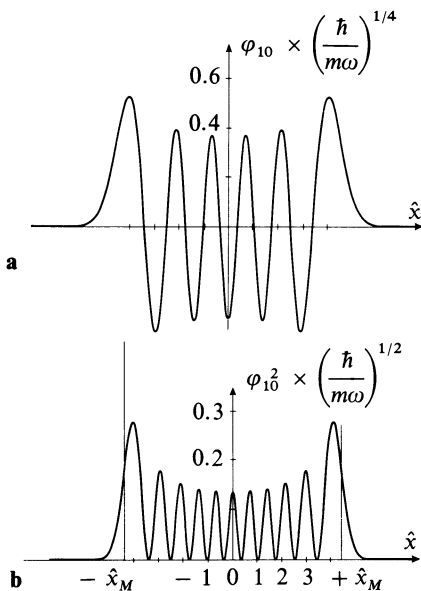


FIGURE 6

Shape of the wave function (fig. a) and of the probability density (fig. b) for the $n = 10$ level of a harmonic oscillator.

The Harmonic Oscillator

<http://hyperphysics.phy-astr.gsu.edu/hbase/quantum/hosc.html>

<http://www.falstad.com/qm1d/>

http://www-personal.umich.edu/~lorenz/java_applets/spaceholder/applets/SHO-QM-example.html?D1=5

<http://www.quantum-physics.polytechnique.fr/en/>

Hermite Polynomials

<http://mathworld.wolfram.com/HermitePolynomial.html>

<http://www.efunda.com/math/Hermite/index.cfm>

<http://www.sci.wsu.edu/idea/quantum/hermite.htm>

<http://functions.wolfram.com/Polynomials/>

Coherent States

<http://cat.sckans.edu/physics/Quantum%20Wave%20Packet.htm>