

### Overview of Topics

- ① Nonlinear (black body) radiation
- ② Heat Eq<sup>n</sup> in 2D & 3D
- ③ 2D Laplace?

# Heat Equation in 2D or 3D (or ND...!)

Same idea of conservation:

$$\text{Rate of change of heat energy (in time)} = \underset{\substack{\text{energy} \\ \text{Heat flux through boundary of volume}}}{\text{Heat flux through boundary of volume}} + \underset{\substack{\text{Heat energy generated} \\ \cancel{\text{within}} \text{ over entire volume}}}{\text{Heat energy generated over entire volume}}$$

Note: Now we are using a control volume!

(This should be familiar from the derivation of mass conservation)  
from L23 in MES64

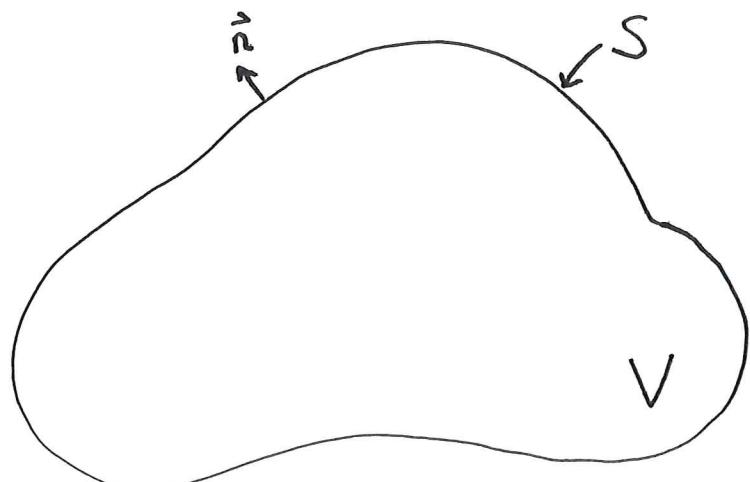
$$\frac{d}{dt} \iiint_V c_p u(\vec{x}, t) dV = - \iint_S \vec{q}_v(\vec{x}, t) \cdot \hat{n} dS + \iiint_V Q(\vec{x}, t) dV$$

Scalar field  
(single # at each spatial location  $\times$ )

Vector field:  $\vec{q}_v = \nabla u$

Scalar field

(continued)



$$\frac{d}{dt} \iiint_V c_p u(\vec{x}, t) dV = - \iint_S \vec{q}(\vec{x}, t) \cdot \vec{n} dS + \iiint_V Q(\vec{x}, t) dV$$

Now, use Gauss's Divergence Theorem to convert flux integral  
into volume integral of divergence!

$$\iint_S \vec{F} \cdot \vec{n} dS = \iiint_V \nabla \cdot \vec{F} dV$$

$$\frac{d}{dt} \iiint_V c_p u(\vec{x}, t) dV = - \iiint_V \nabla \cdot \vec{q}(\vec{x}, t) dV + \iiint_V Q(\vec{x}, t) dV \quad \left. \right\} \text{All Volume Integrals!}$$

$$\iiint_V \left[ c_p \frac{\partial u}{\partial t} + \nabla \cdot \vec{q} - Q \right] dV = 0$$

Since this is true for all volumes  $V$ :

$$c_p \frac{\partial u}{\partial t} + \nabla \cdot \vec{q} - Q = 0$$

Fourier's Law of  
Heat Conduction

$$\vec{q} = -K \nabla u$$

$$\begin{aligned} \nabla \cdot \vec{q} &= -K \nabla \cdot (\nabla u) = -K \begin{bmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{bmatrix} \cdot \begin{bmatrix} \frac{\partial u}{\partial x} \\ \frac{\partial u}{\partial y} \\ \frac{\partial u}{\partial z} \end{bmatrix} \\ &= -K [u_{xx} + u_{yy} + u_{zz}] \\ &= -K \nabla^2 u. \end{aligned}$$

$$c_p \frac{\partial u}{\partial t} - K \nabla^2 u - Q = 0$$

The general heat equation for a system with no sources ( $\mathbf{Q} = \mathbf{0}$ ) and constant thermal parameters ( $c, \rho, K$ ) is:

$$\boxed{\frac{\partial u}{\partial t} = \alpha^2 \nabla^2 u}$$

some positive #.

Equilibrium (Laplace!)

In 2D:  $u_t = \alpha^2 (u_{xx} + u_{yy})$

$$u_{xx} + u_{yy} = 0$$

In 3D:  $u_t = \alpha^2 (u_{xx} + u_{yy} + u_{zz})$

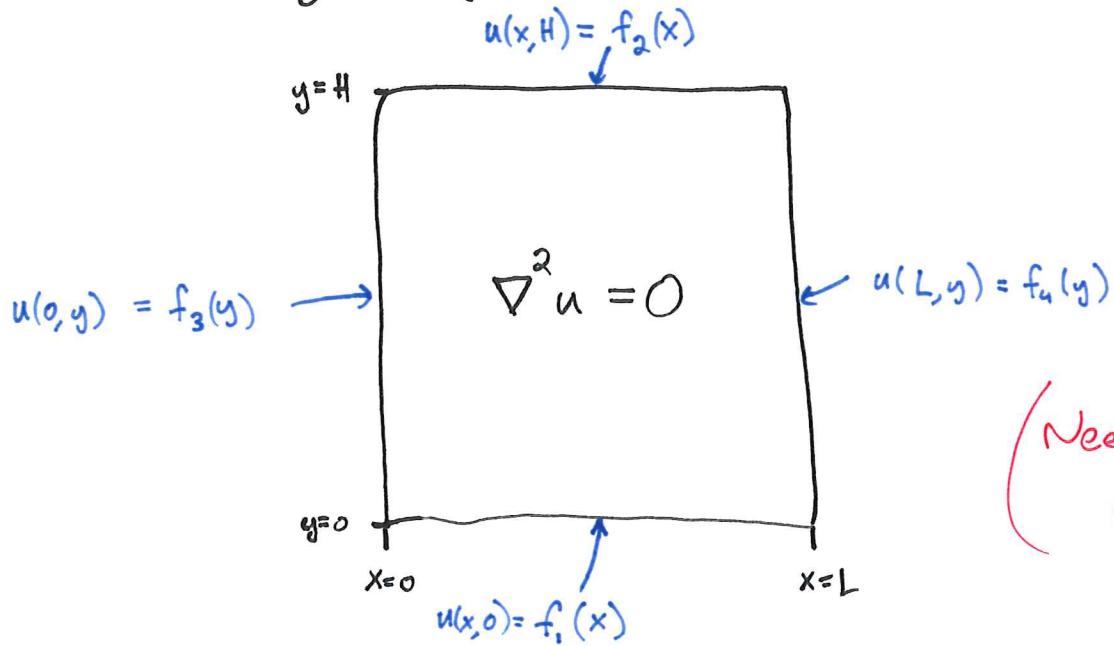
$$u_{xx} + u_{yy} + u_{zz} = 0$$

In ND:  $u_t = \alpha^2 \nabla^2 u$

$$\nabla^2 u = 0$$

$$u_t = \alpha^2 \sum_{k=1}^N \frac{\partial^2 u}{\partial x_k^2}$$

Let's try Laplace's Equation in 2D:



Because  $\nabla^2 u = 0$  is linear, we can solve four different systems with simpler BC's and sum the solutions:

$$\begin{array}{c}
 f_1 \quad f_2 \\
 \boxed{\phantom{f_1 f_2}} \quad f_3 \quad f_4 \\
 f_1 \quad f_2
 \end{array} = 
 \begin{array}{c}
 0 \quad 0 \\
 \boxed{\phantom{0 0}} \quad 0 \quad 0 \\
 0 \quad f_1 \quad 0
 \end{array} + 
 \begin{array}{c}
 f_2 \quad 0 \\
 0 \quad 0 \\
 0 \quad 0
 \end{array} + 
 \begin{array}{c}
 0 \quad 0 \\
 f_3 \quad 0 \quad 0 \\
 0 \quad 0
 \end{array} + 
 \begin{array}{c}
 0 \quad 0 \\
 0 \quad 0 \quad f_4 \\
 0 \quad 0
 \end{array}$$

$$u(x,y) = u_1(x,y) + u_2(x,y) + u_3(x,y) + u_4(x,y)$$

Let's work out one of these problems:

$$\begin{array}{c}
 0 \\
 \boxed{\nabla^2 u_3 = 0} \\
 0
 \end{array}$$