## Finite Volume Methods for Hyperbolic Problems

## Multidimensional Hyperbolic Problems

- Derivation of conservation law
- Hyperbolicity
- Advection
- Gas dynamics and acoustics
- Shear waves


## Derivation of conservation law

$$
\frac{d}{d t} \iint_{\Omega} q(x, y, t) d x d y=\text { net flux across } \partial \Omega
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Net flux is determined by integrating the flux of $q$ normal to $\partial \Omega$ around this boundary.

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$f(q)=$ flux of $q$ in the $x$-direction, $g(q)=$ flux of $q$ in the $y$-direction,
(both per unit length in orthog direction, per unit time),

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$\vec{n}(s)=\left(n^{x}(s), n^{y}(s)\right)$ outward-pointing unit normal $(x(s), y(s))$.
Flux at $(x(s), y(s))$ in the direction $\vec{n}(s)$ :

$$
\vec{n}(s) \cdot \vec{f}(q(x(s), y(s)))=f(q) n^{x}(s)+g(q) n^{y}(s)
$$

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If $q$ is smooth: divergence theorem $\Longrightarrow$

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\frac{d}{d t} \iint_{\Omega} q(x, y, t) d x d y=-\iint_{\Omega} \vec{\nabla} \cdot \vec{f}(q) d x d y
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where the divergence of $\vec{f}$ is

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True for any $\Omega \Longrightarrow \quad q_{t}+\vec{\nabla} \cdot \vec{f}(q)=0$. (PDE form)

## First order hyperbolic PDE in 2 space dimensions

General conservation law: $\quad q_{t}+f(q)_{x}+g(q)_{y}=0$
Quasi-linear form: $\quad q_{t}+f^{\prime}(q) q_{x}+g^{\prime}(q) q_{y}=0$

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Constant coefficient linear system: $q_{t}+A q_{x}+B q_{y}=0$
where $q \in \mathbb{R}^{m}, f(q)=A q, g(q)=B q$ and $A, B \in \mathbb{R}^{m \times m}$.

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Then plane wave propagating in any direction satisfies 1D hyperbolic equation.

## Plane wave solutions

Suppose

$$
\begin{aligned}
q(x, y, t) & =\breve{q}(x \cos \theta+y \sin \theta, t) \\
& =\breve{q}(\xi, t) .
\end{aligned}
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Then:

$$
\begin{aligned}
& q_{x}(x, y, t)=\cos \theta \breve{q}_{\xi}(\xi, t) \\
& q_{y}(x, y, t)=\sin \theta \breve{q}_{\xi}(\xi, t)
\end{aligned}
$$

SO

$$
q_{t}+A q_{x}+B q_{y}=\breve{q}_{t}+(A \cos \theta+B \sin \theta) \breve{q}_{\xi}
$$

and the 2 d problem reduces to the 1 d hyperbolic equation

$$
\breve{q}_{t}(\xi, t)+(A \cos \theta+B \sin \theta) \breve{q}_{\xi}(\xi, t)=0 .
$$

## Advection in 2 dimensions

Constant coefficient: $\quad q_{t}+u q_{x}+v q_{y}=0$
In this case solution for arbitrary initial data is easy:

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q(x, y, t)=q(x-u t, y-v t, 0)
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Data simply shifts at constant velocity $(u, v)$ in $x-y$ plane.

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Variable coefficient:
Conservation form: $\quad q_{t}+(u(x, y, t) q)_{x}+(v(x, y, t) q)_{y}=0$
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Equivalent only if flow is divergence-free (incompressible):

$$
\nabla \cdot \vec{u}=u_{x}(x, y, t)+v_{y}(x, y, t)=0 \quad \forall t \geq 0
$$

## Gas dynamics in 2D

$$
\begin{aligned}
& \rho(x, y, t)=\text { mass density } \\
& \rho(x, y, t) u(x, y, t)=x \text {-momentum density } \\
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If pressure $=P(\rho)$, e.g. isothermal or isentropic:

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\begin{aligned}
\rho_{t}+(\rho u)_{x}+(\rho v)_{y} & =0 \\
(\rho u)_{t}+\left(\rho u^{2}+p\right)_{x}+(\rho u v)_{y} & =0 \\
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For any $\theta$, the matrix $f^{\prime}(q) \cos \theta+g^{\prime}(q) \sin \theta$ has eigenvalues

$$
\breve{u}-c, \breve{u}, \breve{u}+c
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where $c=\sqrt{P^{\prime}(\rho)}$ and $\breve{u}=u \cos \theta+v \sin \theta$.

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Full Euler equations: 1 more equation for Energy
For any $\theta$, the matrix $f^{\prime}(q) \cos \theta+g^{\prime}(q) \sin \theta$ has eigenvalues

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\breve{u}-c, \breve{u}, \breve{u}+c \quad \text { Euler: another wave with } \lambda=\breve{u}
$$

where $c=\sqrt{P^{\prime}(\rho)}$ and $\breve{u}=u \cos \theta+v \sin \theta$.

## Solution of plane wave Riemann problem in 2D



Jump in $v$ from $v_{\ell}$ to $v_{r}$ propagates with the contact discontinuity

## Acoustics in 2 dimensions

Linearize about $u=0, v=0$ and $p=$ perturbation in pressure:

$$
\begin{aligned}
p_{t}+K_{0}\left(u_{x}+v_{y}\right) & =0 \\
\rho_{0} u_{t}+p_{x} & =0 \\
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Note: pressure responds to compression or expansion and so $p_{t}$ is proportional to divergence of velocity.

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Gives hyperbolic system $q_{t}+A q_{x}+B q_{y}=0$ with
$q=\left[\begin{array}{l}p \\ u \\ v\end{array}\right], \quad A=\left[\begin{array}{ccc}0 & K_{0} & 0 \\ 1 / \rho_{0} & 0 & 0 \\ 0 & 0 & 0\end{array}\right], \quad B=\left[\begin{array}{ccc}0 & 0 & K_{0} \\ 0 & 0 & 0 \\ 1 / \rho_{0} & 0 & 0\end{array}\right]$.

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Plane waves:

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A \cos \theta+B \sin \theta=\left[\begin{array}{ccc}
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Eigenvalues: $\lambda^{1}=-c_{0}, \quad \lambda^{2}=0, \quad \lambda^{3}=+c_{0}$
where $c_{0}=\sqrt{K_{0} / \rho_{0}}$ is independent of angle $\theta$.
Isotropic: sound propagates at same speed in any direction.

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where $c_{0}=\sqrt{K_{0} / \rho_{0}}$ is independent of angle $\theta$.
Isotropic: sound propagates at same speed in any direction.
Note: Zero wave speed for "shear wave" with variation only in velocity in direction $(-\sin \theta, \cos \theta)$.

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In this case, decouples into scalar advection equation for each component of $w$ :
$w_{t}^{p}+\lambda^{p} w_{x}^{p}+\mu^{p} w_{y}^{p}=0 \Longrightarrow w^{p}(x, y, t)=w^{p}\left(x-\lambda^{p} t, y-\mu^{p} t, 0\right)$.
Note: In this case information propagates only in a finite number of directions $\left(\lambda^{p}, \mu^{p}\right)$ for $p=1, \ldots, m$.

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This is not true for most coupled systems, e.g. acoustics.

## Acoustics in 2 dimensions

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\begin{aligned}
& p_{t}+K_{0}\left(u_{x}+v_{y}\right)=0 \\
& \rho_{0} u_{t}+p_{x}=0 \\
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& A=\left[\begin{array}{ccc}
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1 / \rho_{0} & 0 & 0 \\
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\end{array}\right], \quad R^{x}=\left[\begin{array}{rrr}
-Z_{0} & 0 & Z_{0} \\
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Solving $q_{t}+A q_{x}=0$ gives pressure waves in $(p, u)$.

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$$
B=\left[\begin{array}{ccc}
0 & 0 & K_{0} \\
0 & 0 & 0 \\
1 / \rho_{0} & 0 & 0
\end{array}\right] \quad R^{y}=\left[\begin{array}{rrr}
-Z_{0} & 0 & Z_{0} \\
0 & 1 & 0 \\
1 & 0 & 1
\end{array}\right]
$$

Solving $q_{t}+B q_{y}=0$ gives pressure waves in $(p, v)$.

