

Finite Volume Methods for Hyperbolic Problems

Linear Systems – Riemann Problems

- Riemann problems
- Riemann problem for advection
- Riemann problem for acoustics
- Phase plane

The Riemann problem

The **Riemann problem** consists of the hyperbolic equation under study together with initial data of the form

$$q(x, 0) = \begin{cases} q_l & \text{if } x < 0 \\ q_r & \text{if } x \geq 0 \end{cases}$$

Piecewise constant with a single jump discontinuity from q_l to q_r .

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The Riemann problem is fundamental to understanding

- The mathematical theory of hyperbolic problems,
- Godunov-type finite volume methods

Why? Even for nonlinear systems of conservation laws, the Riemann problem can often be solved for general q_l and q_r , and consists of a set of waves propagating at constant speeds.

The Riemann problem for advection

The **Riemann problem** for the advection equation $q_t + uq_x = 0$ with

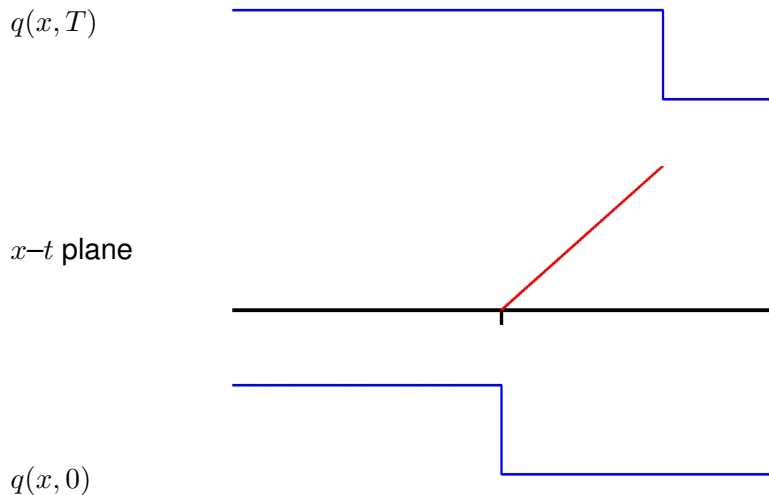
$$q(x, 0) = \begin{cases} q_l & \text{if } x < 0 \\ q_r & \text{if } x \geq 0 \end{cases}$$

has solution

$$q(x, t) = q(x - ut, 0) = \begin{cases} q_l & \text{if } x < ut \\ q_r & \text{if } x \geq ut \end{cases}$$

consisting of a single wave of strength $\mathcal{W}^1 = q_r - q_l$ propagating with speed $s^1 = u$.

Riemann solution for advection



Discontinuous solutions

Note: The Riemann solution is not a classical solution of the PDE $q_t + uq_x = 0$, since q_t and q_x blow up at the discontinuity.

Integral form:

$$\frac{d}{dt} \int_{x_1}^{x_2} q(x, t) dx = uq(x_1, t) - uq(x_2, t)$$

Integrate in time from t_1 to t_2 to obtain

$$\begin{aligned} \int_{x_1}^{x_2} q(x, t_2) dx - \int_{x_1}^{x_2} q(x, t_1) dx \\ = \int_{t_1}^{t_2} uq(x_1, t) dt - \int_{t_1}^{t_2} uq(x_2, t) dt. \end{aligned}$$

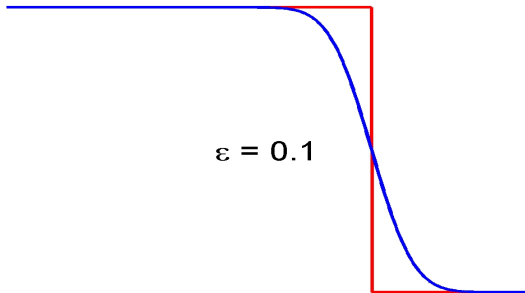
The Riemann solution satisfies the given initial conditions and this integral form for all $x_2 > x_1$ and $t_2 > t_1 \geq 0$.

Discontinuous solutions

Vanishing Viscosity solution: The Riemann solution $q(x, t)$ is the limit as $\epsilon \rightarrow 0$ of the solution $q^\epsilon(x, t)$ of the parabolic advection-diffusion equation

$$q_t + uq_x = \epsilon q_{xx}.$$

For any $\epsilon > 0$ this has a classical smooth solution:

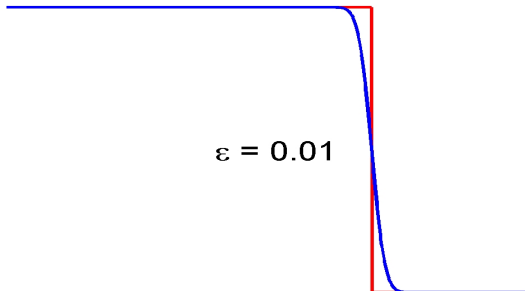


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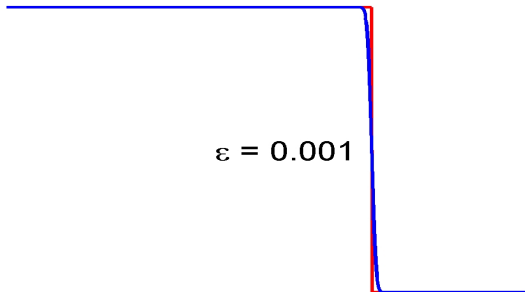


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Riemann Problems and Jupyter Solutions

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bookstore.siam.org/fa16/bonus

View static version of notebooks at:
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Eigenvectors for acoustics

$$A = \begin{bmatrix} 0 & K_0 \\ 1/\rho_0 & 0 \end{bmatrix}$$

Eigenvectors:

$$r^1 = \begin{bmatrix} -\rho_0 c_0 \\ 1 \end{bmatrix}, \quad r^2 = \begin{bmatrix} \rho_0 c_0 \\ 1 \end{bmatrix}.$$

Check that $A r^p = \lambda^p r^p$, where

$$\lambda^1 = -c_0, \quad \lambda^2 = +c_0.$$

with $c_0 = \sqrt{K_0/\rho_0} \implies K_0 = \rho_0 c_0^2$.

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Let $Z_0 = \rho_0 c_0 = \sqrt{K_0 \rho_0} =$ **impedance**.

Physical meaning of eigenvectors

Eigenvectors for acoustics:

$$r^1 = \begin{bmatrix} -\rho_0 c_0 \\ 1 \end{bmatrix} = \begin{bmatrix} -Z_0 \\ 1 \end{bmatrix}, \quad r^2 = \begin{bmatrix} \rho_0 c_0 \\ 1 \end{bmatrix} = \begin{bmatrix} Z_0 \\ 1 \end{bmatrix}.$$

Consider a pure 1-wave (simple wave), at speed $\lambda^1 = -c_0$,

If $\overset{\circ}{q}(x) = \bar{q} + \overset{\circ}{w}^1(x)r^1$ then

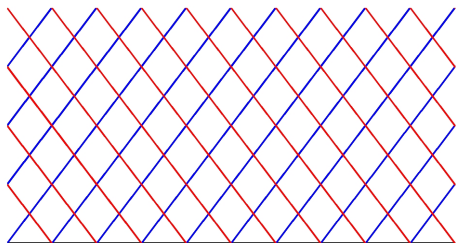
$$q(x, t) = \bar{q} + \overset{\circ}{w}^1(x - \lambda^1 t)r^1$$

Variation of q , as measured by q_x or $\Delta q = q(x + \Delta x) - q(x)$ is proportional to eigenvector r^1 , e.g.

$$q_x(x, t) = \overset{\circ}{w}_x^1(x - \lambda^1 t)r^1$$

Linear acoustics — characteristics

$$\begin{aligned} q(x, t) &= w^1(x + ct, 0)r^1 + w^2(x - ct, 0)r^2 \\ &= \frac{-\overset{\circ}{p}(x + ct)}{2Z_0} \begin{bmatrix} -Z_0 \\ 1 \end{bmatrix} + \frac{\overset{\circ}{p}(x - ct)}{2Z_0} \begin{bmatrix} Z_0 \\ 1 \end{bmatrix}. \end{aligned}$$



For IBVP on $a < x < b$, must specify one incoming boundary condition at each side: $w^2(a, t)$ and $w^1(b, t)$

Riemann Problem

Special initial data:

$$q(x, 0) = \begin{cases} q_l & \text{if } x < 0 \\ q_r & \text{if } x > 0 \end{cases}$$

Example: Acoustics with bursting diaphragm ($u_l = u_r = 0$)



Pressure:



Acoustic waves propagate with speeds $\pm c$.

Riemann Problem

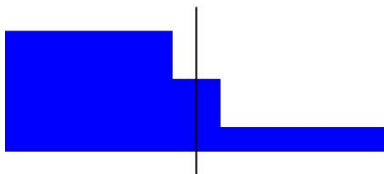
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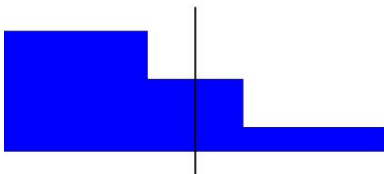
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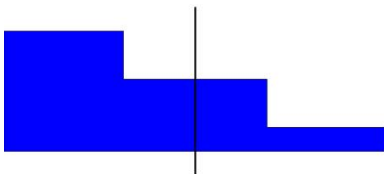
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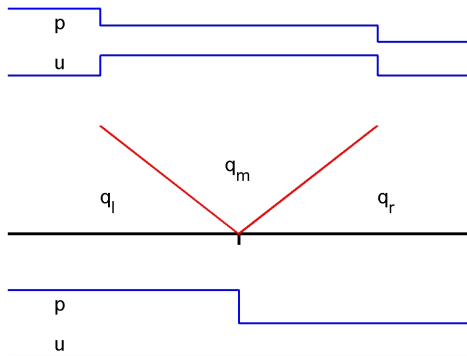
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Riemann Problem for acoustics

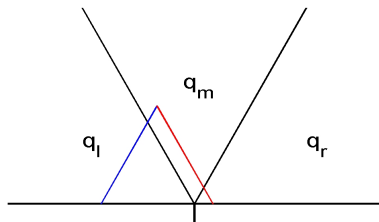
Waves propagating in $x-t$ space:



Left-going wave $\mathcal{W}^1 = q_m - q_l$ and
right-going wave $\mathcal{W}^2 = q_r - q_m$ are eigenvectors of A .

Riemann Problem for acoustics

In $x-t$ plane:



$$q(x, t) = w^1(x + ct, 0)r^1 + w^2(x - ct, 0)r^2$$

Decompose q_l and q_r into eigenvectors:

$$q_l = w_l^1 r^1 + w_l^2 r^2$$

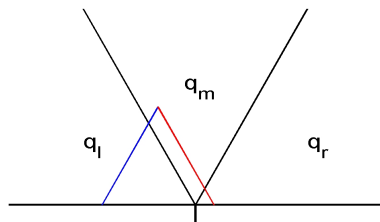
$$q_r = w_r^1 r^1 + w_r^2 r^2$$

Then

$$q_m = w_r^1 r^1 + w_l^2 r^2$$

Riemann Problem for acoustics

In $x-t$ plane:



Decompose $q_r - q_l$ into eigenvectors: Solve $R\alpha = \Delta q$

$$\begin{aligned}q_r - q_l &= (w_r^1 - w_l^1)r^1 + (w_r^2 - w_l^2)r^2 \\ &= \alpha^1 r^1 + \alpha^2 r^2 = \mathcal{W}^1 + \mathcal{W}^2.\end{aligned}$$

Then

$$\begin{aligned}q_m &= w_r^1 r^1 + w_l^2 r^2 \\ &= q_l + \alpha^1 r^1 = q_r - \alpha^2 r^2.\end{aligned}$$

Riemann solution for acoustics

$$r^1 = \begin{bmatrix} -\rho c \\ 1 \end{bmatrix} = \begin{bmatrix} -Z \\ 1 \end{bmatrix}, \quad r^2 = \begin{bmatrix} \rho c \\ 1 \end{bmatrix} = \begin{bmatrix} Z \\ 1 \end{bmatrix}.$$

Solving $R\alpha = \Delta q$ gives:

$$\alpha^1 = \frac{-\Delta p + Z\Delta u}{2Z}, \quad \alpha^2 = \frac{\Delta p + Z\Delta u}{2Z},$$

so

$$q_m = q_l + \alpha^1 r^1 = \frac{1}{2} \begin{bmatrix} (p_l + p_r) - Z(u_r - u_l) \\ (u_l + u_r) - (p_r - p_l)/Z \end{bmatrix}.$$

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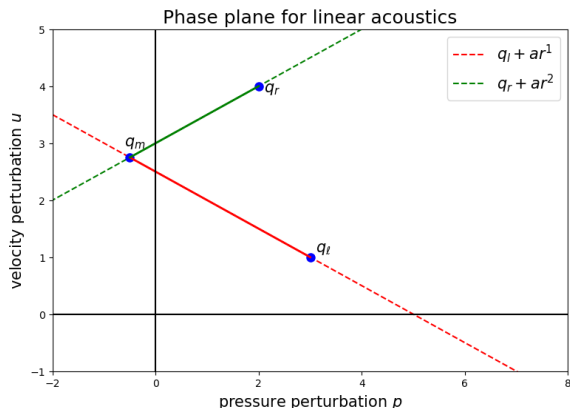
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Ex: shock tube with $u_l = u_r = 0$:

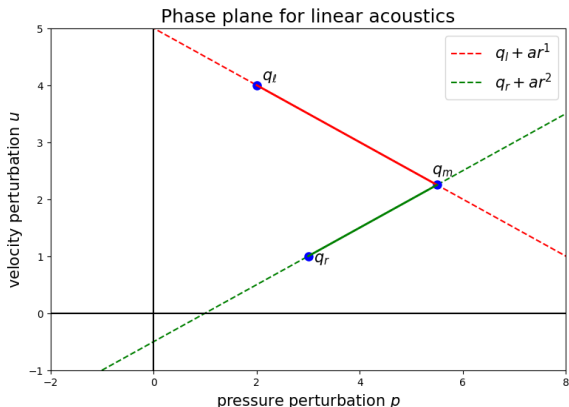
$$q_m = q_l + \alpha^1 r^1 = \frac{1}{2} \begin{bmatrix} (p_l + p_r) \\ -(p_r - p_l)/Z \end{bmatrix}.$$

Phase plane solution to Riemann problem



q_l and q_m are connected by a multiple of r^1
 q_m and q_r are connected by a multiple of r^2

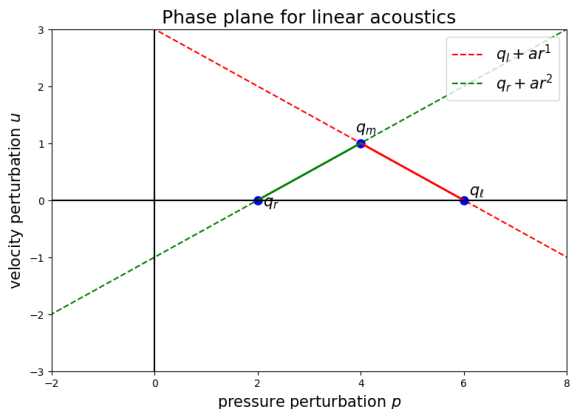
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Note that swapping q_l and q_r changes the solution!

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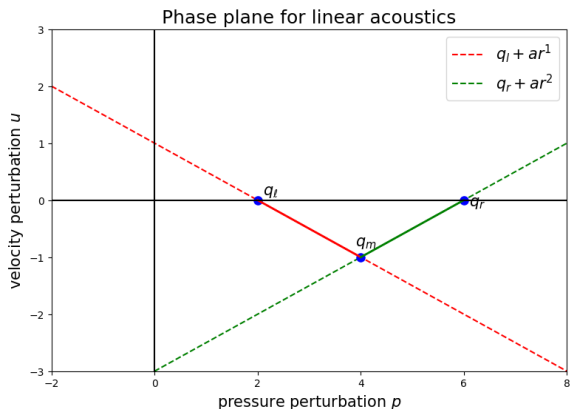


“Shock tube” solution with $u_\ell = u_r = 0$.

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Riemann solution for a linear system

Linear hyperbolic system: $q_t + Aq_x = 0$ with $A = R\Lambda R^{-1}$.
General Riemann problem data $q_l, q_r \in \mathbb{R}^m$.

Decompose jump in q into eigenvectors:

$$q_r - q_l = \sum_{p=1}^m \alpha^p r^p$$

Note: the vector α of eigen-coefficients is

$$\alpha = R^{-1}(q_r - q_l) = R^{-1}q_r - R^{-1}q_l = w_r - w_l.$$

Riemann solution consists of m waves $\mathcal{W}^p \in \mathbb{R}^m$:

$$\mathcal{W}^p = \alpha^p r^p, \quad \text{propagating with speed } s^p = \lambda^p.$$

Phase space

For a system of m equations, phase space is m -dimensional.

Solving the Riemann problem finds a path from q_ℓ to q_r that generally has m segments, each in the direction of an eigenvector (for a linear system; curves more generally).

If $\lambda^1 \leq \lambda^2 \leq \dots \leq \lambda^m$, then first segment from q_ℓ to $q_\ell + \alpha^1 r^1$, next segment goes to $q_\ell + \alpha^1 r^1 + \alpha^2 r^2$, etc.

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Unique such path provided eigenvectors are linearly independent. $q_\ell + \alpha^1 r^1 + \alpha^2 r^2 + \dots + \alpha^m r^m = q_r$.

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Visualization is most useful when $m = 2$ (phase plane).

But sometimes illuminating to project phase space onto a two-dimensional plane.