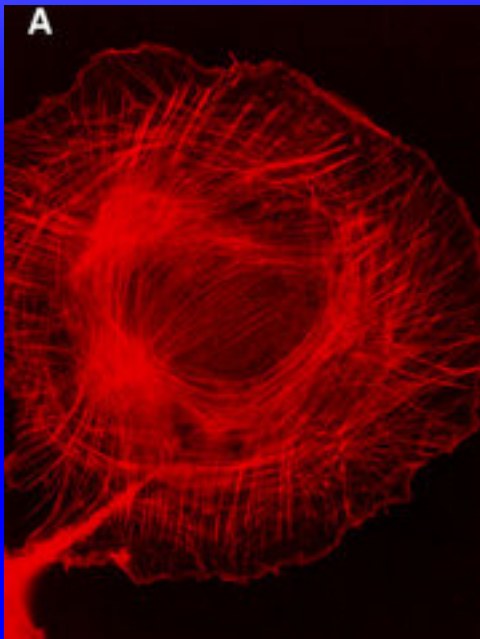


Session 16

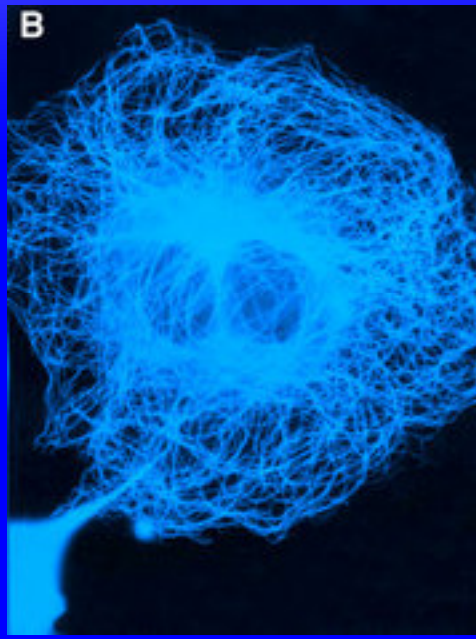
# SEMIFLEXIBLE POLYMER NETWORKS

# Recall: Cytoskeleton

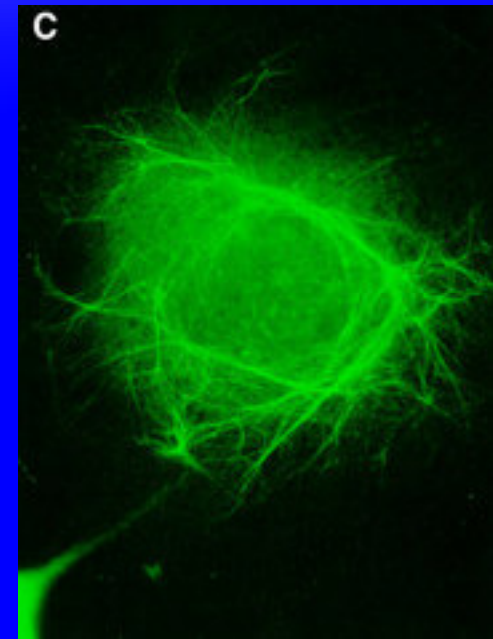
- ◆ Three fundamental cytoskeletal filaments



Actin



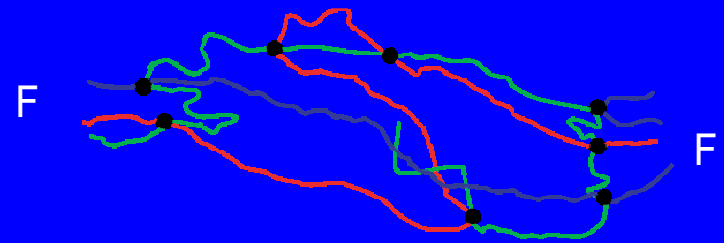
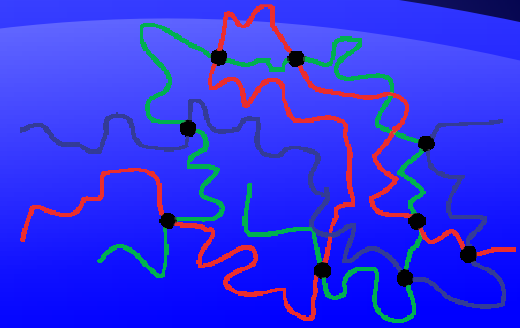
Microtubules (MT)



Intermediate Filaments (IF)

# Rubber (Elastomer)

- ◆ Random coiling of polymer filaments
- ◆ Cross-linking between filaments (black dots)
- ◆ Entropic Spring
  - ◆ Stretching force cause order in filaments
  - ◆ Reduces entropy
  - ◆ Generates of heat
  - ◆ (What if you add heat back in?)



$$F = dU/dl - T dS/dl$$

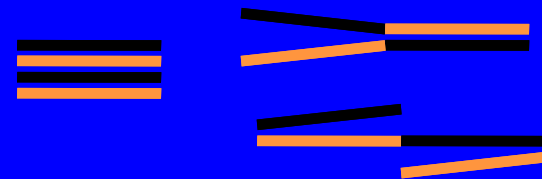
$$F = 0 - T dS/dl$$

# Entropy of Filaments

- ◆ Consider four-segment polymer

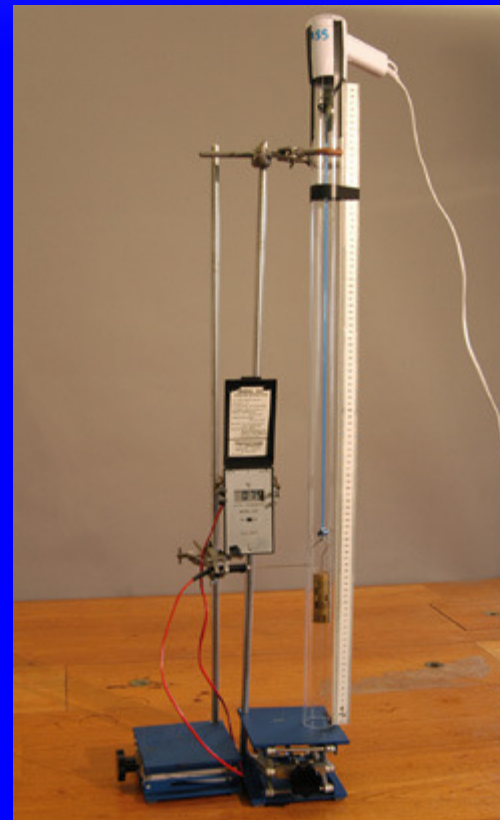
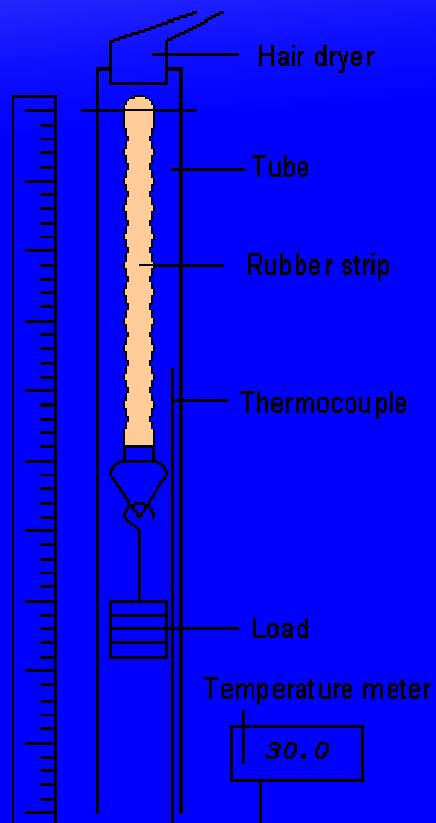


- ◆ There are 16 unique configuration states
  - ◆  $S \approx \ln(16) = 2.77$
- ◆ There tension is applied such that there cannot be zero separation from end-to-end
  - ◆  $S \approx \ln(10) = 2.30$
- ◆ Therefore, tension reduces entropy



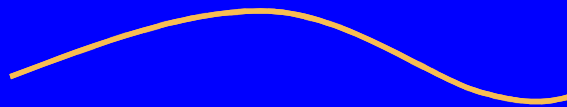
# Fun Experiment

- ◆ Entropic stiffness is proportional to temperature



# Persistence Length $l_p$

- ◆ Simply, filaments in thermal equilibrium in a liquid solution will appear
  - ◆ Straight over lengths  $< l_p$
  - ◆ Contorted randomly over lengths  $> l_p$



Hi  $l_p$



Lo  $l_p$

- ◆ In vitro solution measurements:
  - ◆ DNA: 50 nm
  - ◆ F-actin: 17  $\mu\text{m}$
  - ◆ Microtubule: 1 mm

# Persistence Length $l_p$

- ◆ Formally,

$$l_p = \frac{\kappa}{kT}$$

where

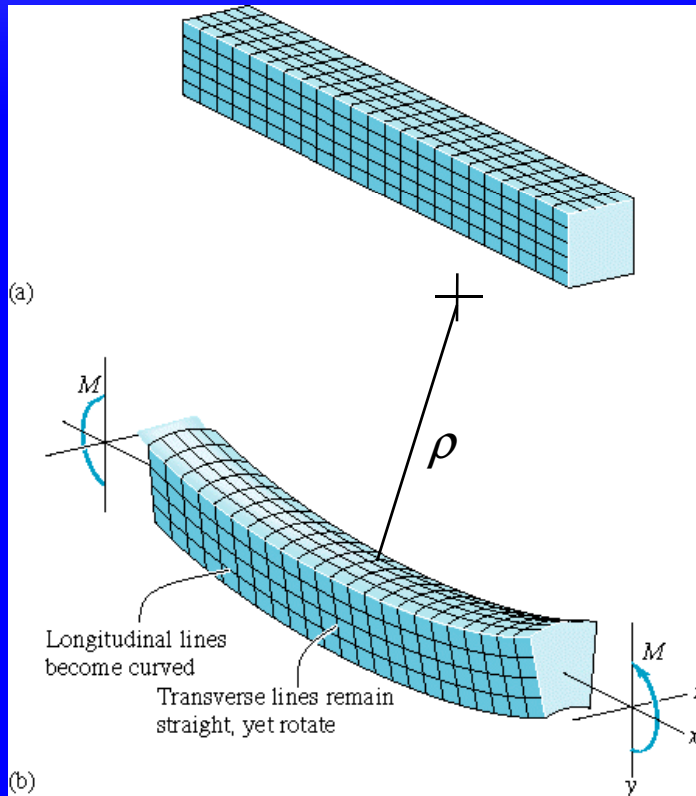
$\kappa = E_Y I$  is bending stiffness

$k = 1.38 \times 10^{-23}$  J/K is Boltzmann constant

T is temperature in K

# Bending Energy

- ◆ Applied forces lead to moment  $M$  in rod
- ◆  $M$  acts to deforms straight shape to curved



- ◆ Strain Energy

$$W = \frac{1}{2} F_x \Delta x$$

- ◆ Strain Energy per volume

$$U = \frac{W}{V} = \frac{1}{2} \sigma_x \epsilon_{xx}$$

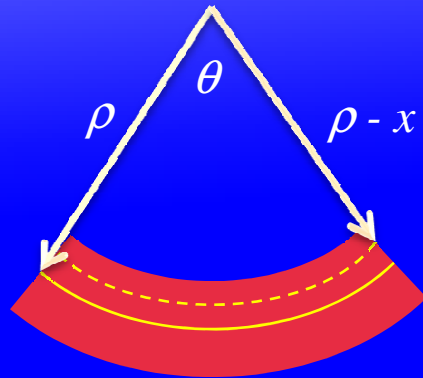
- ◆ By Hooke's Law

$$U = \frac{1}{2} E_Y \epsilon_{xx}^2$$



# Bending Curvature

- ◆ Interior bending lengths



$$L_{NA} = \rho\theta$$

$$L' = (\rho - x)\theta$$

- ◆ Strain

$$\epsilon_{xx} = \frac{(L' - L)}{L} = -\frac{x}{\rho}$$

- ◆ Thus...

$$U = \frac{1}{2} E_Y \left( \frac{x}{\rho} \right)^2$$

# Bending Energy

- ◆ Strain energy per unit length

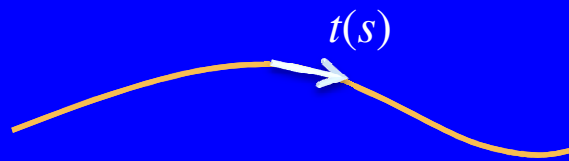
$$\begin{aligned} E &= \int_A \frac{1}{2} E_Y \frac{x^2}{\rho^2} dA \\ &= \frac{1}{2} \frac{E_Y}{\rho^2} \int_A x^2 dA \\ &= \frac{1}{2} \frac{E_Y I_y}{\rho^2} \\ &= \frac{\kappa}{2} \frac{1}{\rho^2} \end{aligned}$$

# Worm-like Chain Model

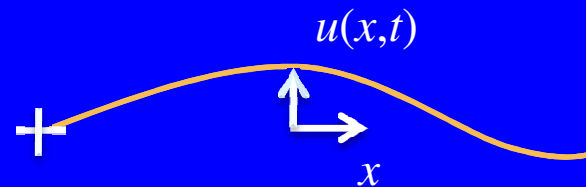
- ◆ Energy can be described as

$$H_{bend} = \frac{\kappa}{2} \int_l \frac{1}{\rho^2} ds$$

- ◆ Radius of curvature



$$\frac{1}{\rho} = \frac{\partial \bar{t}}{\partial s}$$



$$\frac{1}{\rho} \approx \frac{\partial^2 u}{\partial x^2}$$

# Worm-like Chain Model

- ◆ We can express as

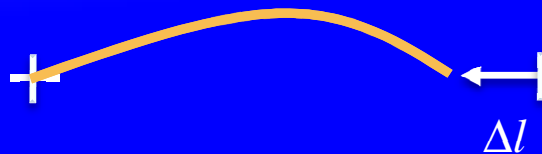
$$H_{bend} = \frac{\kappa}{2} \int_l \left( \frac{\partial \bar{t}}{\partial s} \right)^2 ds$$

- ◆ Or

$$H_{bend} = \frac{\kappa}{2} \int_l \left( \frac{\partial^2 u}{\partial x^2} \right)^2 dx$$

# Force – Extension Relationship

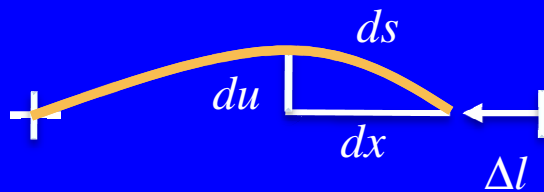
- ◆ Tension to extend filament measured by amount of extension along a line



- ◆ Extension response dominated by entropy
- ◆ At any finite temperature, there is contraction due to thermal fluctuations that make polymer deviate from straight line

# Force – Extension Behavior

- ◆ Semiflexible chain stretched by tension  $\tau$ 
  - ◆ Energy of bending  $H_{bend}$
  - ◆ Energy of contraction against  $\tau$
- ◆ Derive the “shortening” in the filament:  
*Length is unchanged by  $kT$ ,  $\tau$*



$$\begin{aligned}\Delta l &= \int (ds - dx) \\ &= \int \left( \sqrt{dx^2 + du^2} - dx \right) \\ &= \int \left( \left( \sqrt{1 + \left( \frac{\partial u}{\partial x} \right)^2} \right) - 1 \right) dx \\ &\approx \frac{1}{2} \int \left( \frac{\partial u}{\partial x} \right)^2 dx\end{aligned}$$

# Deformation Behavior

- ◆ Strain energy from stretching

$$H_{stretch} = \frac{1}{2} \int_l (\tau \cdot \Delta l) dx = \frac{1}{2} \int_l \tau \left( \frac{\partial u}{\partial x} \right)^2 dx$$

- ◆ Thus, total energy

$$\begin{aligned} H &= H_{bend} + H_{stretch} \\ &= \frac{1}{2} \int_l \left( \kappa \left( \frac{\partial^2 u}{\partial x^2} \right)^2 + \tau \left( \frac{\partial u}{\partial x} \right)^2 \right) dx \end{aligned}$$

- ◆ Note, there exists transverse motion direction version

# Spring Constants

- ◆ It can be shown ...
  - ◆ Transverse spring constant

$$k_{sp,bend} \simeq \frac{kTl_p}{l^3}$$

- ◆ Longitudinal spring constant

$$k_{sp,stretch} \simeq \frac{kTl_p^2}{l^4} = k_{sp,bend} \frac{l_p}{l}$$

- ◆ For  $l < l_p$ , longitudinal compliance is smaller than transverse



# Force – Extension Behavior

- ◆ It can be shown
  - ◆ Considering the balance with thermal energy
  - ◆ Adding transverse motion direction, and
  - ◆ Utilizing statistical mechanics concepts

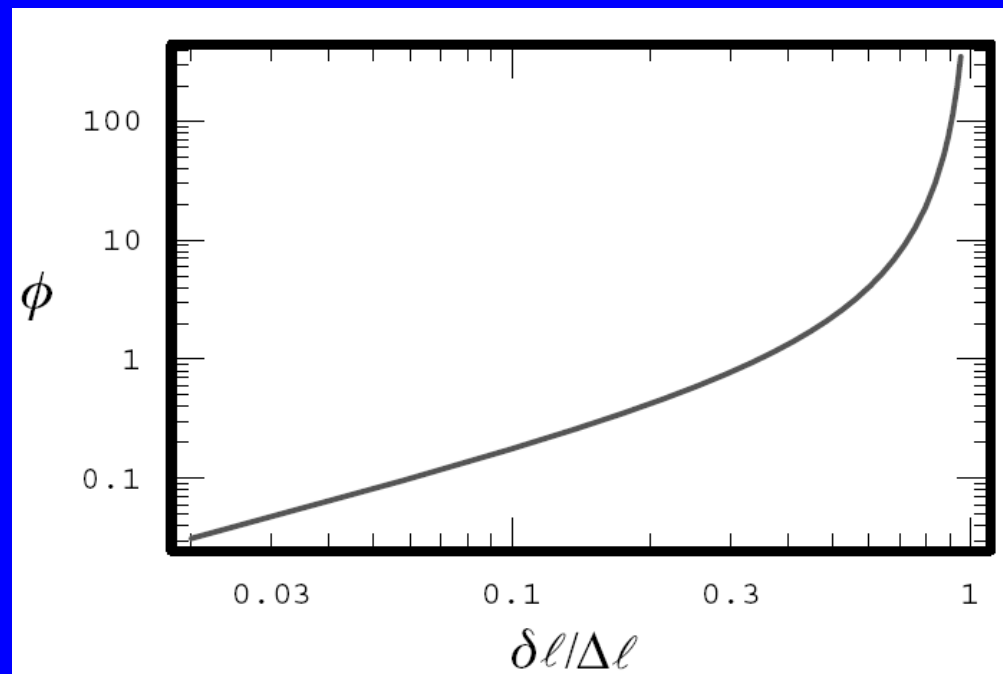
$$\delta l = \frac{kT}{\pi^2} \frac{l^2}{\kappa} \sum_n \frac{\phi}{n^2 (n^2 + \phi)}$$

with dimensionless force

$$\phi = \tau l^2 / \kappa \pi^2$$

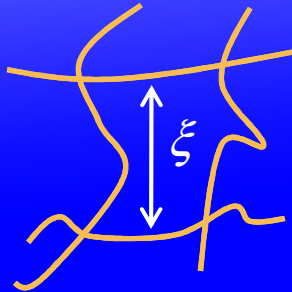
# Force - Extension

- ◆ Non-linear behavior
  - ◆ Linear at small force
  - ◆ Strain stiffening at large forces



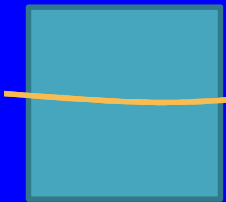
# Polymer Interactions

- ◆ Solution of polymers have mesh size dimension  $\xi$



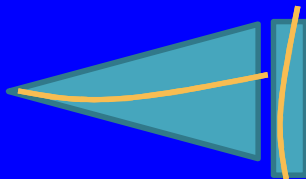
- ◆ Typical spacing between filaments
- ◆ Estimated by volume fraction  $\psi$
- ◆ For rigid rods (lengths  $< l_p$ )

$$\psi = \frac{V_{\text{filament}}}{V_{\text{mesh}}} = \frac{a^2 \xi}{\xi^3}$$



- ◆ Accounting for thermal fluctuations

$$\psi = \frac{V_{\text{transverse}}}{V_{\text{longitudinal}}} = \frac{a^2 \delta u}{\frac{1}{3} \pi l \delta u^2} = \frac{a^2}{l \delta u} = \frac{a^2}{l_e \delta u}$$



# Shear Modulus for Solutions

- ◆ It can be shown that for entropy only

$$G \sim kT / (\xi^2 l_e) \sim \psi^{7/5}$$

- ◆ Temperature dependence
- ◆ Filament density dependence

# Shear Modulus for Solutions

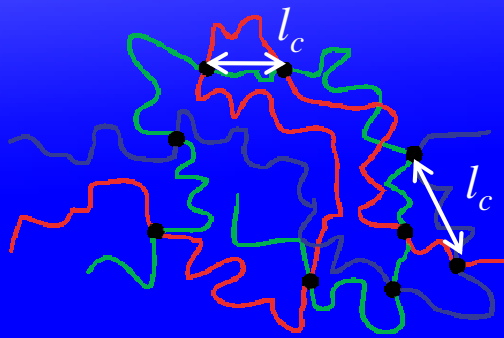
- ◆ High frequency testing
  - ◆ Filaments not able to relax from high bending modes
  - ◆ Increased stiffness from less compliant filaments

$$G(\omega) = \frac{1}{15} \rho \kappa l_p (-2i\zeta / \kappa)^{3/4} \omega^{3/4} - i\omega\eta$$

$\rho$  is polymer concentration,  $\zeta$  is hydrodynamic drag (per unit length),  $\eta$  is viscosity

# Shear Modulus for Networks

- ◆ Consider cross-linking distance  $l_c$



- ◆ Deformation types
  - ◆ Affine network models: uniform rotation or stretching
  - ◆ Non-affine models: macroscopic strains vary from one region to another

# Shear Modulus for Networks

- ◆ Affine, thermal-entropic (AT)

$$G_{NA} \sim \kappa / \xi^4 \sim \psi^2$$

- ◆ Modulus depends strongly on x-linking
- ◆ Non-affine (NA)
  - ◆ Low poly conc, high x-linking (low  $l_c$ )

$$G_{AT} \sim \kappa l_p / \xi^2 l_c^2$$

- ◆ Affine, mechanical (AM)
  - ◆ Filament segments (small  $l_c$ ) behave as rigid rods with modulus  $\mu$

$$G_{AM} \sim \mu / \xi^2 \sim \psi$$

