

Critical Temperature for Deconfinement in the $(2+1)$ -d Georgi-Glashow Model

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Outline

- ◆ Definition of the theory
- ◆ Motivation
- ◆ The simple case: $SU(2)$
- ◆ What happens for $SU(N)$?
- ◆ Summary

The theory: Pure gauge SU(N) + adjoint scalars

◆ Lagrangian

$$\mathcal{L} = \frac{1}{2} (D_\mu \phi^a)^2 + \lambda (\phi^a \phi^a - v^2)^2 + \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a$$

Adjoint scalars (N^2-1)

$$\phi^a T^a \in ASU(N)$$

Covariant derivative

$$D_\mu^a = \partial_\mu + ad(A_\mu)^a$$

Field strength tensor

$$F_{\mu\nu} = F_{\mu\nu}^a T^a = [D_\mu, D_\nu] = \partial_\mu A_\nu - \partial_\nu A_\mu + [A_\mu, A_\nu]$$

The theory: Pure gauge SU(N) + adjoint scalars

◆ Lagrangian

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◆ Perturbatively...

	SU(2)	SU(N)	
massive bosons	$W^\pm = A^1 \pm iA^2$	$N^2 - N$	$m_W = gv$
massless photon	A^3	$N - 1$	(for maximal sym. breaking)
In the Unitary gauge:	$\phi^a = \delta^{a3} (v + \sigma(x))$		

◆ Weak coupling...

$$m_W \sim m_H \gg g^2$$

Motivation: why is this interesting?

- These theories have topologically non-trivial classical solutions of finite action (a.k.a. solitons/instantons).
- In $(2+1)$ -d, instantons provide a mechanism for confinement of charge, and this happens at weak coupling.
 - One can prove, e.g....
 - ▲ Confinement
 - ▲ T_c for Deconfinement

... analytically !

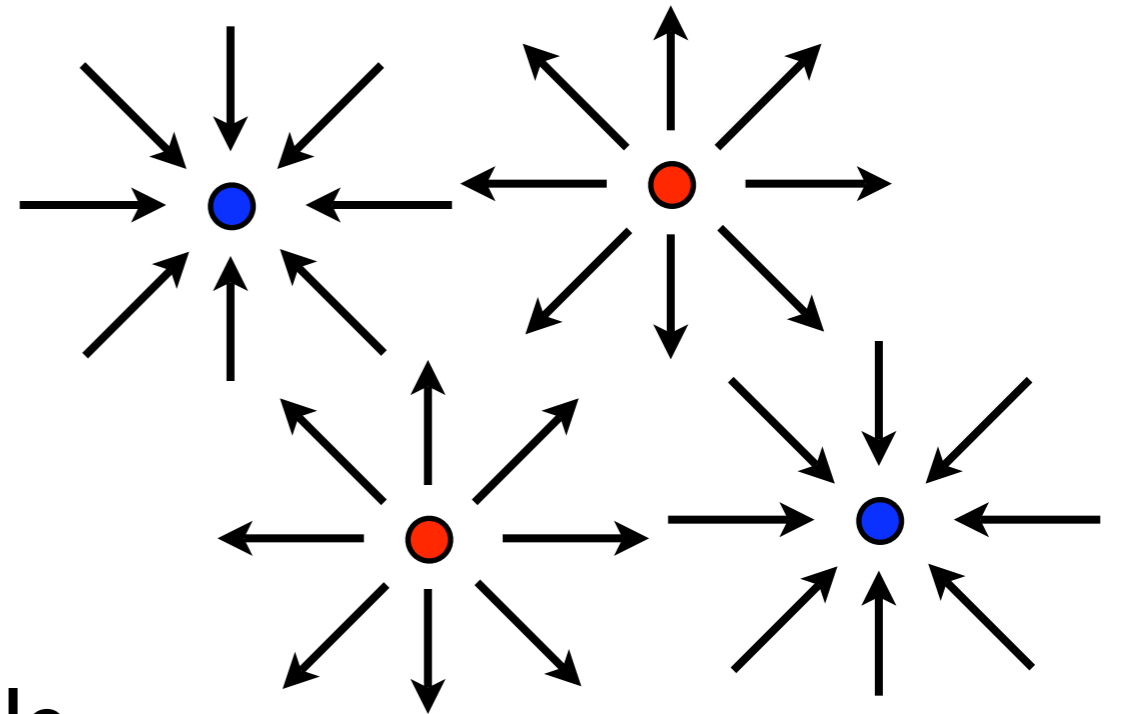
Confinement in SU(2)

I. Polyakov's proof:

'Instantons generate a non-perturbative mass for the photon'

Classical solutions (photon part)

$$F_{\mu\nu}^3 = \frac{e}{4\pi} \epsilon_{\mu\nu\lambda} \frac{x_\lambda}{x^2} \quad e = \frac{4\pi}{g}$$



Write the partition function in a saddle-point approximation as a sum over multi-instanton configurations.

→ Coulomb gas (of instantons) representation.

Confinement in SU(2)

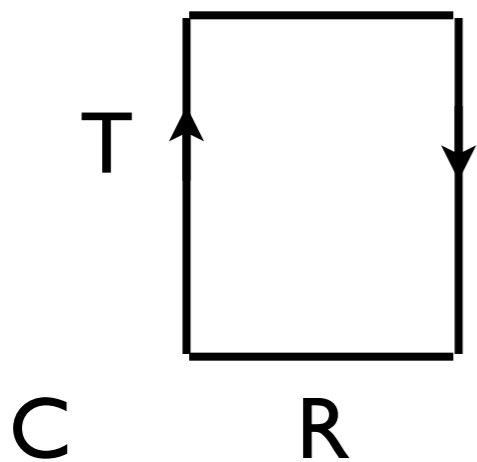
→ Map Coulomb gas onto Sine-Gordon model.

$$\mathcal{L}_{eff} \propto (\partial_\mu \varphi)^2 - M^2 \cos \varphi \qquad S_0 = \frac{4\pi v}{g} \epsilon \left(\frac{m_H}{m_W} \right)$$

$$M^2 \propto \frac{m_w^{7/2}}{g^3} e^{-S_0}$$

$$m_\gamma \sim e^{-S_0/2}$$

→ Compute Wilson loop



$$F(C) = e^{-E(R)T} = e^{-\gamma A}$$

$$\gamma \propto g^2 M$$

$$E(R) = \gamma R$$

Confinement !

Deconfinement in SU(2)

2. Deconfinement transition

(Dunne et al.; Kovchegov & Son).

→ Again map onto Sine-Gordon model...
(plus W bosons, in dimensional reduction)

$$H = \int dx \left[\frac{1}{2}(\partial_x \Phi)^2 + \frac{1}{2}(\partial_x \Theta)^2 + 2\zeta_0 \cos \beta \Phi + 2\tilde{\zeta}_0 \cos \tilde{\beta} \Theta \right]$$

$$\zeta_0 \sim T^2 e^{-S_0},$$

$$\tilde{\zeta}_0 \sim T^2 e^{-m_W/T}$$

$$\beta = e\sqrt{T}, \quad \tilde{\beta} = \frac{g}{\sqrt{T}}$$

$$S_0 = \frac{4\pi v}{g} \epsilon \left(\frac{m_H}{m_W} \right)$$

$$e = \frac{4\pi}{g}$$

Deconfinement in SU(2)

→ Study RG flow...

$$T_c = \frac{g^2}{4\pi} \frac{\epsilon + 2}{2\epsilon + 1}$$

$$S_0 = \frac{4\pi v}{g} \epsilon \left(\frac{m_H}{m_W} \right)$$

→ Universality class

~~BKT~~ \mathbb{Z}_2

Critical exponents are
as in the 2D Ising model

Numerical results for SU(2)

$$S = \beta \sum_{x, \mu > \nu} \left(1 - \frac{1}{2} \text{Tr} U_{\mu\nu}(x) \right) + 2 \sum_x \text{Tr}(\Phi(x)\Phi(x))$$

$$- 2\kappa \sum_{x, \mu} \text{Tr}(\Phi(x)U_\mu(x)\Phi(x + \hat{\mu}a)U_\mu^\dagger(x))$$

$$+ \lambda \sum_x (2\text{Tr}(\Phi(x)\Phi(x)) - 1)^2$$

$$\overline{L}_F = \frac{1}{N_s^2} \sum_{\vec{x}} \left(\text{Tr}_F \left(\prod_{t=1}^{N_\tau} U_\tau(\vec{x}, t) \right) \right)$$

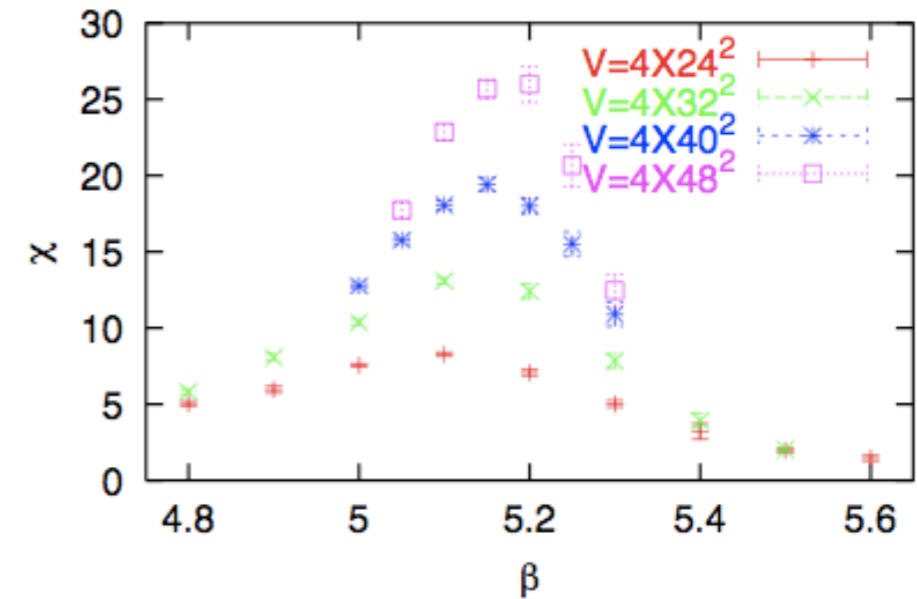
$$\chi_{L_F} = N_s^2 \left(\langle \overline{L}_F^2 \rangle - \langle \overline{L}_F \rangle^2 \right)$$

Z_2 Universality class
confirmed by finite size scaling

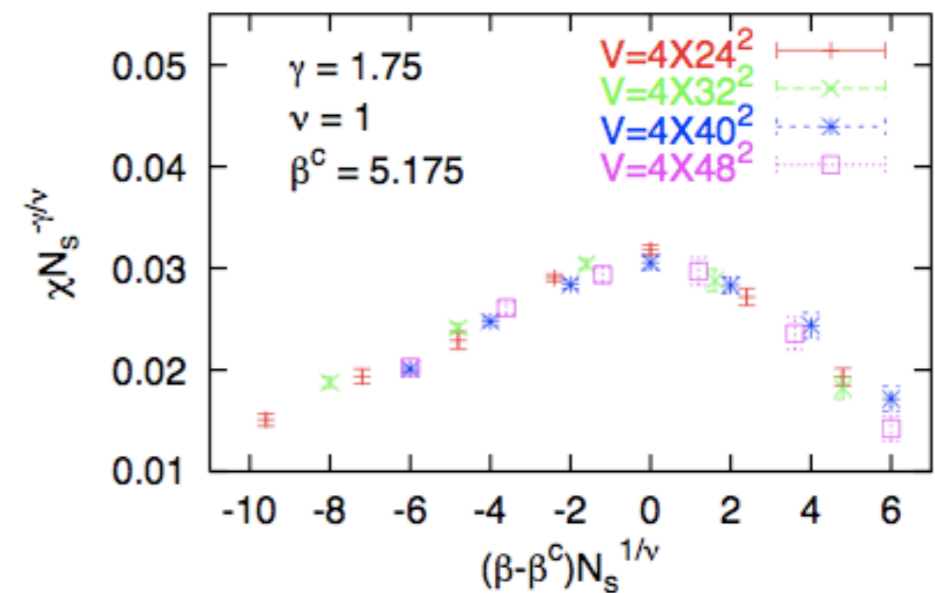


Barresi et al.

$\beta_c = 5.175, x = 0.326696, y = 0.19518$



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What happens for $SU(N)$?

- Put the theory on the lattice.
- Study Polyakov loops and their correlations
- Do finite size scaling analysis

Critical exponents?

- Quantum phase transitions?

Dynamical exponents?

Summary

- ▲ In the $(2+1)$ -d $SU(N)$ Georgi-Glashow model confinement of charge happens at weak coupling and can be proven analytically ('t Hooft & Polyakov).
- ▲ The nature of the finite temperature deconfinement transition is believed to be understood for all N (Svetitsky & Yaffe; Kogan et al.; Lecheminant). Needs numerical check.
- ▲ The critical temperature has only been computed analytically in the simplest case of $SU(2)$ (Son & Kovchegov).
- ▲ The critical temperature for $SU(N)$ remains unknown and is hard to compute analytically.

To be continued...

Thanks!