

10/15/2009

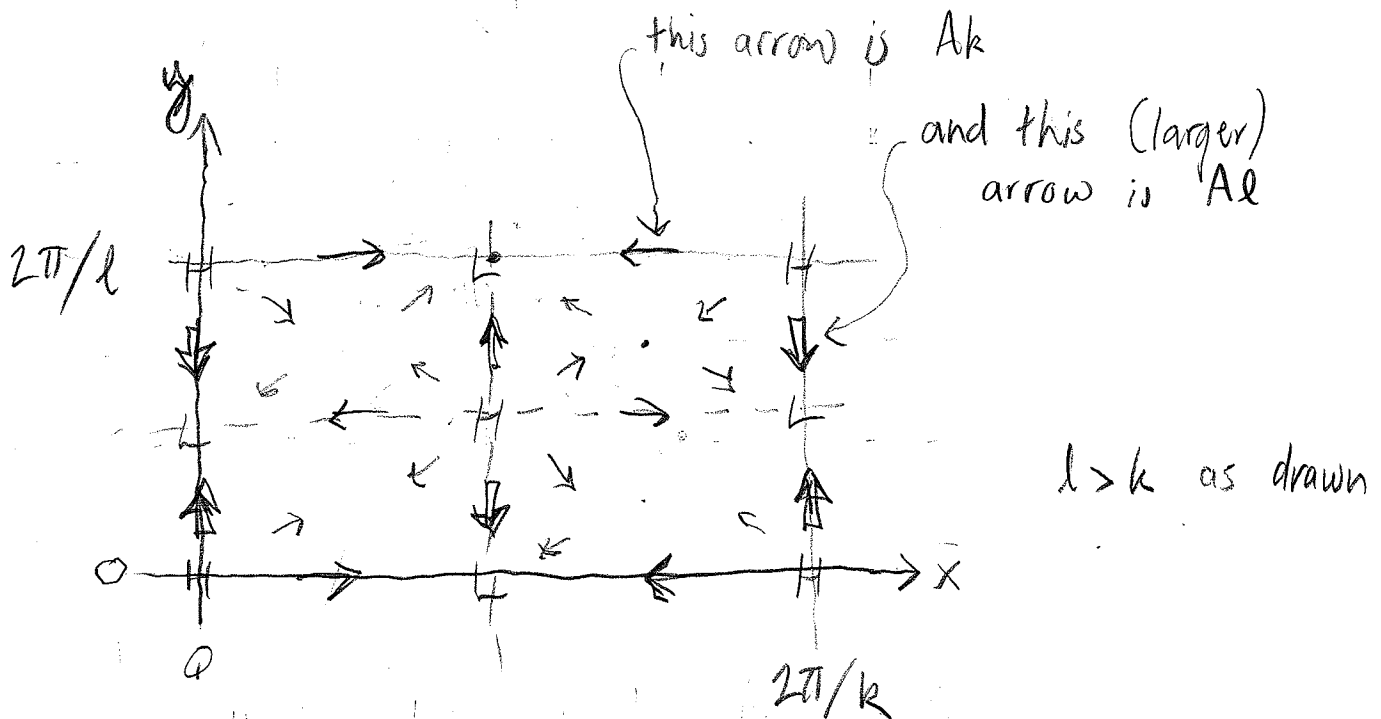
①

FLUIDS 2009

Problem Set #2 Solutions

1.i. since $p = A \cos(kx) \cos(l y)$

$$\Rightarrow -\nabla p = \left(-\frac{\partial p}{\partial x}, -\frac{\partial p}{\partial y} \right) = (A k \sin(kx) \cos(l y), A l \cos(kx) \sin(l y))$$



H & L denote the initial pressure field

1.ii. The divergence is

$$-\nabla \cdot \nabla p = \left(-\frac{\partial}{\partial x} \frac{\partial p}{\partial x} - \frac{\partial}{\partial y} \frac{\partial p}{\partial y} \right) = -(p_{xx} + p_{yy}) = -\nabla^2 p$$

$$= A(k^2 + l^2) \cos kx \cos ly$$

1.iii. Clearly this would not be incompressible if $\underline{u} \propto -\nabla p$

2.i. $\frac{dT}{Dt} = 0$

2.ii. zero

2.iii Note that $\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} = 0$ (*)

Does $\frac{\partial T}{\partial x}$ change with time? No, you can show this by taking $\frac{\partial (*)}{\partial x}$ to form

$$\frac{\partial}{\partial t} \left(\frac{\partial T}{\partial x} \right) + \frac{\partial u}{\partial x} \frac{\partial T}{\partial x} + u \frac{\partial^2 T}{\partial x^2} = 0$$

\swarrow \searrow
 0 0 at $t=0$

$\therefore \frac{\partial T}{\partial x}$ is constant for all time, and $\frac{\partial T}{\partial x} = \frac{T_0}{L}$

so (*) can be written as

$$\frac{\partial T}{\partial t} + \frac{u}{H} \frac{T_0}{L} z = 0 \Rightarrow T_t = - \frac{u T_0}{HL} z$$

(3)

this has solution

$$T = -\frac{\alpha T_0}{HL} z t + \underbrace{f(x, z)}_{T(t=0)}$$

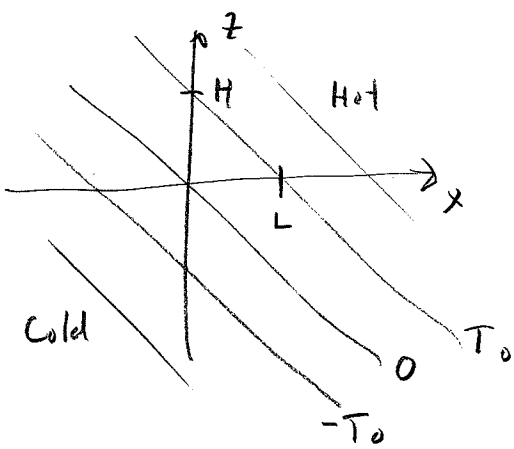
constant of integration

so the full solution is

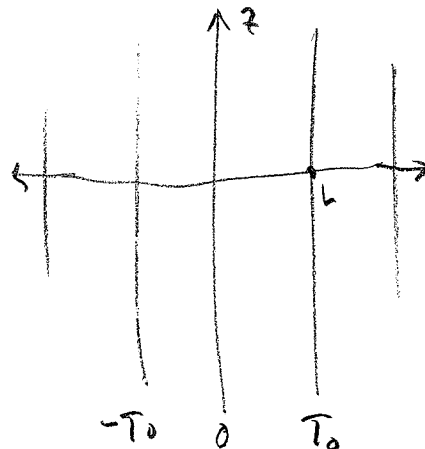
$$T = \frac{T_0}{L} x + \frac{T_0}{H} z \left(1 - \frac{\alpha}{L} t\right)$$

at $z = \frac{H}{2}$:

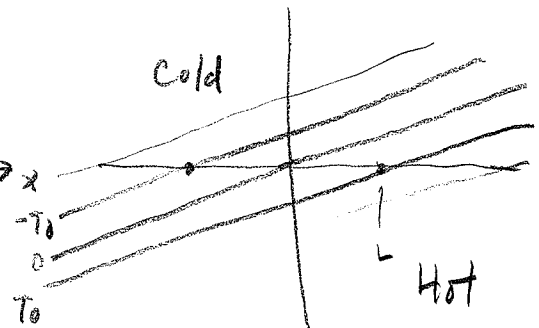
$$\frac{\partial T}{\partial t} = -\frac{T_0 \alpha}{2L}$$



$$t = 0$$



$$t = \frac{L}{\alpha}$$



$$t \gg \frac{L}{\alpha}$$

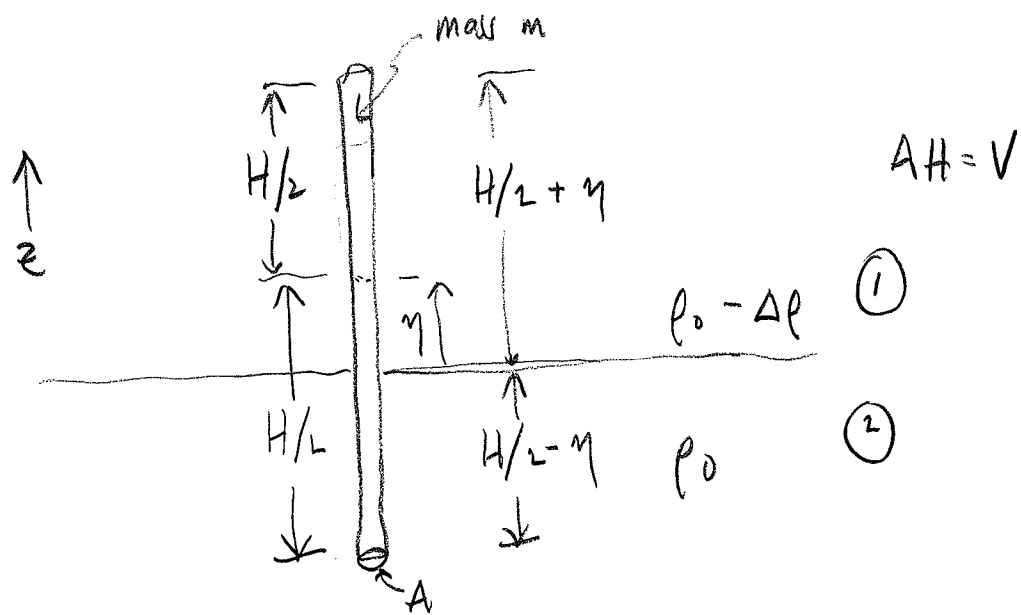
2. iv.

$$\nabla T = \left(\frac{\partial T}{\partial x}, \frac{\partial T}{\partial z} \right) = \left[\frac{T_0}{L}, \frac{T_0}{H} \left(1 - \frac{\alpha}{L} t\right) \right]$$

constant

growing ever more negative in time (after $t = L/\alpha$)

3.i.



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$$\vec{F} = m \vec{a}$$

$$(-) \quad \underbrace{-mg}_{\text{gravity}} + \underbrace{A \left(\frac{H}{2} + \eta \right) (\rho_0 - \Delta \rho) g + A \left(\frac{H}{2} - \eta \right) \rho_0 g}_{\text{Buoyancy} = \int_V \rho \, dV} = m \frac{\partial^2 \eta}{\partial t^2}$$

ρ of the displaced fluid

for steady state $\eta = \eta_{tt} = 0$

$$\Rightarrow m = AH\rho_0 - \frac{\Delta H}{2} \Delta \rho \approx AH\rho_0 \quad (\text{for } \Delta \rho \ll \rho_0)$$

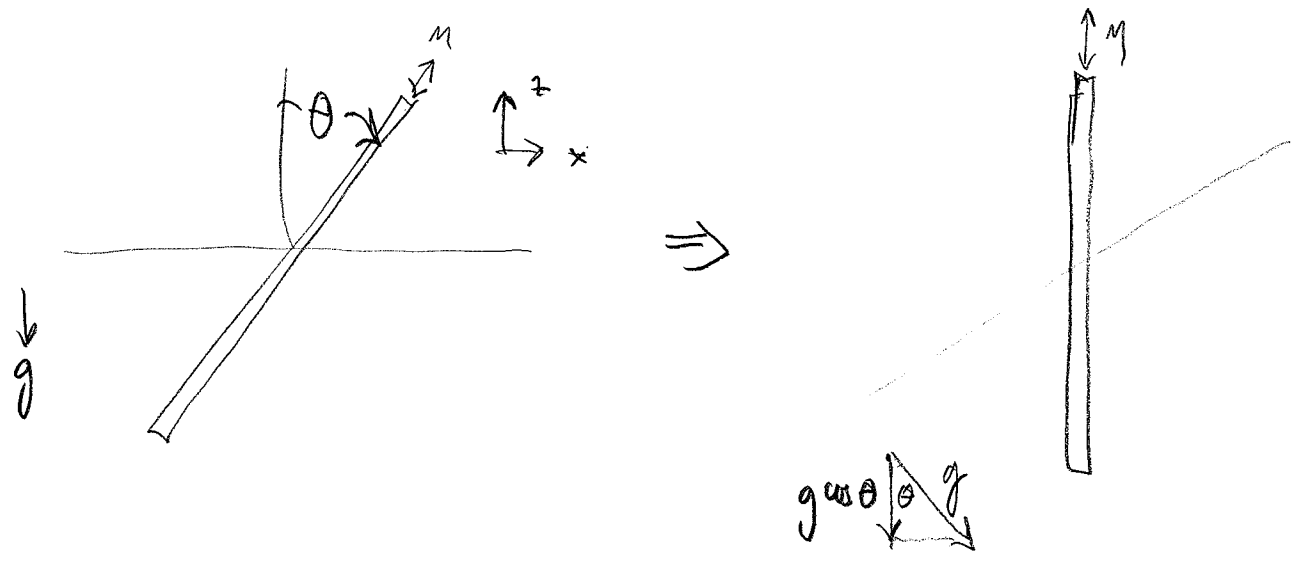
+ then it is easy to simplify (+) to find

$$\eta_{tt} + \frac{g'}{H} \eta = 0$$

where $g' \equiv \frac{g \Delta \rho}{\rho_0}$ is the "reduced gravity"

$$\Rightarrow \eta = \eta_0 \cos \omega t \quad \text{where} \quad \omega = \sqrt{\frac{g'}{H}}$$

3. ii. You can do this most easily by rotating your coordinate system



and the solution is the same as in 3. i.

except
$$\omega = \sqrt{\frac{g' \cos \theta}{H}}$$

3. iii In layer (1)
$$m \eta_{tt} = (\rho_0 - \Delta \rho) V g - m g$$

and in (2)
$$m \eta_{tt} = \rho_0 V g - m g$$

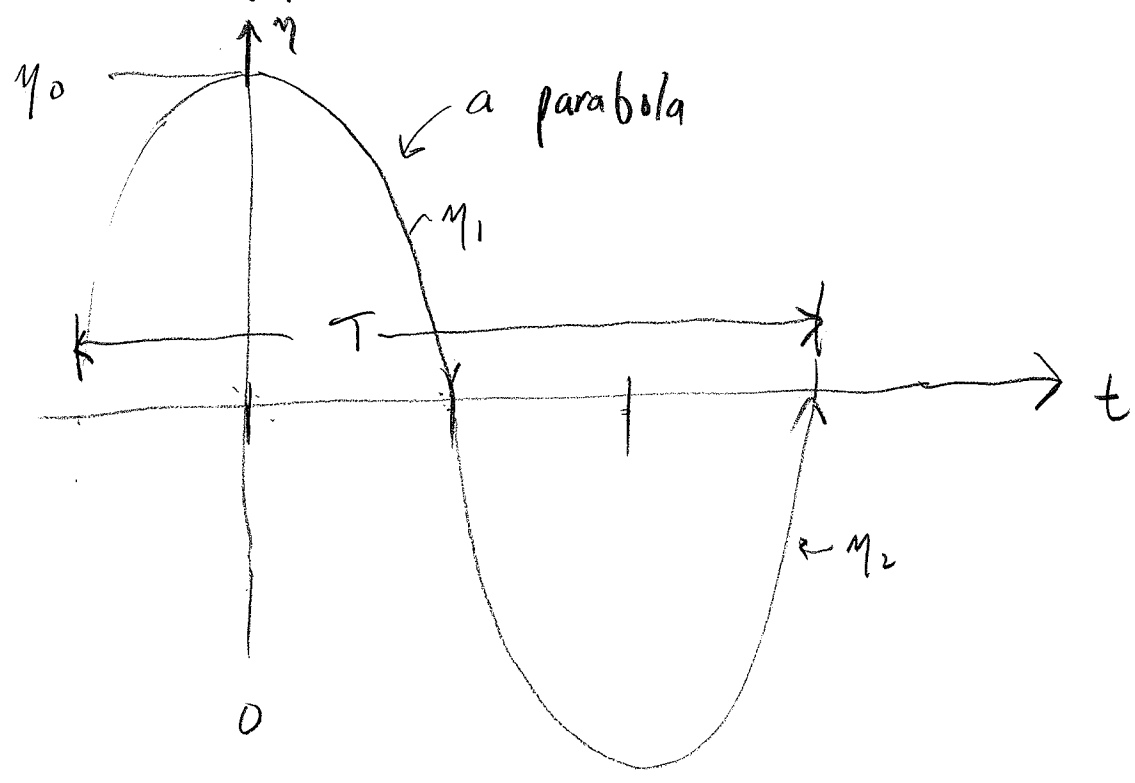
and
$$m = V \rho_0 - V \frac{\Delta \rho}{2}$$

so the equations become

① $\eta_{tt} = -g'/2$

② $\eta_{tt} = g'/2$

which has solution



$$\left. \begin{aligned} \eta_1 &= \eta_0 - \frac{g'}{2} t^2 \\ \eta_2 &= \frac{g'}{2} \left(t - \frac{T}{2}\right)^2 - \eta_0 \end{aligned} \right\} \text{and so on}$$

where $T = 4\sqrt{\frac{2\eta_0}{g'}} = 80 \text{ sec. for } \eta_0 = 2 \text{ m.}$