# AMath 574 February 28, 2011

#### Today:

- Another example nonlinear system: Burgers' + Advection
- Shallow water Riemann solution

#### Next Monday:

- Finite volume methods
- Approximate Riemann solvers

Reading: Chapter 15

Another example of a nonlinear system:

$$q = \begin{bmatrix} u \\ v \end{bmatrix}, \qquad f(q) = \begin{bmatrix} \frac{1}{2}(u^2) \\ (u+1)v \end{bmatrix}.$$

This is simply Burgers' equation

$$u_t + \frac{1}{2}(u^2)_x = 0$$

coupled to conservative advection

$$v_t + ((u+1)v)_x = 0$$

But... Advection velocity u + 1 comes from solution of Burgers' equation.

Solving  $u_t + \frac{1}{2}(u^2)_x = 0$  gives rarefaction wave (if  $u_l < u_r$ ) or shock wave with speed  $s^1 = \frac{1}{2}(u_l + u_r)$  (if  $u_l > u_r$ ).

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Advection equation can be rewritten as

$$v_t + (u+1)v_x = -u_x v$$

and characteristic theory shows that

$$\frac{d}{dt}v(X(t),t) = -u_x(X(t),t)v(X(t),t)$$

along the curve X'(t) = u(X(t), t) + 1.

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In regions where u is constant:

Characteristics are straight lines,

 $u_x = 0 \implies v$  is constant.

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If delta moves a different speed than advection velocity, this leads to a jump in v at the shock location.

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If delta moves a different speed than advection velocity, this leads to a jump in v at the shock location.

Resonant case: If shock moves at same speed as advection velocity then delta function is stationary relative to advecting v and we expect solution to blow up!

Reconsider as nonlinear system:

$$q = \begin{bmatrix} u \\ v \end{bmatrix}, \qquad f(q) = \begin{bmatrix} \frac{1}{2}(u^2) \\ (u+1)v \end{bmatrix}.$$

Jacobian matrix:

$$f'(q) = \left[ \begin{array}{cc} u & 0 \\ v & u+1 \end{array} \right].$$

Always hyperbolic since  $u \neq u + 1$ .

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$$\begin{split} \lambda^1 &= u, \quad r^1 = \begin{bmatrix} 1 \\ -v \end{bmatrix}, \qquad \nabla \lambda^1 \cdot r^1 \equiv 1, \quad \text{genuinely nonlinear} \\ \lambda^2 &= u+1, \quad r^2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}, \qquad \nabla \lambda^2 \cdot r^2 \equiv 0, \quad \text{linearly degenerate} \end{split}$$

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$$\begin{split} \tilde{u}'(\xi) &= 0 & \implies \tilde{u}(\xi) = u_* \\ \tilde{v}'(\xi) &= v(\xi) \implies \tilde{v}(\xi) = v_* e^{\xi} \end{split}$$

Integral curves are vertical lines.

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Integral curves are vertical lines.

These lines are also contours of  $\lambda^2$  (linearly degenerate!)

We'll see later these are also the Hugoniot loci for 2-waves.

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$$\begin{split} \tilde{u}'(\xi) &= 1 & \Longrightarrow \ \tilde{u}(\xi) = u_* + \xi \\ \tilde{v}'(\xi) &= -v(\xi) \ \Longrightarrow \ \tilde{v}(\xi) = v_* e^{-\xi} \end{split}$$

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States q and  $q_*$  must satisfy Rankine-Hugoniot jump condition:

$$f(q) - f(q_*) = s(q - q_*)$$

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One solution:

 $u = u_*$  (and jump in v arbitrary)  $\implies$  vertical lines

These are Hugoniot loci for 2-waves.

2-waves are discontinuities in v alone, speed  $s = u_* + 1$ (determined from second equation of R-H conditions).

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Second solution:

$$s = s^1 = \frac{1}{2}(u + u_*) \implies$$
 shock waves in Burgers' equation

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Relation between v and u across shock:

Second equation of R-H relation:

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$$\implies v = \left(\frac{1 + \frac{1}{2}(u_* - u)}{1 - \frac{1}{2}(u_* - u)}\right)v_*$$

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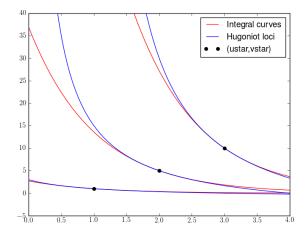
$$(u+1)v - (u_*+1)v_* = s(v-v_*) = \frac{1}{2}(u+u_*)(v-v_*)$$

$$\implies v = \left(\frac{1 + \frac{1}{2}(u_* - u)}{1 - \frac{1}{2}(u_* - u)}\right)v_* \approx e^{u_* - u}v_*$$

The Hugoniot locus agrees to  $\mathcal{O}(|u_* - u|^3)$  with integral curve.

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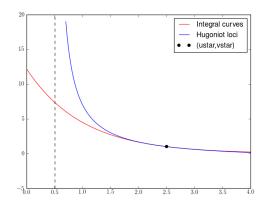
### Burgers' + advection: Phase plane



### Burgers' + advection: Phase plane

But note that

$$v = \left(\frac{1 + \frac{1}{2}(u_* - u)}{1 - \frac{1}{2}(u_* - u)}\right)v_* \to \infty \quad \text{as } u \to u_* - 2$$



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To be discussed on the board...

See also the description and codes at http://www.clawpack.org/links/burgersadv