# Global Analysis w/ Invariant Mänifold Tube Transport 

## Organization of Talk

$\square$ Theory of Invariant Manifold Tubes

- Restricted Three Body Problem (R3BP)
$\square$ Computational Methods
- High order normal form expansions
- Monte Carlo sampling of energy surface
$\square$ Chemical Reaction Dynamics
- Electron scattering in the Rydberg atom
- Planar scattering of $\mathrm{H}_{2} \mathrm{O}$ with $\mathrm{H}_{2}$


## A Historical Perspective

- Appleyard [1970]: Invariant sets near unstable Lagrange points of R3BP.
- First picture of transport tube.



## Invariant Manifold Tubes

$\square$ What are tubes and where do they live?

- Geometry
$\square$ What do tubes do (prediction/control)?
- Dynamics


Figure from Gomez, Koon, Lo, Marsden, Masdemont, \& Ross 2001

## Restricted Three Body Problem




Figures from Marsden and Ross 2006
$\square$ Left: Fixed points viewed in rotating frame $\square$ Right: Hill's Region (potential energy surface)

## Low Energy Saddle Points



Effective Potential


Level set shows the Hill's region Figure from Koon, Lo, Marsden, \& Ross 2000
$\square$ Reduce out rotations and work at fixed ang. mom.
$\square \mathrm{L}_{1} \& \mathrm{~L}_{2}$ are low energy saddle points

- mediate transport from inner and outer realms


# $L_{1.2 .3}$ are Rank-1 Saddles 




$x$ (nondimensional units, rotating frame)

Figure from Koon, Lo, Marsden \& Ross 1999a

$$
H_{2}=\lambda q_{1} p_{1}+\frac{1}{2} \omega_{1}\left(q_{2}^{2}+p_{2}^{2}\right)+\frac{1}{2} \omega_{2}\left(q_{3}^{2}+p_{3}^{2}\right)
$$



planar oscillations projection

vertical oscillations projection

Figure from Gomez, Koon, Lo, Marsden, Masdemont, \& Ross 2001

## Rank-1 Saddle Geometry



planar oscillations projection

vertical oscillations projection
saddle projection
Figure from Gomez, Koon, Lo, Marsden, Masdemont, \& Ross 2001
$\square$ Energy is shared between saddle and two centers $S^{3} \cong\left\{\frac{\omega_{1}}{2}\left(q_{2}^{2}+p_{2}^{2}\right)+\frac{\omega_{2}}{2}\left(q_{3}^{2}+p_{3}^{2}\right)=H-\lambda q_{1} p_{2}\right\}$
 $+$

## Orbit Structures




Figure from Koon, Lo, Marsden, \& Ross 1999b
$\square$ Conley [1968]: Low energy transit orbits
$\square$ McGehee [1969]: Homoclinic orbits

## Symbolic Dynamics



Figure from Koon, Lo, Marsden, \& Ross 1999b
$\square$ Symbolic/horseshoe dynamics
Thm. [Koon, Lo, Marsden, Ross, Chaos 2000]:

- There is an orbit with any admissible itinerary
- Example: (...,X,J,S,S,J,X,...)


# Manifold Tube Intersections 



Figure from Gomez, Koon, Lo, Marsden, Masdemont, \& Ross 2001

## Patched Three Body Problem



Figure from Gomez, Koon, Lo, Marsden, Masdemont, \& Ross 2001
$\square$ Jupiter Icy Moons Orbiter (JIMO)
$\square$ Arbitrarily many flyby's of each moon

## Normal Forms

$\square$ Integrable approximation to chaotic dynamics
$\square$ Linearize Vector Field at fixed pt.

- $\dot{z}=D \mathbb{J} \nabla H(z)=A z ; \quad z=(q, p)$
- Matrix $A$ has eigenvalues $\pm \lambda, \pm i \omega_{1}, \pm i \omega_{2}, \ldots, \pm i \omega_{n}$
- $\pm i \omega_{k}$ corresponds to elliptic motion (center)
$\circ \pm \lambda$ corresponds to hyperbolic motion (saddle)
- Transport is governed by $\pm \lambda$ direction
$\square$ NF decouples saddle \& center modes to high order


Center Projections

Saddle Proiection
Figure from Gomez, Koon, Lo, Marsden, Masdemont, \& Ross 2001

# Normal Form at Rank-1 Saddle 

$\square$ Quadratic Normal Form:

$$
H_{2}=\lambda q_{1} p_{1}+i \frac{\omega_{1}}{2}\left(q_{2}^{2}+p_{2}^{2}\right)+i \frac{\omega_{2}}{2}\left(q_{3}^{2}+p_{3}^{2}\right)
$$

$\square$ Successive transformations eliminate $n^{\text {th }}$ order terms

- Computations use Lie Transform method:

$$
\hat{H}=H+\{H, G\}+\frac{1}{2!}\{\{H, G\}, G\}+\frac{1}{3!}\{\{\{H, G\}, G\}, G\}+
$$

- Each change depends only on $A$
- Kill all terms $q_{1}^{i} p_{1}^{j}$ for $i \neq j$
$\square$ Action-angle variables $\left(I=q_{1} p_{1}, \theta_{k}\right)$
- $I=0$ is reduction to center manifold
- $I=\epsilon$ nudges orbits in saddle direction
$\square$ Implemented for 3DOF systems by A. Jorba (1999)


## Other Methods

$\square$ Global Analysis of Invariant Objects (GAIO)

- Transfer operators on box subdivisions
- Tree structured box elimination
$\square$ Statistical Sampling of Trajectories
- Monte Carlo sample initial conditions from phase space box surrounding tubes
- Integrate forwards and backwards to determine which tubes the i.c. are in
- After a relatively small number of samples one obtains a good estimate of volume ratios
- Applies well to higher dimensional systems (~5 or ~10 DOF)


## Overview of Method

$\square$ Identify Saddle/TS \& Hill Region
$\square$ Find Box Bounding Reactive Trajectories (outcut)

- in \& out cuts make "airlock"
- Monte Carlo sample energy surface in box
$\square$ Integrate traj's into bound state until escape



## Bounding Box Method



Integrate Trajectories Backwards Until Out Cut


Refine Bounding Box
Until It Contains All
Reactive Trajectories

## Sampling the Energy Surface

- Randomly select points in bounding box
- Project (using momentum variables) until intersects energy surface

Bounding Box


## Transition State Theory



Reaction Coordinate

saddle projection

planar oscillations projection

vertical oscillations projection
$\square$ Transition State: Joins Reactants \& Products

- Bottleneck near rank-1 saddle
- Opens for energies larger than saddle
$\square$ TST Assumes Unstructured Phase Space - Even Chaotic Phase Space is Structured


## What is a Scattering Reaction?



Bound State

## Unbound State

$\square$ Bound vs. Unbound States (Hill Region)
$\square$ Zero Angular Momentum not always valid

## Example - Rydberg Atom

$$
\begin{gathered}
H=\frac{1}{2}\left(p_{x}^{2}+p_{y}^{2}+p_{z}^{2}\right)+\frac{1}{2}\left(x p_{y}-y p_{x}\right)+\frac{1}{8}\left(x^{2}+y^{2}\right) \\
-\epsilon x-\frac{1}{\sqrt{x^{2}+y^{2}+z^{2}}}
\end{gathered}
$$









- $\sim 3$ minutes :: $4,000 \mathrm{pts}::<.5 \%$ error
- $\sim 1$ hour :: 140,000 pts :: < . $1 \%$ error
- $\sim 2$ days :: 1,000,000 pts ::


# Planar Scattering of $\mathrm{H}_{2} \mathrm{O}-\mathrm{H}_{2}$ 

$$
H=\frac{p_{R}^{2}}{2 m}+\frac{\left(p_{\theta}-p_{\alpha}\right)^{2}}{2 m R^{2}}+\frac{\left(p_{\alpha}-p_{\beta}\right)^{2}}{2 I_{a}}+\frac{p_{\beta}^{2}}{2 I_{b}}+V
$$

- $V=$ dipole/quadrupole + dispersion + induction + Leonard-Jones. (Wiesenfeld, 2003)
- Reduce out $\theta$ and work on $p_{\theta} \equiv J>0$ level set.

Body Frame
Lab Frame



Fixed Axis Frame


## $\mathrm{H}_{2} \mathrm{O}-\mathrm{H}_{2}$ Saddles




planar oscillations projection

vertical oscillations projection
saddle projection
Linearization near rank-1 saddle

## $\mathrm{H}_{2} \mathrm{O}-\mathrm{H}_{2}$ Hill Region

## 





## $\mathrm{H}_{2} \mathrm{O}-\mathrm{H}_{2}$ Collision Dynamics




$\square$Unrealistic Potential?
$\square$ Numerically Volatile Collisions
$\square$ Is Non-Scattering Reaction Occurring?
$\square$ More Realistic Potential Surface (Wiesenfeld)

# $\mathrm{H}_{2} \mathrm{O}-\mathrm{H}_{2}$ Lifetime Distribution 




CollisionDistributionfor $\mathrm{H} 2 \mathrm{O}-\mathrm{H} 2(<1000)$


- Locally structured (fine scale)
- Globally RRKM (coarse scale)
- Does structure persist w/ error in energy samples?


## Gaussian Energy Sampling

$\square$ Experimental verification of lifetime distribution

- Fixed energy slice is not realistic
- Gaussian around target energy is more physical
- Do nonRRKM features persist?

LifetimeDistribution $\epsilon=.58$


Energy Distributionof SamplePoints


$$
\begin{aligned}
& \geq 3 \text { DOF Rydberg Analog } \\
& H=\frac{1}{2}\left(p_{x}^{2}+p_{y}^{2}+p_{z}^{2}+p_{w}^{2}\right)+\frac{1}{2}\left(x p_{y}-y p_{x}\right)+\frac{1}{8}\left(x^{2}+y^{2}\right) \\
& -\epsilon x-\frac{}{\sqrt{x^{2}+y^{2}+z^{2}+w^{2}}}
\end{aligned}
$$








## Comparison of Methods

$\square$ High Order Normal Form Expansion

- Compute Transit Tubes directly
- NF expansion becomes involved for > 3 DOF
$\square$ Almost Invariant Set Methods (GAIO)
- Transfer operators on box subdivisions
- Increasing memory demands w/ higher DOF
$\square$ Bounding Box Method
- Lifetime Distribution essentially 1D problem
- Scales well to higher DOF systems
- Integration \& sampling become bottleneck


## Future Work

$\square$ Tighter Bounding Box

$\square$ Variational Integrator

- Larger time steps, faster runtime
- Computes collisions more accurately
- Bulk of computation is integration
$\square$ Asteroid Capture Rates


## Conclusions \& Open Questions

$\square$ Conclusions

- Bounding Box Method is very efficient
- Requires minor modification for new systems
- Remains fast for high DOF systems
$\square$ Next Steps
- Apply method to higher DOF chemical system
- Obtain experimental verification of method
$\square$ Open Problems
- Is there an estimate for how small energy must be for linear dynamics to persist?
- Perron-Frobenius operator (coarse grained reaction coordinate)
- Apply tube dynamics to stochastic models
- Solve Rank-2 sampling problem (non-compact)


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## The End



## Questions...

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