

## Energy Methods

Conservation of energy:

If we have no losses in the system  
"mechanical" energy is conserved:

$$T + V = \text{constant}$$

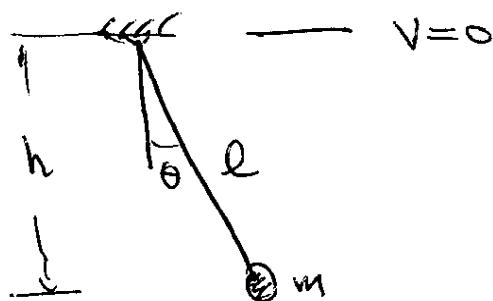
$T$  - kinetic energy

$V$  - Potential energy

Then  $\frac{d}{dt}(T+V) = 0$

"no change with time"

Ex Pendulum:



Potential Energy:  $V = -mgh = -mg l \cos \theta$

Kinetic Energy:  $T = \frac{1}{2}mv^2 = \frac{1}{2}m(l\dot{\theta})^2$

$$\text{so } \frac{d}{dt} \left( \frac{1}{2} m l^2 \dot{\theta}^2 - m g l \cos \theta \right) = 0$$

$$m l^2 \ddot{\theta} + m g l \sin \theta \dot{\theta} = 0$$

$$m l^2 \ddot{\theta} + m g l \sin \theta = 0$$

$$\ddot{\theta} + \frac{g}{l} \sin \theta = 0$$

The system is nonlinear because of the  $\sin \theta$  term

For small motions,  $|\theta| \ll 1$

then  $\sin \theta \approx \theta$

$$\text{so } \ddot{\theta} + \frac{g}{l} \theta = 0$$

$$\text{so } \omega_n = \sqrt{\frac{g}{l}}$$

## Lagrange Equation

Define  $L = T - V$  - the Lagrangian

The the equations of motion are

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}_k} \right) - \frac{\partial L}{\partial x_k} = Q_k \quad k=1, \dots, n$$

where  $Q_k = \sum_{i=1}^N \vec{F}_i \cdot \frac{d\vec{r}_i}{dx_k}$

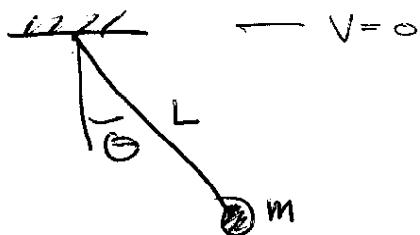
where  $\vec{F}_i$  = nonconservative forces acting  
on the  $i$ th particle (mass)

$\vec{r}_i = \vec{r}_i(x_1, x_2, x_3, \dots, t)$  Position vector  
to the  $i$ th particle

$x_k$  = the generalized coordinate

$N$  is the total number of masses

$n$  is the degree of freedom

Ex

Here  $x_1 = \theta$ ,  $n = 1$  (one degree of freedom)

$$V = -mgL \cos \theta$$

$$T = \frac{1}{2}mv^2 = \frac{1}{2}m(L\dot{\theta})^2$$

$$L = T - V = \frac{1}{2}m(L\dot{\theta})^2 + mgL \cos \theta$$

so Lagrange equation becomes

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{\theta}}\right) - \frac{\partial L}{\partial \theta} = 0$$

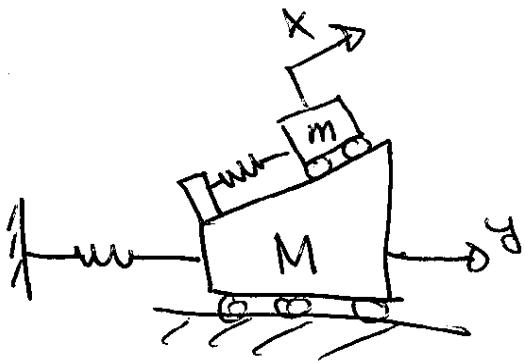
$$\frac{d}{dt}(mL^2\ddot{\theta}) - mgL(-\sin \theta) = 0$$

$$mL^2\ddot{\theta} + mgL \sin \theta = 0$$

$$\ddot{\theta} + \frac{g}{L} \sin \theta = 0 \quad \checkmark$$

or  $\ddot{\theta} + \frac{g}{L} \theta = 0$  for  $|\theta| \ll 1$

(3)

Ex 2No gravity,  $g=0$ Here  $x_1 = y$ ,  $x_2 = x$ ,  $n = 2$ 

$$V = \frac{1}{2}k_1y^2 + \frac{1}{2}k_2(x - y\cos\theta)^2$$

$$T = \frac{1}{2}m\dot{x}^2 + \frac{1}{2}M\dot{y}^2$$

$$L = \frac{1}{2}m\dot{x}^2 + \frac{1}{2}M\dot{y}^2 - \frac{1}{2}k_1y^2 - \frac{1}{2}k_2(x - y\cos\theta)^2$$

$$\frac{\partial}{\partial t}\left(\frac{\partial L}{\partial \dot{y}}\right) - \frac{\partial L}{\partial y} = 0 \Rightarrow$$

$$\frac{\partial}{\partial t}(M\cdot\ddot{y}) + k_1y + k_2(x - y\cos\theta)(-\cos\theta) = 0$$

$$\Rightarrow \boxed{M\ddot{y} + (k_1 + k_2\cos^2\theta)\cdot\ddot{y} - k_2\cos\theta\dot{x} = 0} \quad (1)$$

Also

$$\frac{\partial}{\partial t}\left(\frac{\partial L}{\partial \dot{x}}\right) - \frac{\partial L}{\partial x} = 0$$

$$\frac{\partial}{\partial t}(m\ddot{x}) + k_2(x - y\cos\theta)\cdot 1 = 0$$

$$\boxed{m\ddot{x} + k_2x - k_2\cos\theta\cdot\ddot{y} = 0} \quad (2) \Leftarrow$$

or in matrix form, Eqn (1) and (2)  $\Rightarrow$ 

$$\begin{pmatrix} M & 0 \\ 0 & m \end{pmatrix} \begin{pmatrix} \ddot{x} \\ \ddot{y} \end{pmatrix} + \begin{pmatrix} k_1 + k_2\cos^2\theta & -k_2\cos\theta \\ -k_2\cos\theta & k_2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$