

Numerical Methods (See, e.g., *Mathematical Methods of Physics*, by Mathews & Walker, *Computational Physics*, by Koonin, and *Numerical Recipes* by Press)

Imagine, as is often the case, that we know that the derivative of a function is given in terms of another function(al) of the original function and the free variable, e.g.,

$$\dot{y}(t) = f(t, y). \quad (\text{NM.1})$$

We also know that $y(t_0) = y_0$ and we want to use this information to (numerically) find y at other values of t , i.e., numerically solve the above equation. Previously we have considered cases where there is no t dependence on the right-hand-side and proceeded by separating variables and writing

$$dt = \frac{dy}{f(y)} \Rightarrow t - t_0 = \int_{y_0}^y \frac{dy'}{f(y')}. \quad (\text{NM.2})$$

Even in this case we may not be able to perform the integral analytically and want to use numerical methods. A common method for numerically estimating an integral is to use Simpson's rule. Suppose we want to evaluate

$$I = \int_a^b f(x) dx, \quad (\text{NM.3})$$

with $f(x)$ a known (but presumably complicated) function. Simpson says that we should divide the interval $\Delta x = b - a$ into n ($n = \text{even}$) smaller equal intervals, $h = \Delta x/n$ and calculate

$$I \approx \frac{h}{3} \left[f(a) + 4f(a+h) + 2f(a+2h) + 4f(a+3h) + 2f(a+4h) \right. \\ \left. + \dots + 2f(b-2h) + 4f(b-h) + f(b) \right]. \quad (\text{NM.4})$$

We could apply this method above to find the time to go from y_0 to y_1 , i.e., $h = (y_1 - y_0)/n$. Using such a method one can test the reliability of the result by

looking at how rapidly the numerical result varies as n is varied. Clearly one expects arbitrarily accurate results as $n \rightarrow \infty$.

Returning to the original problem, Eq. (NM.1), a typical approach would be to use the Runge-Kutta method. By using Taylor series expansion we can show that a good estimate of $y(t_0 + \delta)$ (to order δ^5) is provided by

$$\begin{aligned}
 y(t_0 + \delta) &\simeq y_0 + \frac{1}{6}(k_1 + 2k_2 + 2k_3 + k_4) \left[+O(\delta^5) \right], \\
 k_1 &= \delta f(t_0, y_0), \\
 k_2 &= \delta f\left(t_0 + \frac{\delta}{2}, y_0 + \frac{k_1}{2}\right), \\
 k_3 &= \delta f\left(t_0 + \frac{\delta}{2}, y_0 + \frac{k_2}{2}\right), \\
 k_4 &= \delta f(t_0 + \delta, y_0 + k_3).
 \end{aligned}
 \tag{NM.5}$$

This result is, in fact quite similar to the Simpson result. To see where this comes from we define $f_0 = f(t_0, y_0)$, $f'_0 = \partial f / \partial y|_{(t_0, y_0)}$, $\dot{f}_0 = \partial f / \partial t|_{(t_0, y_0)}$, etc. With this notation we can expand the above expressions (in general to arbitrary order, but here to order δ^4)

$$\begin{aligned}
k_1 &= \delta f_0, \\
k_2 &\approx \delta \left[f_0 + \frac{\delta}{2} \dot{f}_0 + \frac{k_1}{2} f_0' + \frac{\delta^2}{4 \cdot 2} \ddot{f}_0 + 2 \frac{\delta k_1}{4 \cdot 2} \dot{f}_0' + \frac{k_1^2}{4 \cdot 2} f_0'' \right. \\
&\quad \left. + \frac{\delta^3}{8 \cdot 6} \ddot{\dot{f}}_0 + \frac{3}{8 \cdot 6} \delta^2 k_1 \ddot{f}_0' + \frac{3}{8 \cdot 6} \delta k_1^2 \dot{f}_0'' + \frac{k_1^3}{8 \cdot 6} f_0''' \right] + O(\delta^5) \\
&= \delta \left[f_0 + \frac{\delta}{2} \dot{f}_0 + \frac{\delta}{2} f_0 f_0' + \frac{\delta^2}{8} \ddot{f}_0 + \frac{\delta^2}{4} f_0 \dot{f}_0' + \frac{\delta^2}{8} f_0^2 f_0'' \right. \\
&\quad \left. + \frac{\delta^3}{48} \ddot{\dot{f}}_0 + \frac{1}{16} \delta^3 f_0 \ddot{f}_0' + \frac{1}{16} \delta^3 f_0^2 \dot{f}_0'' + \frac{\delta^3}{48} f_0^3 f_0''' \right] + O(\delta^5), \tag{NM.6} \\
&= \delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f_0') + \frac{\delta^3}{8} (\ddot{f}_0 + 2 f_0 \dot{f}_0' + f_0^2 f_0'') \\
&\quad + \frac{\delta^4}{48} (\ddot{\dot{f}}_0 + 3 f_0 \ddot{f}_0' + 3 f_0^2 \dot{f}_0'' + f_0^3 f_0''') + O(\delta^5),
\end{aligned}$$

$$\begin{aligned}
k_3 &\approx \delta \left[f_0 + \frac{\delta}{2} \dot{f}_0 + \frac{k_2}{2} f_0' + \frac{\delta^2}{4 \cdot 2} \ddot{f}_0 + 2 \frac{\delta k_2}{4 \cdot 2} \dot{f}_0' + \frac{k_2^2}{4 \cdot 2} f_0'' \right. \\
&\quad \left. + \frac{\delta^3}{8 \cdot 6} \ddot{f}_0 + \frac{3}{8 \cdot 6} \delta^2 k_2 \ddot{f}_0' + \frac{3}{8 \cdot 6} \delta k_2^2 \dot{f}_0'' + \frac{k_2^3}{8 \cdot 6} f_0''' \right] + O(\delta^5) \\
&= \delta \left[f_0 + \frac{\delta}{2} \dot{f}_0 + \frac{1}{2} \left\{ \delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f_0') + \frac{\delta^3}{8} (\ddot{f}_0 + 2 f_0 \dot{f}_0' + f_0^2 f_0'') \right\} f_0' \right. \\
&\quad + \frac{\delta^2}{8} \ddot{f}_0 + \frac{\delta}{4} \left\{ \delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f_0') \right\} \dot{f}_0' + \frac{1}{8} \left\{ \delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f_0') \right\}^2 f_0'' \\
&\quad \left. + \frac{\delta^3}{48} \ddot{f}_0 + \frac{1}{16} \delta^2 \{ \delta f_0 \} \ddot{f}_0' + \frac{1}{16} \delta \{ \delta f_0 \}^2 \dot{f}_0'' + \frac{\delta^3}{48} f_0^3 f_0''' \right] + O(\delta^5), \\
&= \delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f_0') + \frac{\delta^3}{8} \{ \ddot{f}_0 + 2 f_0 \dot{f}_0' + 2 (\dot{f}_0 + f_0 f_0') f_0' + f_0^2 f_0'' \} \\
&\quad + \frac{\delta^4}{16} \{ \ddot{f}_0 f_0' + 4 f_0 \dot{f}_0' \dot{f}_0' + 3 f_0^2 f_0' f_0'' + 2 \dot{f}_0 \dot{f}_0' + 2 f_0 \dot{f}_0 f_0'' \\
&\quad + \frac{1}{3} \ddot{f}_0 + f_0 \ddot{f}_0' + f_0^2 \dot{f}_0'' + \frac{1}{3} f_0^3 f_0''' \} + O(\delta^5), \tag{NM.7}
\end{aligned}$$

$$\begin{aligned}
k_4 &\approx \delta \left[f_0 + \delta \dot{f}_0 + k_3 f'_0 + \frac{\delta^2}{2} \ddot{f}_0 + 2 \frac{\delta k_3}{2} \dot{f}'_0 + \frac{k_3^2}{2} f''_0 \right. \\
&\quad \left. + \frac{\delta^3}{6} \ddot{f}_0 + \frac{3}{6} \delta^2 k_3 \ddot{f}_0 + \frac{3}{6} \delta k_3^2 \dot{f}''_0 + \frac{k_3^3}{6} f'''_0 \right] + O(\delta^5) \\
&= \delta \left[f_0 + \delta \dot{f}_0 + \left(\delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f'_0) \right. \right. \\
&\quad \left. \left. + \frac{\delta^3}{8} \left\{ \ddot{f}_0 + 2 f_0 \dot{f}'_0 + 2 (\dot{f}_0 + f_0 f'_0) f'_0 + f_0^2 f_0'' \right\} \right) f'_0 \right. \\
&\quad \left. + \frac{\delta^2}{2} \ddot{f}_0 + \delta \left(\delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f'_0) \right) \dot{f}'_0 + \frac{\left(\delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f'_0) \right)^2}{2} f''_0 \right. \\
&\quad \left. + \frac{\delta^3}{6} \ddot{f}_0 + \frac{\delta^2}{2} (\delta f_0) \ddot{f}_0 + \frac{\delta}{2} (\delta f_0)^2 \dot{f}''_0 + \frac{(\delta f_0)^3}{6} f'''_0 \right] + O(\delta^5) \\
&= \delta f_0 + \delta^2 (\dot{f}_0 + f_0 f'_0) + \frac{\delta^3}{2} \left\{ \dot{f}_0 f'_0 + f_0 (f'_0)^2 + \ddot{f}_0 + 2 f_0 \dot{f}'_0 + f_0^2 f_0'' \right\} \\
&\quad + \frac{\delta^4}{8} \left(\ddot{f}_0 f'_0 + 6 f_0 \dot{f}'_0 f'_0 + 2 (\dot{f}_0 (f'_0)^2 + f_0 (f'_0)^3) + 5 f_0^2 f_0' f_0'' + 4 \dot{f}_0 \dot{f}'_0 \right. \\
&\quad \left. + 4 f_0 \dot{f}_0 f_0'' + \frac{4}{3} \ddot{f}_0 + 4 f_0 \ddot{f}'_0 + 4 f_0^2 \dot{f}''_0 + \frac{4}{3} f_0^3 f_0''' \right) + O(\delta^5). \tag{NM.8}
\end{aligned}$$

Thus, if we take the Runge-Kutta combination for the change in the function, we have

$$\begin{aligned}
\frac{1}{6} (k_1 + 2k_2 + 2k_3 + k_4) &= \delta f_0 + \frac{\delta^2}{2} (\dot{f}_0 + f_0 f'_0) \\
&\quad + \frac{\delta^3}{6} \left(\ddot{f}_0 + 2 f_0 \dot{f}'_0 + f_0 (f'_0)^2 + \dot{f}_0 f'_0 + f_0^2 f_0'' \right) \\
&\quad + \frac{\delta^4}{24} \left(\ddot{f}_0 + \dot{f}_0 f'_0 + 3 f_0 \ddot{f}'_0 + 3 \dot{f}_0 \dot{f}'_0 + 5 f_0 f_0' \dot{f}'_0 + 3 f_0^2 \dot{f}''_0 \right. \\
&\quad \left. + 3 f_0 \dot{f}_0 f_0'' + \dot{f}_0 (f'_0)^2 + 4 f_0^2 \dot{f}'_0 f_0'' + f_0 (f'_0)^3 + f_0^3 f_0''' \right). \tag{NM.9}
\end{aligned}$$

On the other hand we find by (tediously) taking derivatives directly

$$\begin{aligned}
 \dot{y}(t_0) &= f_0, \\
 \ddot{y}(t_0) &= \dot{f}_0 + f_0 f'_0, \\
 \ddot{y}(t_0) &= \ddot{f}_0 + 2f_0 \dot{f}'_0 + f_0 (f'_0)^2 + \dot{f}_0 f'_0 + f_0^2 f''_0, \\
 \dddot{y}(t_0) &= \dddot{f}_0 + \ddot{f}_0 f'_0 + 3f_0 \ddot{f}'_0 + 3\dot{f}_0 \dot{f}'_0 + 5f_0 f'_0 \dot{f}'_0 + 3f_0^2 \dot{f}''_0 \\
 &\quad + 3f_0 \dot{f}_0 f''_0 + \dot{f}_0 (f'_0)^2 + 4f_0^2 \dot{f}_0 f''_0 + f_0 (f'_0)^3 + f_0^3 f'''_0.
 \end{aligned}
 \tag{NM.10}$$

Thus, as advertised, the Runge-Kutta expression is the appropriate expansion, *i.e.*, the Taylor series. Clearly the Runge-Kutta notation is much more compact!