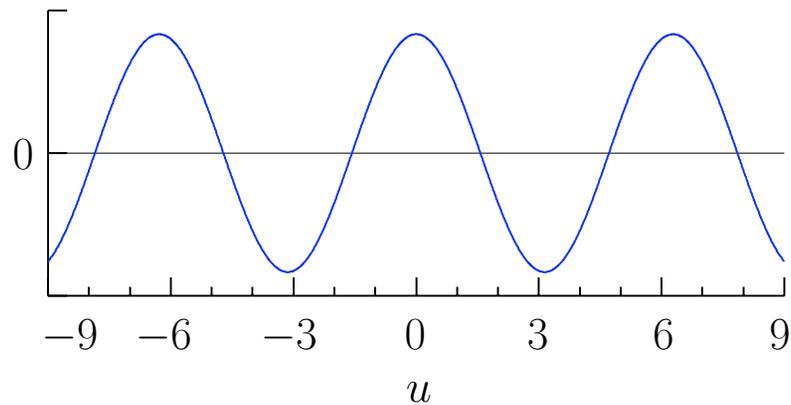


# Introduction to Wavelets: Overview

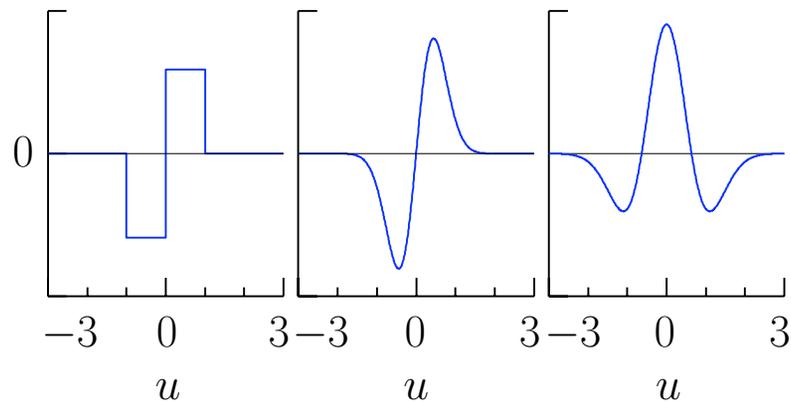
- wavelets are analysis tools for time series and images
- as a subject, wavelets are
  - relatively new (1983 to present)
  - a synthesis of old/new ideas
  - keyword in 50,000+ articles and books since 1989  
(an inundation of material!!!)
- broadly speaking, there have been two waves of wavelets
  - continuous wavelet transform (1983 and on)
  - discrete wavelet transform (1988 and on)
- will introduce subject via CWT & then concentrate on DWT

# What is a Wavelet?

- sines & cosines are ‘big waves’

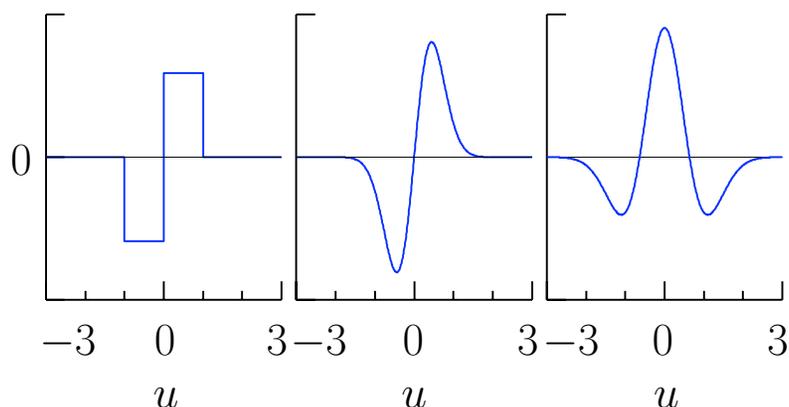


- wavelets are ‘small waves’ (left-hand is Haar wavelet  $\psi^{(H)}(\cdot)$ )



## Technical Definition of a Wavelet: I

- real-valued function  $\psi(\cdot)$  defined over real axis is a wavelet if
  1. integral of  $\psi^2(\cdot)$  is unity:  $\int_{-\infty}^{\infty} \psi^2(u) du = 1$   
(called ‘unit energy’ property, with apologies to physicists)
  2. integral of  $\psi(\cdot)$  is zero:  $\int_{-\infty}^{\infty} \psi(u) du = 0$   
(technically, need an ‘admissibility condition,’ but this is almost equivalent to integration to zero)

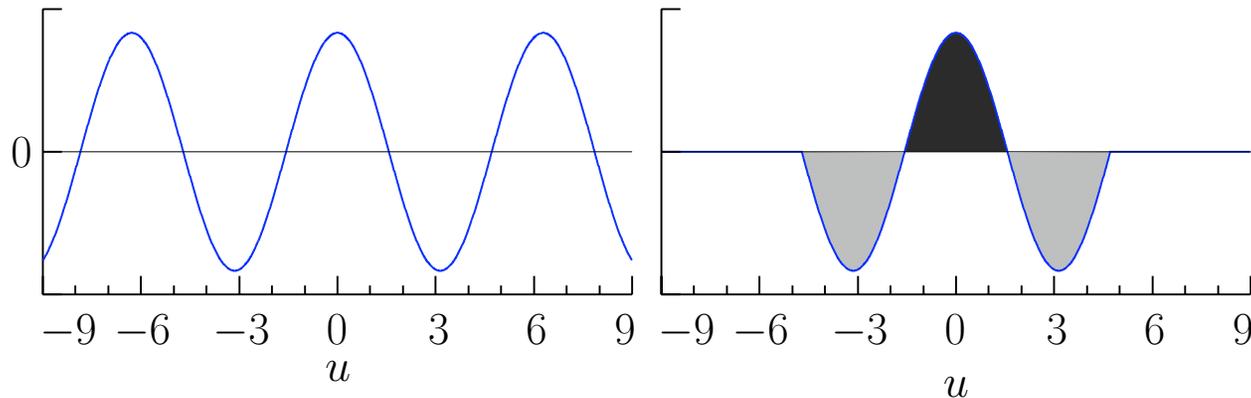


## Technical Definition of a Wavelet: II

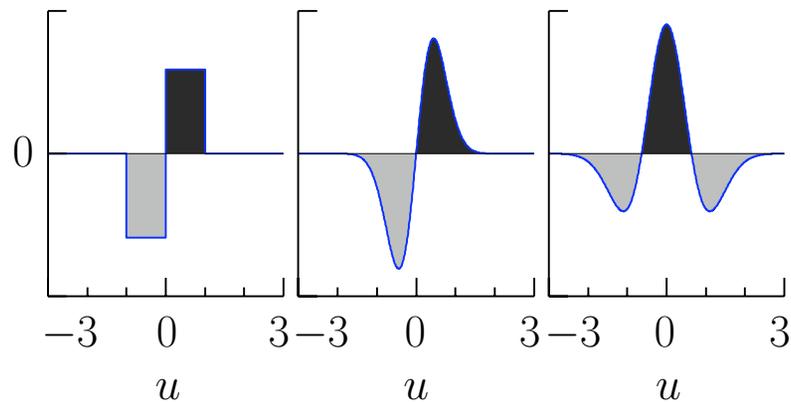
- $\int_{-\infty}^{\infty} \psi^2(u) du = 1$  &  $\int_{-\infty}^{\infty} \psi(u) du = 0$  give a wavelet because:
  - by property 1, for every small  $\epsilon > 0$ , have
$$\int_{-\infty}^{-T} \psi^2(u) du + \int_T^{\infty} \psi^2(u) du < \epsilon$$
for some finite  $T$
  - ‘business’ part of  $\psi(\cdot)$  is over interval  $[-T, T]$
  - width  $2T$  of  $[-T, T]$  might be huge, but will be insignificant compared to  $(-\infty, \infty)$
  - by property 2,  $\psi(\cdot)$  is balanced above/below horizontal axis
- matches intuitive notion of a ‘small’ wave

## Two Non-Wavelets and Three Wavelets

- two failures:  $f(u) = \cos(u)$  & same limited to  $[-3\pi/2, 3\pi/2]$ :



- Haar wavelet  $\psi^{(H)}(\cdot)$  and two of its friends:



## What is Wavelet Analysis?

- wavelets tell us about variations in local averages
- to quantify this description, let  $x(\cdot)$  be a ‘signal’
  - real-valued function of  $t$  defined over real axis
  - will refer to  $t$  as time (but it need not be such)
- consider ‘average value’ of  $x(\cdot)$  over  $[a, b]$ :

$$\frac{1}{b-a} \int_a^b x(t) dt$$

## Approximating Average Value of a Signal

- can approximate integral using Riemann sum
  - break  $[a, b]$  into  $N$  subintervals of equal width  $(b - a)/N$
  - sample  $x(\cdot)$  at midpoint of each subinterval:

$$x_j = x \left( a + \left[ j + \frac{1}{2} \right] \frac{b-a}{N} \right), \quad j = 0, 1, \dots, N - 1$$

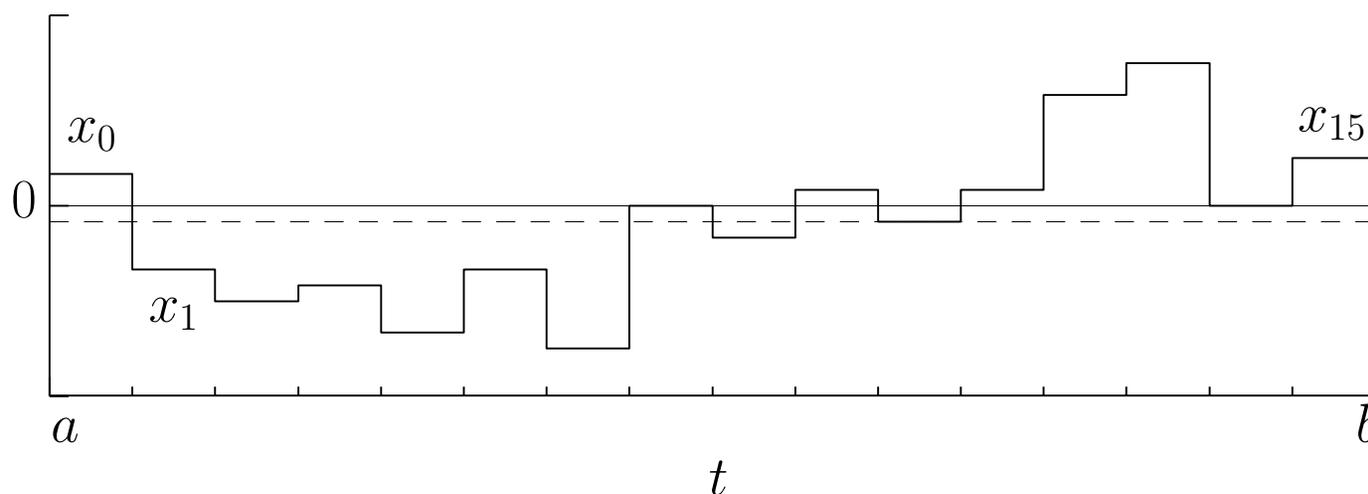
- Riemann sum = sum of  $x_j$ 's  $\times$  width  $(b - a)/N$
- yields approximation to average value of  $x(\cdot)$  over  $[a, b]$ :

$$\frac{1}{b - a} \int_a^b x(t) dt \approx \frac{1}{b - a} \left( \frac{b - a}{N} \sum_{j=0}^{N-1} x_j \right) = \frac{1}{N} \sum_{j=0}^{N-1} x_j$$

- average value of  $x(\cdot) \approx$  sample mean of sampled values

## Example of Average Value of a Signal

- let  $x(\cdot)$  be step function taking on values  $x_0, x_1, \dots, x_{15}$  over 16 equal subintervals of  $[a, b]$ :



- here we have

$$\frac{1}{b-a} \int_a^b x(t) dt = \frac{1}{16} \sum_{j=0}^{15} x_j = \text{height of dashed line}$$

## Average Values at Different Scales and Times

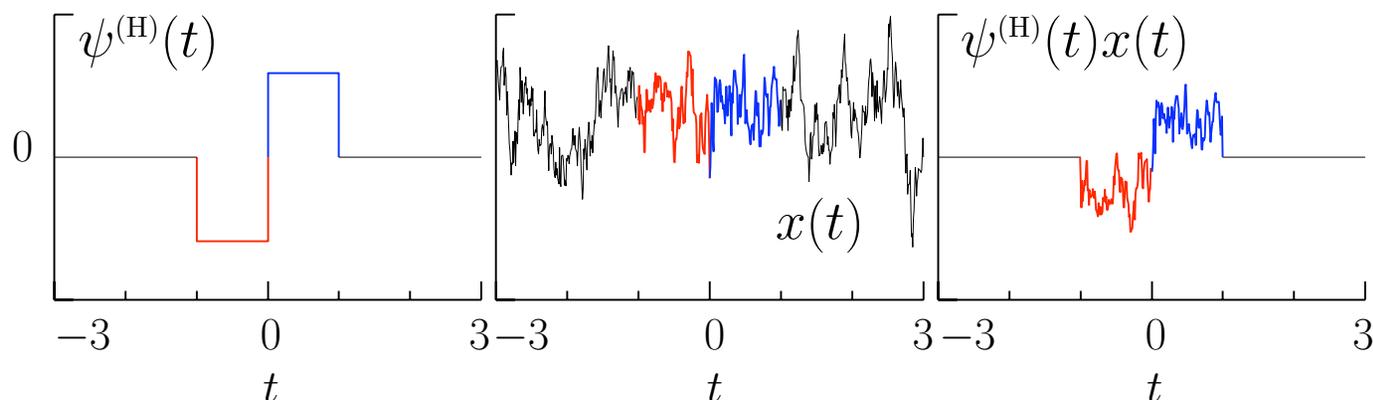
- define the following function of  $\lambda$  and  $t$

$$A(\lambda, t) \equiv \frac{1}{\lambda} \int_{t-\frac{\lambda}{2}}^{t+\frac{\lambda}{2}} x(u) du$$

- $\lambda$  is width of interval – referred to as ‘scale’
- $t$  is midpoint of interval
- $A(\lambda, t)$  is average value of  $x(\cdot)$  over scale  $\lambda$  centered at  $t$
- average values of signals have wide-spread interest
  - one second average temperatures over forest
  - ten minute rainfall rate during severe storm
  - yearly average temperatures over central England

## Defining a Wavelet Coefficient $W$

- multiply Haar wavelet & time series  $x(\cdot)$  together:



- integrate resulting function to get ‘wavelet coefficient’  $W(1, 0)$ :

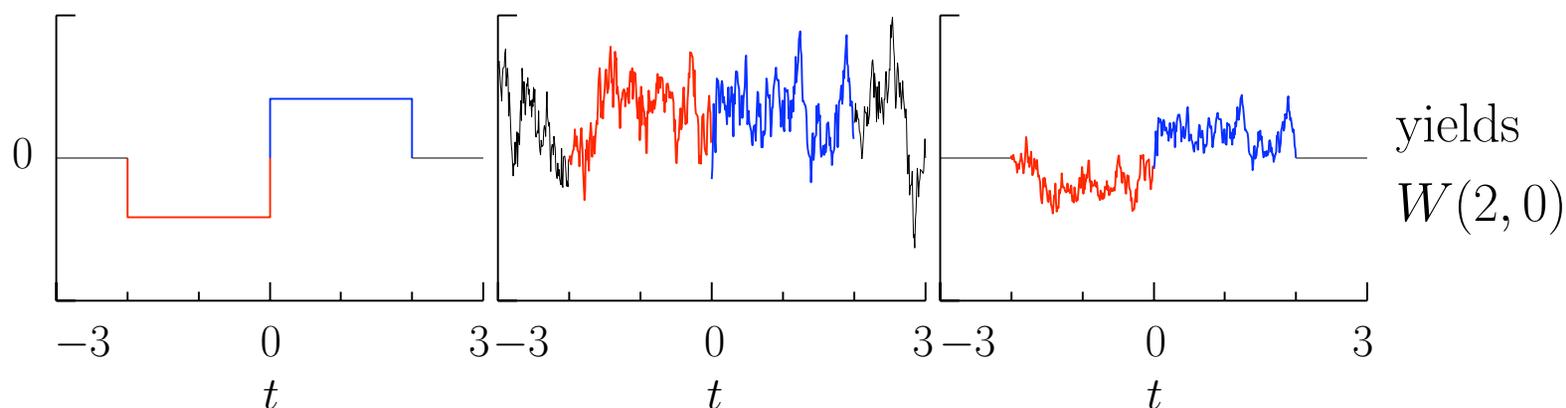
$$\int_{-\infty}^{\infty} \psi^{(H)}(t)x(t) dt = W(1, 0)$$

- to see what  $W(1, 0)$  is telling us about  $x(\cdot)$ , note that

$$W(1, 0) \propto \frac{1}{1} \int_0^1 x(t) dt - \frac{1}{1} \int_{-1}^0 x(t) dt = A(1, \frac{1}{2}) - A(1, -\frac{1}{2})$$

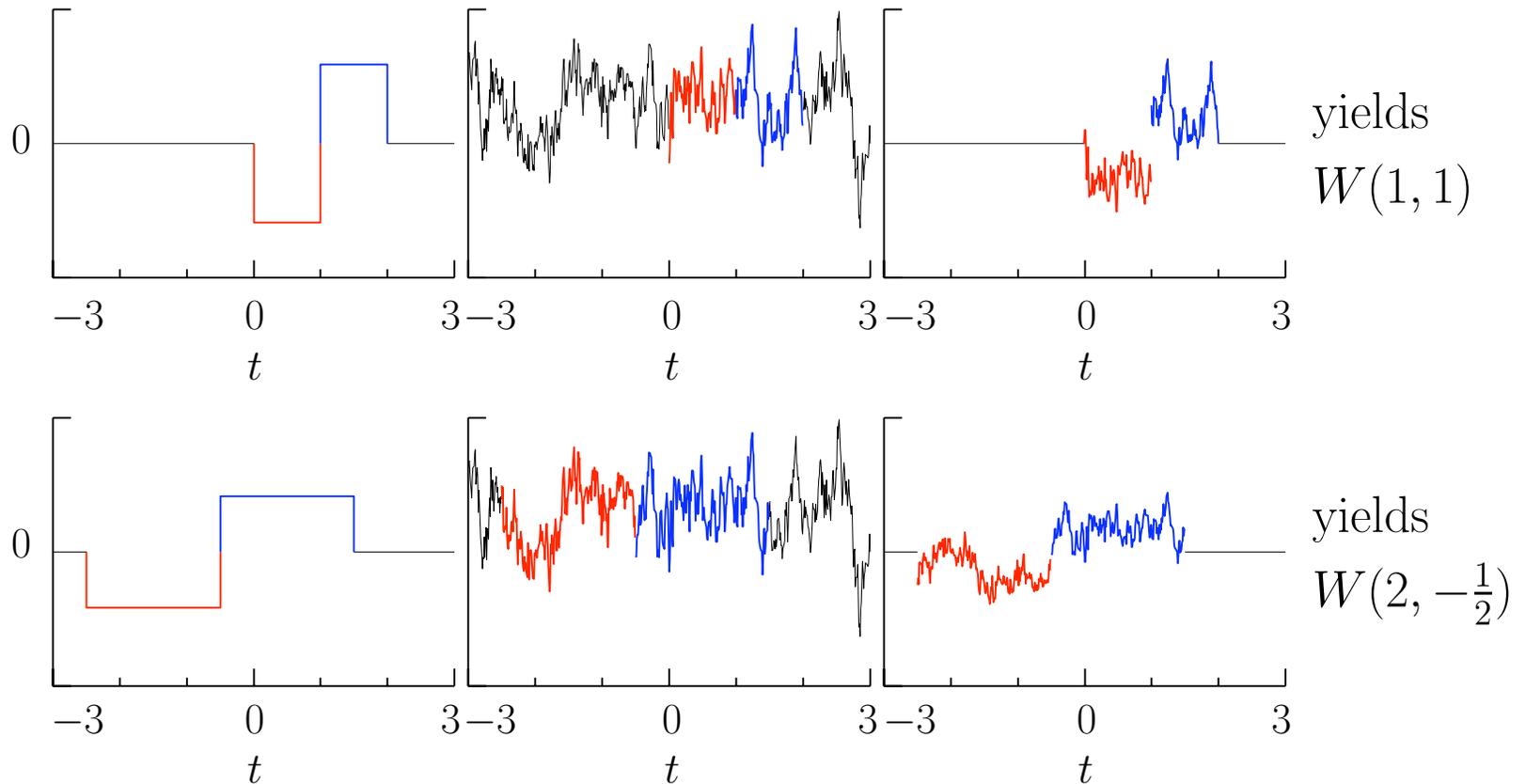
## Defining Wavelet Coefficients for Other Scales

- $W(1, 0)$  proportional to difference between averages of  $x(\cdot)$  over  $[-1, 0]$  &  $[0, 1]$ , i.e., two unit scale averages before/after  $t = 0$ 
  - ‘1’ in  $W(1, 0)$  denotes scale 1 (width of each interval)
  - ‘0’ in  $W(1, 0)$  denotes time 0 (center of combined intervals)
- stretch or shrink wavelet to define  $W(\tau, 0)$  for other scales  $\tau$ :



# Defining Wavelet Coefficients for Other Locations

- relocate to define  $W(\tau, t)$  for other times  $t$ :



## Haar Continuous Wavelet Transform (CWT)

- for all  $\tau > 0$  and all  $-\infty < t < \infty$ , can write

$$W(\tau, t) = \frac{1}{\sqrt{\tau}} \int_{-\infty}^{\infty} x(u) \psi^{(\text{H})} \left( \frac{u-t}{\tau} \right) du$$

- $\frac{u-t}{\tau}$  does the stretching/shrinking and relocating
- $\frac{1}{\sqrt{\tau}}$  needed so  $\psi_{\tau,t}^{(\text{H})}(u) \equiv \frac{1}{\sqrt{\tau}} \psi^{(\text{H})} \left( \frac{u-t}{\tau} \right)$  has unit energy
- since it also integrates to zero,  $\psi_{\tau,t}^{(\text{H})}(\cdot)$  is a wavelet
- $W(\tau, t)$  over all  $\tau > 0$  and all  $t$  is Haar CWT for  $x(\cdot)$
- analyzes/breaks up/decomposes  $x(\cdot)$  into components
  - associated with a scale and a time
  - physically related to a difference of averages

## Other Continuous Wavelet Transforms: I

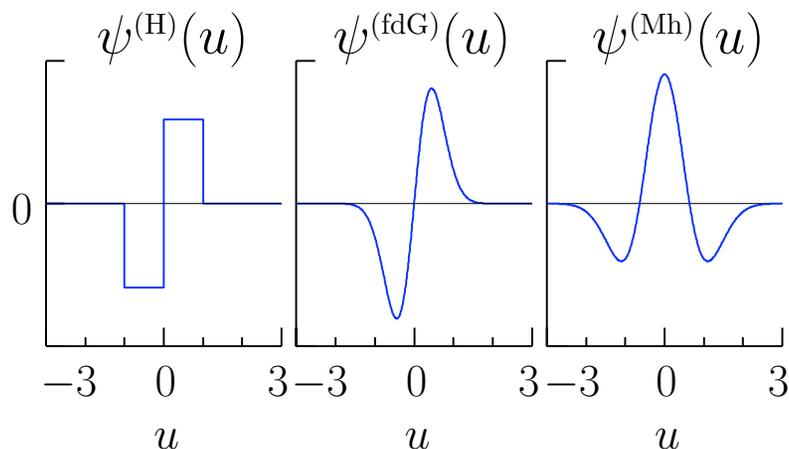
- can do the same for wavelets other than the Haar
- start with basic wavelet  $\psi(\cdot)$
- use  $\psi_{\tau,t}(u) = \frac{1}{\sqrt{\tau}}\psi\left(\frac{u-t}{\tau}\right)$  to stretch/shrink & relocate
- define CWT via

$$W(\tau, t) = \int_{-\infty}^{\infty} x(u)\psi_{\tau,t}(u) du = \frac{1}{\sqrt{\tau}} \int_{-\infty}^{\infty} x(u)\psi\left(\frac{u-t}{\tau}\right) du$$

- analyzes/breaks up/decomposes  $x(\cdot)$  into components
  - associated with a scale and a time
  - physically related to a difference of *weighted* averages

## Other Continuous Wavelet Transforms: II

- consider two friends of Haar wavelet



- $\psi^{(fdG)}(\cdot)$  proportional to 1st derivative of Gaussian PDF
- ‘Mexican hat’ wavelet  $\psi^{(Mh)}(\cdot)$  proportional to 2nd derivative
- $\psi^{(fdG)}(\cdot)$  looks at difference of adjacent weighted averages
- $\psi^{(Mh)}(\cdot)$  looks at difference between weighted average and sum of weighted averages occurring before & after

## First Scary-Looking Equation

- CWT equivalent to  $x(\cdot)$  because we can write

$$x(t) = \int_0^\infty \left[ \frac{1}{C\tau^2} \int_{-\infty}^\infty W(\tau, u) \frac{1}{\sqrt{\tau}} \psi \left( \frac{t-u}{\tau} \right) du \right] d\tau,$$

where  $C$  is a constant depending on specific wavelet  $\psi(\cdot)$

- can synthesize (put back together)  $x(\cdot)$  given its CWT; i.e., nothing is lost in reexpressing signal  $x(\cdot)$  via its CWT
- regard stuff in brackets as defining ‘scale  $\tau$ ’ signal at time  $t$
- says we can reexpress  $x(\cdot)$  as integral (sum) of new signals, each associated with a particular scale
- similar additive decompositions will be one central theme

## Second Scary-Looking Equation

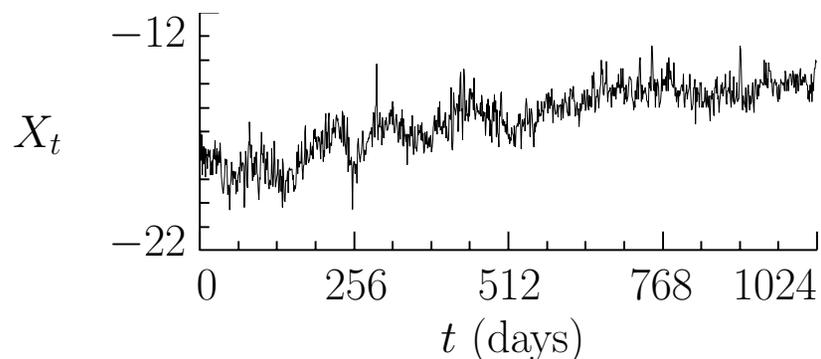
- energy in  $x(\cdot)$  is reexpressed in CWT because

$$\text{energy} = \int_{-\infty}^{\infty} x^2(t) dt = \int_0^{\infty} \left[ \frac{1}{C\tau^2} \int_{-\infty}^{\infty} W^2(\tau, t) dt \right] d\tau$$

- can regard  $x^2(t)$  versus  $t$  as breaking up the energy across time (i.e., an ‘energy density’ function)
- regard stuff in brackets as breaking up the energy across scales
- says we can reexpress energy as integral (sum) of components, each associated with a particular scale
- function defined by  $W^2(\tau, t)/C\tau^2$  is an energy density across both time and scale
- similar energy decompositions will be a second central theme

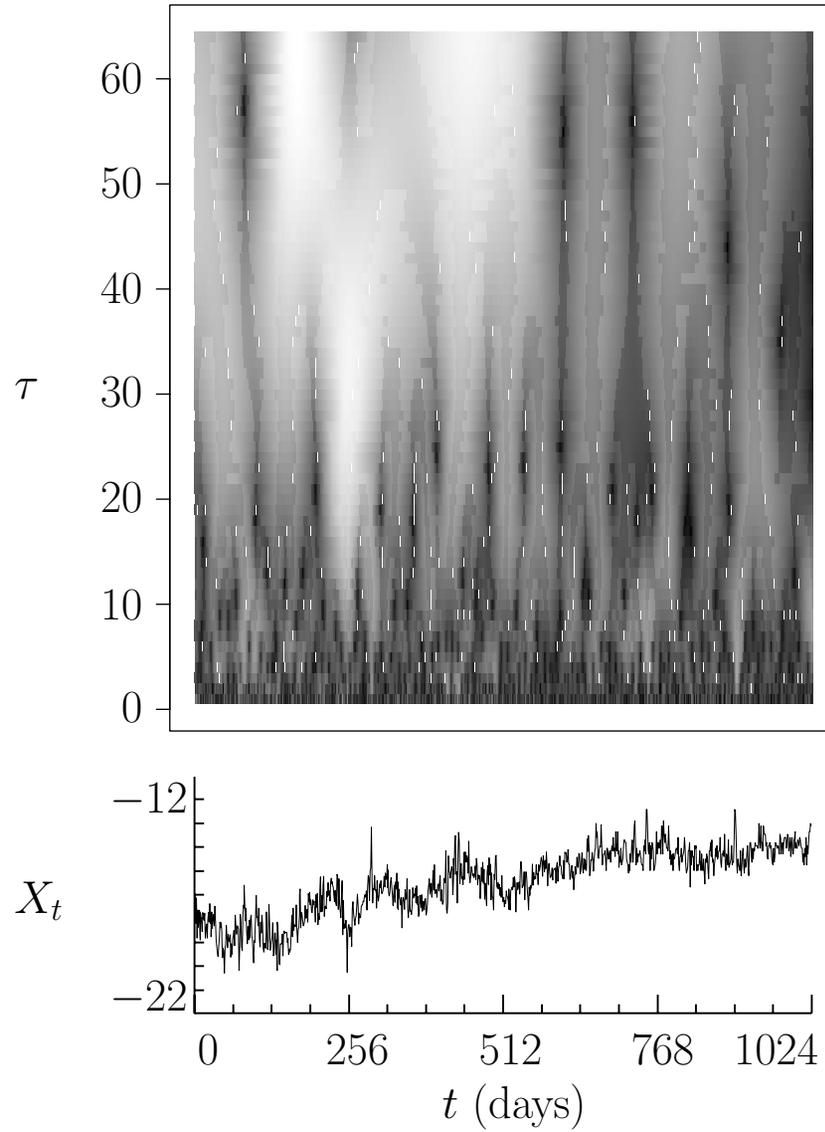
## Example: Atomic Clock Data

- example: average daily frequency variations in clock 571

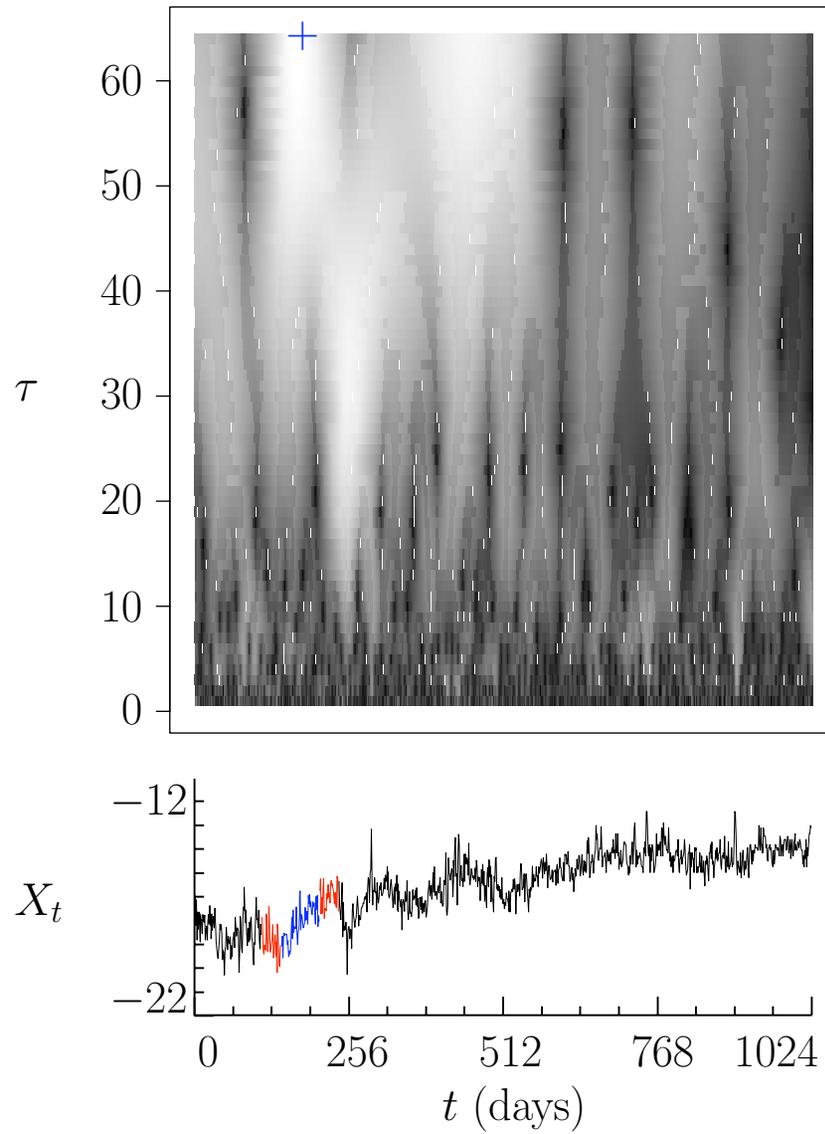


- $t$  is measured in days (one measurement per day)
- plot shows  $X_t$  versus integer  $t$
- $X_t = 0$  would mean that clock 571 could keep time perfectly
- $X_t < 0$  implies that clock is losing time systematically
- can easily adjust clock if  $X_t$  were constant
- inherent quality of clock related to changes in averages of  $X_t$

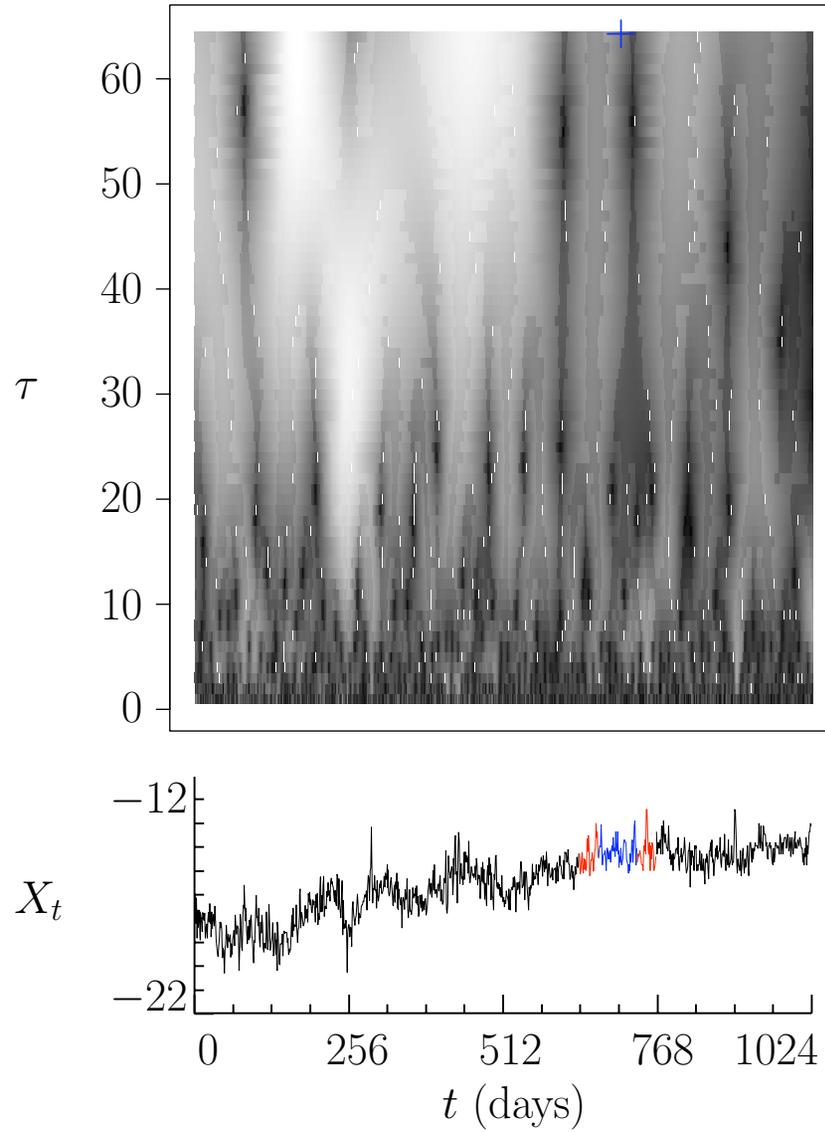
# Mexican Hat CWT of Clock Data: I



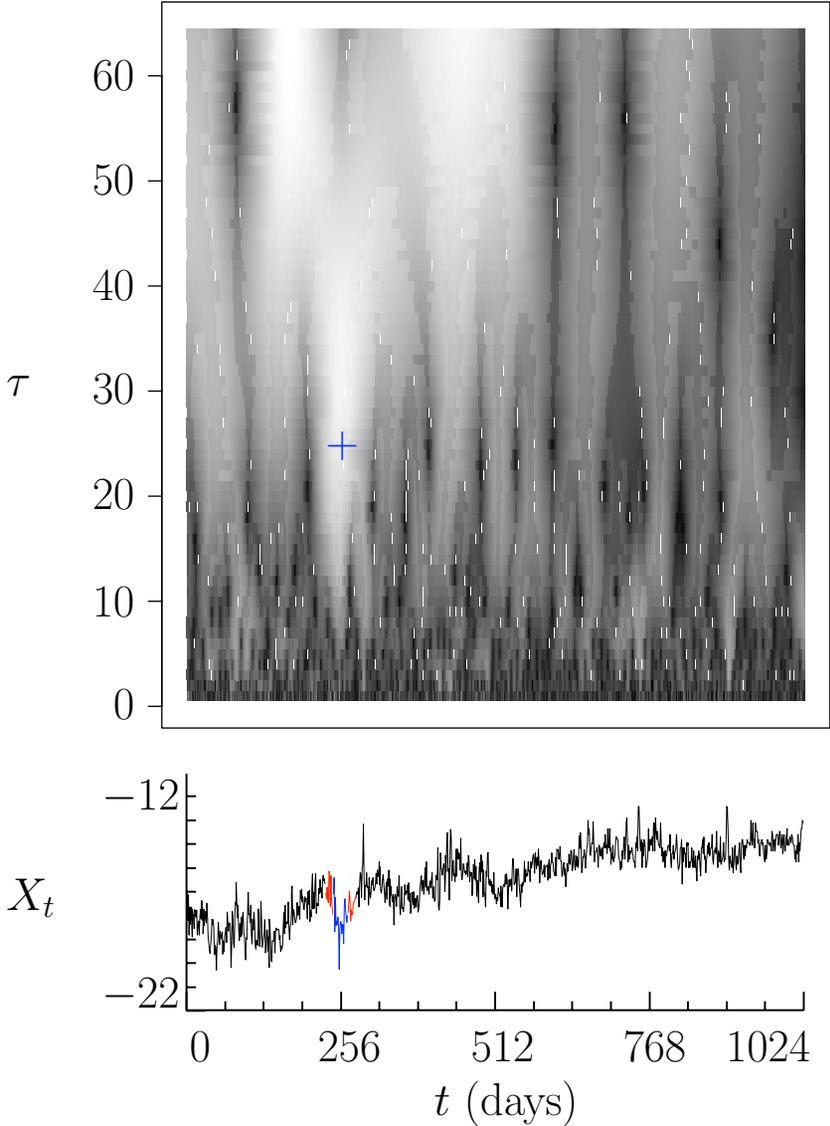
## Mexican Hat CWT of Clock Data: II



# Mexican Hat CWT of Clock Data: III

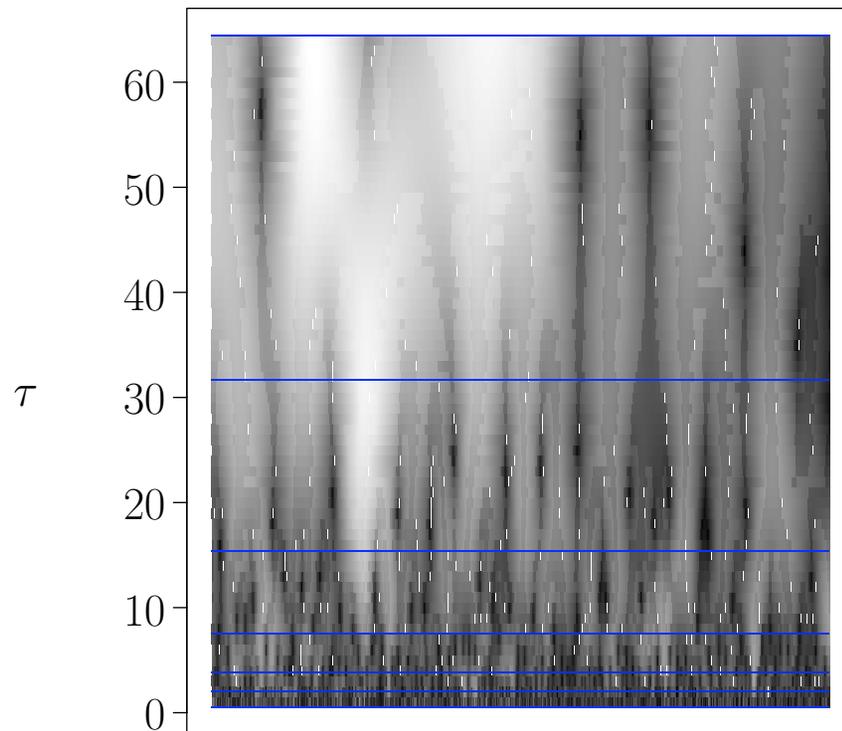


# Mexican Hat CWT of Clock Data: IV



## Beyond the CWT: the DWT

- can often get by with subsamples of  $W(\tau, t)$
- leads to notion of discrete wavelet transform (DWT)  
(can regard as discretized ‘slices’ through CWT)



## Rationale for the DWT

- DWT has appeal in its own right
  - most time series are sampled as discrete values (can be tricky to implement CWT)
  - can formulate as orthonormal transform (makes meaningful statistical analysis possible)
  - tends to decorrelate certain time series
  - standardization to dyadic scales often adequate
  - generalizes to notion of wavelet packets
  - can be faster than the fast Fourier transform
- will concentrate primarily on DWT for remainder of course

## Addendum on First Scary-Looking Equation: I

- can synthesize signal  $x(\cdot)$  from its CWT  $W(\cdot, \cdot)$ :

$$x(t) = \int_0^\infty \left[ \frac{1}{C\tau^2} \int_{-\infty}^\infty W(\tau, u) \frac{1}{\sqrt{\tau}} \psi\left(\frac{t-u}{\tau}\right) du \right] d\tau, \quad (*)$$

where  $C$  is a constant depending on specific wavelet  $\psi(\cdot)$

- Q: what is the constant  $C$  all about?
- as mentioned on overhead I-3, for a function  $\psi(\cdot)$  to be a wavelet, it must satisfy a so-called ‘admissibility condition’
- to state admissibility condition, let  $\Psi(\cdot)$  denote Fourier transform of  $\psi(\cdot)$  (assumed to be a square-integrable function):

$$\Psi(f) = \int_{-\infty}^\infty \psi(u) e^{-i2\pi fu} du$$

## Addendum on First Scary-Looking Equation: II

- admissibility condition says that

$$C \equiv \int_0^\infty \frac{|\Psi(f)|^2}{f} df \text{ must be such that } 0 < C < \infty$$

(note: above implies that  $\psi(\cdot)$  must integrate to zero)

- $C$  above is same  $C$  appearing in (\*)
- as to *why*  $C$  appears, need to work through proof of (\*), which is not trivial
  - see Mallat, 1998, §4.3 for a clear proof
  - proof in the wavelet literature due to Grossman and Morlet, 1984, who discuss why admissibility condition is needed
  - Grossman and Morlet's result actually appeared earlier in 1964 paper by Calderón