

University of Washington
Department of Chemistry
Chemistry 452/456
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Lecture 10.1 7/13/11

A. Free Energy and Changes in Composition: The Chemical Potential

- We have thus far considered free energy changes in single component systems, i.e. free energy changes accompanying a changes in pressure of an ideal gas. But many chemical and biochemical systems are naturally composed of a number of components and the composition of such systems may change as a result of chemical reactions, physical transport, etc. In such cases we must consider how the free energy changes when compositions changes. Consider the reaction $n_A A \rightarrow n_B B$. The free energy change is given by

$$\begin{aligned}
 dG &= \left(\frac{\partial G}{\partial P} \right)_{T, n_A, n_B} dP - \left(\frac{\partial G}{\partial T} \right)_{P, n_A, n_B} dT + \left(\frac{\partial G}{\partial n_A} \right)_{P, T, n_B} dn_A + \left(\frac{\partial G}{\partial n_B} \right)_{P, T, n_A} dn_B \\
 &= VdP - SdT + \mu_A dn_A + \mu_B dn_B
 \end{aligned}$$

- In general $dG = VdP - SdT + \sum_i \mu_i dn_i$
- The terms $\mu_i = \left(\frac{\partial G}{\partial n_i} \right)_{P, T, n_j}$ are called a chemical potentials.
- The chemical potential of a pure substance is simply the molar Gibbs energy. For a pure substance: $\mu = \left(\frac{\partial G}{\partial n} \right)_{P, T} = \frac{G}{n} = \bar{G}$
- We can relate partial molar quantities using the same derivative equations that we used for thermodynamic state functions. Again for a pure substance: $\left(\frac{\partial \bar{G}}{\partial P} \right)_T = \left(\frac{\partial \mu}{\partial P} \right)_{T, i} = \frac{V}{n} = \bar{V}$.
- Similarly we have $\left(\frac{\partial \bar{G}}{\partial T} \right)_P = \left(\frac{\partial \mu}{\partial T} \right)_P = -\frac{S}{n} = \bar{S}$
- And the chemical potential is an exact differential so that

$$d\mu = \left(\frac{\partial \mu}{\partial T} \right)_P dT + \left(\frac{\partial \mu}{\partial P} \right)_T dP = -\bar{S}dT + \bar{V}dP$$
- This equation gives us a method for calculating the chemical potential of a one component ideal gas. Suppose the temperature is a constant, i.e. $dT=0$. Then:

$$d\mu = -\bar{S}dT + \bar{V}dP = \bar{V}dP$$

$$\therefore \int_{\mu^0}^{\mu_i} d\mu = \mu - \mu^0 = \int_{P^0}^P \bar{V}dP = \int_{P^0}^P \frac{V}{n} dP = RT \int_{P^0}^P \frac{dP}{P} = RT \ln\left(\frac{P}{P^0}\right)$$

$$\therefore \mu = \mu^0 + RT \ln\left(\frac{P}{P^0}\right)$$

- μ^0 is the chemical potential of a standard reference state for the gas which is characterized by the pressure P^0 , which is frequently taken to be 1 bar.

B. Chemical Potential in a Mixture

- So far the chemical potential just seems to be an alternative notation for performing molar Gibbs energy calculations. The real utility of the chemical potential is in describing thermodynamic properties of a mixture. Consider a mixture of ideal gases which satisfies: $PV = \sum_i n_i RT$. Here we define the partial pressures P_i

$$P = \sum_i P_i = \sum_i \chi_i P \quad \text{where the mole fraction } \chi_i = \frac{n_i}{\sum_i n_i} = \frac{n_i}{n}$$

- For the i th component of an ideal gas mixture we define the chemical potential as $d\mu_i = -\bar{S}_i dT + \bar{V}_i dP$. We calculate this chemical potential at constant T , i.e. $dT=0$:

$$\mu_i - \mu_i^0 = \int_{P^0}^{P_i} \bar{V}_i dP = \int_{P^0}^{P_i} \left(\frac{\partial V}{\partial n_i} \right)_{T,P,n_j} dP = \int_{P^0}^{P_i} \frac{RT}{P} \left(\frac{\partial n}{\partial n_i} \right)_{T,P,n_j} dP = RT \ln\left(\frac{P_i}{P^0}\right)$$

$$\therefore \mu_i = \mu_i^0 + RT \ln\left(\frac{P_i}{P^0}\right)$$

C. Criteria for Equilibrium

- The chemical potential is used to express criteria for equilibrium. A very simple example is a single component system containing two phases α and β . At constant pressure and temperature and with the system closed:

$$dG = \mu_\alpha dn_\alpha + \mu_\beta dn_\beta$$

- If we convert β to α we must conserve mass so $dn_\alpha = -dn_\beta$ so that $dG = \mu_\alpha dn_\alpha + \mu_\beta dn_\beta = (\mu_\alpha - \mu_\beta) dn_\alpha$. The direction of the phase change is dictated by the difference $\mu_\alpha - \mu_\beta$. If $\mu_\alpha > \mu_\beta$ then $dG > 0$. If $\mu_\alpha < \mu_\beta$ then $dG < 0$. Clearly the direction of phase change in this closed system is designated by “transport” into the phase with lowest chemical potential. However if $\mu_\alpha - \mu_\beta = 0$ i.e. the chemical potentials are equal, then $dG = 0$ and the system is at equilibrium.

- The criterion for equilibrium in a reacting mixture is somewhat more involved. Assume a chemical reaction $\nu_1 X_1 + \nu_2 X_2 + \dots + \nu_{j-1} X_{j-1} \rightarrow \nu_j X_j + \dots + \nu_{n-1} X_{n-1} + \nu_n X_n$ where ν_i are stoichiometric coefficients. We can represent this reaction as an algebraic equation $\sum_{i=1}^n \nu_i X_i = 0$ where $\nu_i < 0$ for reactants (i.e. $i < j$) and $\nu_i > 0$ for productions (i.e. $i \geq j$).
- Now we define the reaction coordinate ξ which varies between 0 and 1 such that $dn_i = \nu_i d\xi$. Then for $dT=dP=0$

$$dG = \sum_{i=1}^n \mu_i dn_i = \sum_{i=1}^n \mu_i \nu_i d\xi = d\xi \sum_{i=1}^n \mu_i \nu_i$$

$$\therefore \left(\frac{\partial G}{\partial \xi} \right)_{P,T} = \sum_{i=1}^n \mu_i \nu_i$$

- At equilibrium $\left(\frac{\partial G}{\partial \xi} \right)_{P,T} = \sum_{i=1}^n \mu_i \nu_i = 0$ and the criterion for equilibrium in a reacting mixture is: $\sum_{i=1}^n \mu_i \nu_i = 0$. If $\sum_{i=1}^n \mu_i \nu_i \neq 0$ the reacting mixture is not at equilibrium.
- If we use our expression for the chemical potential for components of an ideal gas mixture:

$$\Delta G = \sum_{i=1}^n \nu_i \mu_i = \sum_{i=1}^n \nu_i \left(\mu_i^0 + RT \ln \left(\frac{P_i}{P^0} \right) \right)$$

$$= \sum_{i=1}^n \nu_i \mu_i^0 + RT \sum_{i=1}^n \nu_i \ln \left(\frac{P_i}{P^0} \right) = \Delta G^0 + RT \sum_{i=1}^n \ln \left(\frac{P_i}{P^0} \right)^{\nu_i}$$

$$= \Delta G^0 + RT \ln \prod_{i=1}^n \left(\frac{P_i}{P^0} \right)^{\nu_i}$$

- If the system is at equilibrium

$$\Delta G = \left(\frac{\partial G}{\partial \xi} \right)_{P,T} = 0 = \Delta G^0 + RT \ln \prod_{i=1}^n \left(\frac{P_{i,eq}}{P^0} \right)^{\nu_i}$$

$$\therefore \Delta G^0 = -RT \ln \prod_{i=1}^n \left(\frac{P_{i,eq}}{P^0} \right)^{\nu_i} = -RT \ln K_p$$

where the equilibrium constant $K_p = \prod_{i=1}^n \left(\frac{P_{i,eq}}{P^0} \right)^{\nu_i}$.

- If the system is not at equilibrium:

$$\Delta G = \left(\frac{\partial G}{\partial \xi} \right)_{P,T} = \Delta G^0 + RT \ln \prod_{i=1}^n \left(\frac{P_i}{P^0} \right)^{\nu_i} = \Delta G^0 + RT \ln Q_p$$

Note: Q is a dimensionless quantity called the reaction quotient. K is also dimensionless.