

University of Washington
Department of Chemistry
Chemistry 452/456
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Lecture 17 8/06/14

A. Regular Solutions: A Simple Example of a Real Solution

- The simplest non-ideal solution model that works beyond the Henry's Law model is the regular solution model. The basic assumption of the simplest regular solution model is that when components A and B mix, they mix randomly.

- If that is the case, the for both the simple regular and ideal binary solution

$$\Delta S_{MIX} = -R(\chi_A \ln \chi_A + \chi_B \ln \chi_B) > 0 \quad (17.1)$$

- If mixing is not random then the entropy will deviate from this expression.
- Ideal solution always form when components are mixed. To show this start with

$$\Delta G_{MIX} = \Delta H_{MIX} - T\Delta S_{MIX} \quad (17.2)$$

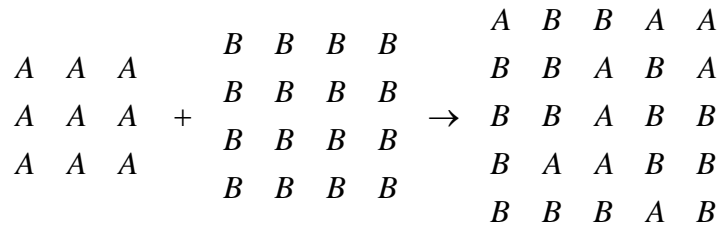
- Recall for an ideal solution...

$$\Delta H_{MIX} = 0 \quad (17.3)$$

- Given 17.1 -17.3 we conclude

$$\Delta G_{MIX} = -T\Delta S_{MIX} = RT(\chi_A \ln \chi_A + \chi_B \ln \chi_B) < 0 \quad (17.4)$$

- To construct a regular solution model we recall an ideal solution is a solution where all interactions between solution components are equal. In a solution composed of two components A and B, call the interaction between two A molecules, ϵ_{AA} , between 2 B molecules ϵ_{BB} , and between an A and a B ϵ_{AB} . If these three interaction energies are equal, the solution is ideal. If they are not the same, the solution's thermodynamic behavior is non-ideal.
- Suppose we imagine mixing N_A molecules A with N_B molecules B to form a solution. The solution can be imagined as N_A A molecules and N_B B molecules arranged on a lattice:



- The energy of such a solution now depends on the number of AA interactions N_{AA} , BB interactions N_{BB} , and AB interactions N_{AB} . In general the energy is

$$U(A, B) = N_{AA}\epsilon_{AA} + N_{AB}\epsilon_{AB} + N_{BB}\epsilon_{BB} \quad (17.5)$$

- This lattice-based theory of solutions is sometimes called regular solution theory. A regular solution is non-ideal but the interactions between molecules are weak, and all molecules are about the same size.

- An important number in the theory of regular solutions is the number of sides for a particular lattice position. This can be imagined to be the number of neighbors a A or B molecules has and in this simple lattice model the number is Z for A and B. The number of neighbors around N_A molecules is therefore ZN_A and can be expressed in terms of the number of AA and AB contacts N_{AA} and N_{AB} :

$$ZN_A = 2N_{AA} + N_{AB} \quad (17.6)$$

- A similar expression exists for the number of neighbors around NB B molecules:

$$ZN_B = 2N_{BB} + N_{AB} \quad (17.7)$$

- Now we solve for N_{AA} and N_{BB} and put the results into the energy equation:

$$\begin{aligned} U(A, B) &= N_{AA}\epsilon_{AA} + N_{AB}\epsilon_{AB} + N_{BB}\epsilon_{BB} \\ &= \left(\frac{ZN_A - N_{AB}}{2}\right)\epsilon_{AA} + N_{AB}\epsilon_{AB} + \left(\frac{ZN_B - N_{AB}}{2}\right)\epsilon_{BB} \\ &= \left(\frac{Z\epsilon_{AA}}{2}\right)N_A + \left(\frac{Z\epsilon_{BB}}{2}\right)N_B + \left(\epsilon_{AB} - \frac{\epsilon_{AA} + \epsilon_{BB}}{2}\right)N_{AB} \end{aligned} \quad (17.8)$$

- The only quantity in this equation that we do not know in general is N_{AB} , because the number of AB contacts depends of the specific arrangement of A and B molecules on the lattice. But we can get this number using statistics. The probability that a given A molecule has a B neighbor is:

$$p_B = \frac{N_B}{N_A + N_B} = \frac{N_B}{N} \quad (17.9)$$

- The number of AB contacts is there the number of neighbors around N_A A molecules, times the probability that any of those neighbors is a B molecule;

$$N_{AB} = ZN_A p_B = \frac{ZN_A N_B}{N} \quad (17.10)$$

- The energy equation is now:

$$\begin{aligned} U(A, B) &= \left(\frac{Z\epsilon_{AA}}{2}\right)N_A + \left(\frac{Z\epsilon_{BB}}{2}\right)N_B + \left(\epsilon_{AB} - \frac{\epsilon_{AA} + \epsilon_{BB}}{2}\right)N_{AB} \\ &= \left(\frac{Z\epsilon_{AA}}{2}\right)N_A + \left(\frac{Z\epsilon_{BB}}{2}\right)N_B + \left(\epsilon_{AB} - \frac{\epsilon_{AA} + \epsilon_{BB}}{2}\right)\left(\frac{ZN_A N_B}{N}\right) \\ &= \left(\frac{Z\epsilon_{AA}}{2}\right)N_A + \left(\frac{Z\epsilon_{BB}}{2}\right)N_B + wRT\left(\frac{N_A N_B}{N}\right) \\ \text{i.e. } w &= \frac{Z}{RT}\left(\epsilon_{AB} - \frac{\epsilon_{AA} + \epsilon_{BB}}{2}\right) \end{aligned} \quad (17.11)$$

- The energy equation has three parts which we write as:

$$U(A, B) = \left(\frac{Z\varepsilon_{AA}}{2} \right) N_A + \left(\frac{Z\varepsilon_{BB}}{2} \right) N_B + wRT \left(\frac{N_A N_B}{N} \right) \quad (17.12)$$

$$= U(A) + U(B) + \Delta U_{MIX}(A, B)$$

- The first two terms are the energies of the separated components. The third term arises from the mixing of the components:

$$\Delta U_{MIX} = wRT \left(\frac{N_A N_B}{N} \right) \quad (17.13)$$

- Recall also that that $\Delta H = \Delta U + \Delta(PV) \approx \Delta U$ if pressure and volumes are kept constant for the solution. Then for a regular solution:

$$\Delta H_{MIX} \approx \Delta U_{MIX} = wRT \left(\frac{N_A N_B}{N} \right) \quad (17.14)$$

- Notice that $w = \frac{Z}{RT} \left(\varepsilon_{AB} - \frac{\varepsilon_{AA} + \varepsilon_{BB}}{2} \right)$ can be positive or negative. So if $w < 0$ the mixing is exothermic and if $w > 0$ the mixing is endothermic.
- Now we want the chemical potential for a solute in a regular solution. Because $\Delta U_{MIX} \approx \Delta H_{MIX}$ then $\Delta A_{MIX} \approx \Delta G_{MIX}$. The Helmholtz (Gibbs) energy of mixing is obtained from:

$$\Delta G_{MIX} \approx \Delta A_{MIX} = \Delta U_{MIX} - T\Delta S_{MIX} = wRT \left(\frac{N_A N_B}{N} \right) + RT(N_A \ln \chi_A + N_B \ln \chi_B)$$

$$= wNRT \chi_A \chi_B + NRT(\chi_A \ln \chi_A + \chi_B \ln \chi_B)$$

$$\therefore \frac{\Delta A_{MIX}}{NRT} = \chi_A \ln \chi_A + \chi_B \ln \chi_B + w\chi_A \chi_B = \chi_A \ln \chi_A + (1 - \chi_A) \ln(1 - \chi_A) + w\chi_A(1 - \chi_A) \quad (17.15)$$

- Clearly the first two terms in the equation correspond to the Helmholtz energy of mixing for an ideal solution (i.e. $w=0$). The third term represents deviations from ideality. We can get the chemical potential for A or B by differentiating with respect to N_A or N_B . Start with

$$\mu_A - \mu_A^0 = \left(\frac{\partial A_{MIX}}{\partial N_A} \right)_{V, T, N_B} \approx \left(\frac{\partial G_{MIX}}{\partial N_A} \right)_{P, T, N_B}$$

$$= \frac{\partial}{\partial N_A} \left(RT \left(N_A \ln \left(\frac{N_A}{N_A + N_B} \right) + N_B \ln \left(\frac{N_B}{N_A + N_B} \right) \right) + wRT \left(\frac{N_A N_B}{N_A + N_B} \right) + \right)$$

$$\begin{aligned}
&= RT \frac{\partial}{\partial N_A} \left(N_A \ln(N_A) + N_B \ln(N_B) - (N_A + N_B) \ln(N_A + N_B) \right) + wRT \left(\frac{N_B}{N_A + N_B} - \frac{N_A N_B}{(N_A + N_B)^2} \right) \\
&= RT \left(\ln(N_A) + 1 - 1 - \ln(N_A + N_B) \right) + wRT \left(\frac{N_B^2}{(N_A + N_B)^2} \right) = RT \ln \chi_A + wRT \chi_B^2
\end{aligned}
\tag{17.16}$$

- It is a lot of calculus but the result is simple and gives the activity coefficients for A and B:

$$\mu_A - \mu_A^0 = RT \ln \chi_A + wRT \chi_B^2 = RT \ln \chi_A + RT \ln \gamma_A$$

$$\therefore \ln \gamma_A = w \chi_B^2$$

Similarly:

(17.17)

$$\mu_B - \mu_B^0 = RT \ln \chi_B + wRT \chi_A^2 = RT \ln \chi_B + RT \ln \gamma_B$$

$$\therefore \ln \gamma_B = w \chi_A^2$$

- We can also consider deviations from ideal behavior produced by properties of the regular solution. The Helmholtz energy can be divided into two parts: a part that represents ideal solution behavior ΔA_{MIX}^{ideal} and a second part that contains w and represents non-ideal conditions. This non-ideal part is sometimes called the excess Helmholtz energy ΔA_{MIX}^E :

$$\begin{aligned}
\Delta A_{MIX} &= \Delta A_{MIX}^{ideal} + \Delta A_{MIX}^E = NRT (\chi_A \ln \chi_A + \chi_B \ln \chi_B) + N \Delta \varepsilon \chi_A \chi_B \\
\text{where: } \Delta \varepsilon &= Z \left(\varepsilon_{AB} - \frac{\varepsilon_{AA} + \varepsilon_{BB}}{2} \right) = wRT
\end{aligned}
\tag{17.18}$$

- Once you have an expression for ΔA_{MIX} you can get an expression for the entropy, enthalpy, and heat capacity for regular solutions:

$$\Delta S_{MIX} = - \frac{\partial \Delta A_{MIX}}{\partial T} = \Delta S_{MIX}^{ideal} + \Delta S_{MIX}^E = -NR (\chi_A \ln \chi_A + \chi_B \ln \chi_B) - N \chi_A \chi_B \frac{\partial \Delta \varepsilon}{\partial T} \tag{17.19}$$

where again there is a part of the entropy that is the ideal entropy of mixing, which assumes random mixing, and a second part that represents non-ideal conditions. This non-ideal part is called the excess entropy of mixing ΔS_{MIX}^E .

- The enthalpy of mixing is easy to get from the fact that the volume is assumed to not change much upon mixing. So for constant volume and pressure

$$\begin{aligned}
\Delta H_{MIX} &= \Delta U_{MIX} + \Delta(PV) = \Delta U_{MIX} = \Delta A_{MIX} + T \Delta S_{MIX} \\
&= N \chi_A \chi_B \left(\Delta \varepsilon - T \frac{\partial \Delta \varepsilon}{\partial T} \right)
\end{aligned}
\tag{17.20}$$

Note the enthalpy of mixing is entirely dependent on non-ideal effect as it must be because for ideal solutions $\Delta H_{MIX}=0$.

- The heat capacity at constant pressure is also easy to

$$\text{get: } \Delta C_{P,MIX} = \left(\frac{\partial \Delta H_{MIX}}{\partial T} \right)_P = -N \chi_A \chi_B T \frac{\partial^2 \Delta \epsilon}{\partial T^2}$$

- Again, the heat capacity of mixing for a regular solution, which is derived from the enthalpy of mixing, depends entirely on non-ideal effects.