

Bose-Einstein Condensation in a Circular Waveguide

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We have produced Bose-Einstein condensates in a ring-shaped magnetic waveguide. The few-millimeter diameter, nonzero-bias ring is formed from a time-averaged quadrupole ring. Condensates that propagate around the ring make several revolutions within the time it takes for them to expand to fill the ring. The ring shape is ideally suited for studies of vorticity in a multiply connected geometry and is promising as a rotation sensor.

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The long range phase coherence of superfluids and superconductors, aside from leading to interesting physical effects, is also of practical importance in allowing for precise measurement devices based on quantum interference. In such devices, spatially separated paths forming a multiply connected geometry are imposed on the macroscopic quantum system. For example, a SQUID magnetometer makes use of a superconducting ring interrupted by Josephson junctions to allow continuous sensitivity to magnetic fields. A similar geometry was used in a superfluid ^3He gyroscope [1].

Dilute gas superfluids have now enabled novel forms of matter-wave interferometry. Precise sensors of rotation, acceleration, and other sources of quantal phases [2,3] using trapped or guided atoms have been envisioned. In particular, the sensitivity of atom-interferometric gyroscopes is proportional to the area enclosed by the closed loop around which atoms are guided [4]. Such considerations motivate the development of closed-loop atom waveguides that enclose a sizable area.

A number of multiply connected trapping geometries for cold atoms have been discussed. Optical traps using high-order Gauss-Laguerre beams were proposed [5,6], and hollow light beams were used to trap nondegenerate atoms in an array of small-radius rings [7]. Large-scale magnetic storage rings were developed for cold neutrons [8] and discussed for atomic hydrogen [9]. More recently, closed-loop magnetic waveguides were demonstrated for laser cooled atoms [10,11]. Unfortunately, these guides are characterized by large variations in the potential energy along the waveguide and by high transmission losses at points where the magnetic field vanishes.

In this Letter, we report the creation of a smooth, stable circular waveguide for ultracold atoms. A simple arrangement of coaxial electromagnetic coils was used to produce a static ring-shaped magnetic trap, which we call the quadrupole ring (Q ring), in which strong transverse confinement is provided by a two-dimensional quadrupole field. Atoms trapped in the Q ring experience large Majorana losses, but we can eliminate such losses with a time-orbiting ring trap (TORT) [12]. In this manner, stable circular waveguides with diameters ranging from 1.2 to

3 mm were produced. Finally, we report on the production of Bose-Einstein condensates (BECs) in a portion of the circular waveguide, and on the guiding of an ultracold atomic beam for several revolutions around the guide. This ring-shaped trap presents opportunities for studies of BECs that are homogeneous in one dimension and therefore of the untruncated propagation of sound waves [13] and solitons [14–16], of persistent currents [17–20], of quantum gases in low dimensions, and of matter-wave interferometry.

To explain the origin of the quadrupole ring trap, we consider a cylindrically symmetric static magnetic field \vec{B}_c in a source-free region. Expanding \vec{B}_c to low order about a point (taken as the origin) on the axis where the field magnitude has a local quadratic minimum, we have

$$\vec{B}_c = B_0 \hat{z} + \frac{B_z''}{2} \left[\left(z^2 - \frac{x^2 + y^2}{2} \right) \hat{z} - z(x\hat{x} + y\hat{y}) \right], \quad (1)$$

where $B_0 > 0$ is the field magnitude at the origin, B_z'' is the

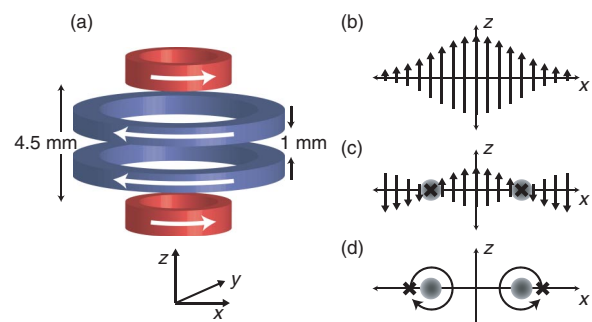


FIG. 1 (color). Forming a circular magnetic waveguide. (a) Four coaxial circular electromagnets (see [21] for details) are used to generate both the static (currents as shown) and rotating fields needed for the waveguide. Axes are indicated; gravity points along $-\hat{z}$. (b) As shown schematically, the field (arrows) from just the two outer coils (curvature coils, outer pair) points axially in the midplane between the coils, with largest fields at the axis. (c) Adding a uniform opposing bias field (using antibias coils, inner pair) produces a ring of field zeros (\times) in the \hat{x} - \hat{y} plane around which weak-field seeking atoms (shaded region) are trapped. (d) Rapidly rotating the field zeros around the trapped atoms produces the TORT.

axial field curvature, and Cartesian coordinates (x, y, z) are chosen so that z is the axial coordinate. The magnetic field magnitude falls to zero in the \hat{x} - \hat{y} plane along a circle of radius $\rho_0 = 2\sqrt{B_0/B_z''}$ centered at the origin. This is the Q ring, a ring-shaped magnetic trap for weak-field seeking atoms. Near the field zeros, the magnetic field has the form of a transverse (radial and axial directions) two-dimensional quadrupole field with gradient $B' = \sqrt{B_0 B_z''}$. Such traps can also be obtained using different electromagnet configurations [12].

In our apparatus, the Q ring is formed using a subset of the coils (the curvature and antibias coils; see Fig. 1) used in our recently demonstrated millimeter-scale Ioffe-Pritchard magnetic trap [21]. Our work is aided, in particular, by the large axial curvatures produced in this trap and by the vertical orientation of the trap axis. These features are relevant for the operation of a Q ring in the presence of gravity, for two reasons. First, trapping atoms in the Q ring requires transverse confinement sufficient to overcome the force of gravity; this places a lower bound on the radius of the Q ring of $\rho_0 > \rho_{\min} \approx 2mg/|\mu|B_z''$ with m the atomic mass, g the acceleration due to gravity, and μ the atomic magnetic moment. Indeed, if ρ_{\min} exceeds the range over which Eq. (1) is valid, typically the distance to the field-producing coils, the formation of a Q ring may be precluded entirely. With $B_z'' = 5300$ G/cm² in our experiments, $\rho_{\min} = 115$ μ m is much smaller than the millimeter dimensions of the electromagnets used for the trap. Second, the vertical orientation of the Q -ring axis allows cold atoms to move slowly along the nearly horizontal waveguide rather than being confined in a deep gravitational well.

Atoms can be localized to a particular portion of the Q ring by application of a uniform sideways (in the \hat{x} - \hat{y} plane) magnetic field; e.g., a weak bias field $B_s \hat{x}$ tilts the Q ring by $\Delta z/\rho_0 = (B_s/B')/\rho_0$ about the \hat{y} axis. This adjustment also adds an azimuthal field of magnitude $B_s |\sin\theta|$, splitting the Q ring into two trap minima at opposite sides of the ring, with θ being the azimuthal angle conventionally defined.

We loaded cold atoms into the Q ring using a procedure similar to previous work [21]. Briefly, about 2×10^9 ⁸⁷Rb atoms in the $|F = 1, m_F = -1\rangle$ hyperfine ground state were loaded into one of two adjacent spherical quadrupole magnetic traps. Using these traps, atoms were transported 3 inches from the loading region to the Q -ring trap region. During this transport, rf evaporative cooling was applied, yielding 2.5×10^7 atoms at a temperature of 60 μ K in a spherical quadrupole trap with an axial field gradient of 200 G/cm. Within 1 s, we then converted the spherical quadrupole to a tilted Q -ring trap produced with $B_z'' = 5300$ G/cm², $B_0 = 22$ G, and a side field of magnitude $B_s = 9.2$ G. This process left 2×10^7 atoms trapped in the Q ring (Fig. 2).

The trapping lifetime of atoms in the Q ring is limited by Majorana losses. In a balanced Q ring, trapped atoms

passing close to the line of zero field, which extends all around the ring, may flip their spins and be expelled from the trap. Extending the treatment by Petrich *et al.* [22] to this scenario, we estimate a Majorana loss rate of $\frac{\hbar^{1/2}}{\pi m^{3/4}} \times \frac{(\mu B')^{3/2}}{(k_B T)^{5/4}} = 6$ s⁻¹ for our trap at a temperature of 60 μ K. In a tilted Q ring, the zero-field region is reduced to just two points at opposite sides of the ring. Majorana losses in a tilted Q ring are thus similar to those in spherical quadrupole traps and much smaller than in a balanced Q ring. We confirmed this qualitative behavior by measuring the lifetime of trapped atoms in balanced and tilted Q -ring traps. In the balanced Q ring, the measured 0.3 s⁻¹ Majorana loss rate was thrice that in a tilted Q ring, while falling far short of the predicted 6 s⁻¹ loss rate, presumably due to residual azimuthal fields.

The high loss rates in the Q ring can be mended in a manner similar to the time-orbiting potential (TOP) trap by which Majorana losses in a spherical quadrupole field were overcome [22]. As proposed by Arnold [12], a TORT with nonzero bias field can be formed by displacing the ring of field zeros away from and then rapidly rotating it around the trapped atoms [Fig. 1(d)]. From Eq. (1), the Q ring can be displaced radially by application of an axial bias field, and displaced along \hat{z} by a cylindrically symmetric spherical quadrupole field. The TORT provides transverse quadratic confinement with an effective field curvature of $B_{\text{eff}}'' = B'^2/2B_{\text{rot}}$, where B_{rot} is the magnitude of the rotating field seen at the trap minimum. Just as the TOP trap depth is limited by the ‘‘circle of death,’’ the TORT trap

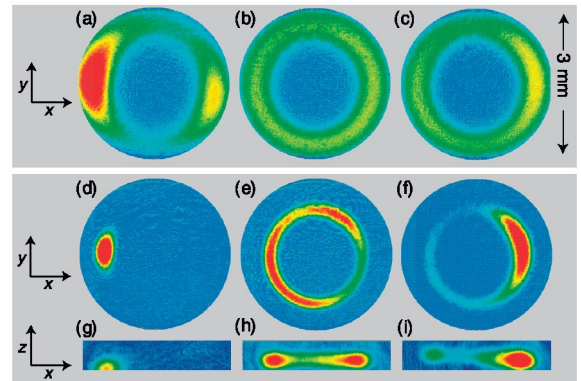


FIG. 2 (color). Atoms in a ring-shaped magnetic trap. Shown are top (a)–(f) and side (g)–(i) absorption images of ultracold ⁸⁷Rb clouds in either a Q ring (a)–(c) or TORT (d)–(i) with applied side field $B_s = 9.2$ (left), 0 (middle), and -2.5 G (right column), respectively, in the \hat{x} direction. Images were taken 2 ms after turning off the traps. The applied field tilts the Q ring or TORT and favors atomic population in one side or another of the trap. For $B_s \sim 0$, the trap lies nearly in the horizontal plane and its azimuthal potential variation is minimized. For the Q ring, $B_0 = 22$ G; for the TORT, $B_0 = 20$ G and $B_{\text{rot}} = 17$ G; and $B_z'' = 5300$ G/cm² for both. The temperature of trapped atoms is 90 μ K in the Q ring, and 10 μ K in the TORT. Resonant absorption ranges from 0 (blue) to $>80\%$ (red).

depth is limited by a “torus of death,” the locus of points at which the magnetic field is zero at some time [23]. This scheme may be applied equally to a tilted Q ring, yielding a tilted TORT and providing a stable trap in which atoms are confined to a portion of the ring. The sideways magnetic field (e.g., along \hat{x}) causes the magnetic potential minimum to vary azimuthally in the tilted TORT as $|\mu|\sqrt{B_{\text{rot}}^2 + B_s^2 \sin^2 \theta}$, while the gravitational potential is the same as that in a Q ring.

The time-varying fields needed to convert our Q ring (or tilted Q ring) traps to TORT (or tilted TORT) traps were obtained by suitably modulating the currents in the four coils used to generate the Q -ring potential. A modulation frequency of 5 kHz was chosen to be much larger than the transverse motional frequencies (< 100 Hz) and also much smaller than the Larmor frequency (> 3 MHz) at the location of the trap minimum. To first switch on the TORT, a rotating field magnitude of $B_{\text{rot}} = 18$ G was used [24].

As shown in Fig. 3, the trap lifetime was dramatically increased by application of the TORT trap. In the first few seconds after switching on the TORT, we observed a fast loss of atoms and a simultaneous drop in their temperature. We ascribe this loss and cooling to the evaporation of atoms from the trapped cloud through the “torus of death.” As the temperature dropped, the evaporation rate diminished and the lifetime of trapped atoms became vacuum limited at 90 s, a value observed both for balanced and for tilted TORT traps.

Given their longevity, it was possible to cool evaporatively the TORT-trapped atoms to the point of quantum degeneracy. Using a tilted TORT with $B_s \sim 9$ G, evaporation was performed in two stages. First, “torus of death” evaporation was applied by ramping down the rotating field strength B_{rot} over 40 s to 4.8 G. The oscillation frequencies in this trap were measured as $\omega_{\perp} = 2\pi \times (87, 74.5)$ Hz in the transverse and $\omega_{\theta} = 2\pi \times 35$ Hz in the azimuthal directions, in agreement with predictions [25]. In the sec-

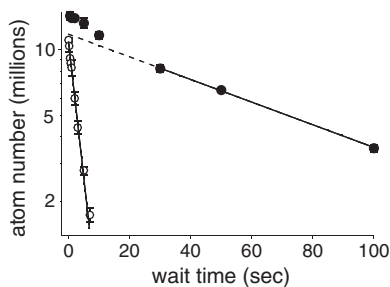


FIG. 3. Elimination of Majorana losses in the TORT. The measured number of trapped atoms in a Q ring (open circles) or TORT (solid circles) trap is shown vs residence time in the trap. Exponential fits indicate a 3 s Majorana-loss-limited lifetime in the Q ring. In the TORT, following an initial (30 s) loss of atoms due to evaporation, a vacuum limited lifetime of 90 s was observed. Settings for B_0 , B_z'' , and B_{rot} are as in Fig. 2, $B_s \sim 0$, and the initial temperature is $60 \mu\text{K}$.

ond stage, rf evaporation was applied for 20 s, yielding clouds of up to 6×10^5 atoms at the Bose-Einstein condensation temperature, and pure BECs of up to 3×10^5 atoms.

Finally, to assess the suitability of the TORT as an atomic waveguide for interferometry, we launched our trapped BECs into closed-loop circular motion along the guide. This was accomplished by reorienting the sideways bias field B_s , inducing the trapped BEC to accelerate toward the newly positioned tilted TORT trap minimum (advanced by an azimuthal angle of about $\pi/4$), while simultaneously reducing the magnitude of B_s to $B_s \sim 0$ and increasing B_{rot} to 12.6 G to produce a well-balanced TORT trap. The TORT was then maintained at this setting, with radius $\rho_0 = 1.25$ mm ($B_0 = 20$ G, $B_z'' = 5300$ G/cm²), and transverse trap frequencies of $\omega_{\perp} \approx 2\pi \times 50$ Hz as measured at the launch point of the atoms. The atoms were allowed to propagate freely around the guide for various guiding times before being observed by absorption imaging. As shown in Fig. 4, the ultracold atomic beam propagated around the circular waveguide at an angular (linear) velocity of 40.5 rad/s (50.6 mm/s). As measured from the azimuthal extent of the atoms for different guiding times, this pulsed atom beam was characterized by an azimuthal rms velocity spread of 1.4 mm/s, equivalent to a longitudinal kinetic temperature

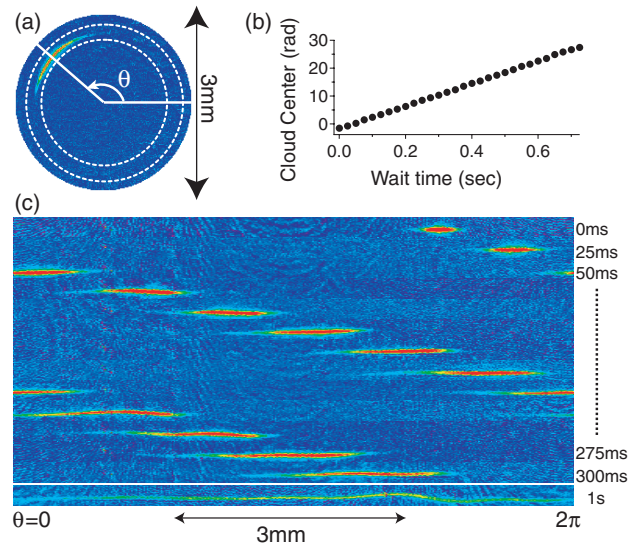


FIG. 4 (color). Circular motion of a quantum degenerate atomic beam in a waveguide. A Bose-Einstein condensate was launched into a balanced TORT and allowed to propagate. (a) Top view in-trap absorption image during the propagation. The mean azimuthal position of the BEC measured from such images is shown in (b). Annular portions (indicated by dashed circles) of top-view images taken at different guiding times are shown in (c) displayed in polar coordinates (radius vs azimuthal angle). The beam advances at an angular velocity of 40.5 rad/s while expanding due to an rms azimuthal velocity spread of 1.4 mm/s. After 1 s, the beam fills the entire guide.

of 22 nK. After about 1 s of guiding, this velocity variation caused the atomic cloud to spread throughout the waveguide, by which point the atoms had traveled $L = 51$ mm along the waveguide, encompassing an area of $A = L\rho_0/2 = 32$ mm². One would expect the front and back ends of the expanding BEC to interfere spatially as they overlap after 1 s of propagation in the guide. However, the fringe periodicity of ~ 600 nm expected for such interference is well below our imaging resolution of several μm , and we observed no such interference.

Many requisite elements for interferometric rotation sensing are still lacking in our system, including a means of in-guide coherent atomic beam splitting [26,27], bi-directional propagation, proper radial waveguiding [28], full characterization of longitudinal coherence in the beam, an assessment of the influence of the time-orbiting field on sensor precision, and atom-interferometric stability. Nevertheless, it is valuable to consider the possible sensitivity of our system if these elements are attained. As limited by atomic shot noise, rotation measurements with an uncertainty of $\Delta\Omega = (\hbar/4mA)N_0^{-1/2} \sim 1 \times 10^{-8}$ rad/s could be made from a single (1 s long) measurement, where $N_0 = 3 \times 10^5$ is the number of atoms used. While this figure is nearly 20 times that of existing atom-based gyroscopes [4], improvements such as launching the atoms at higher velocities, increasing the TORT radius, and increasing the atom number may ultimately yield a useful, compact sensing device.

Other applications of the TORT may include studies of propagation [29–31] and nonlinear dynamics [32] in atomic waveguides. In a TORT potential that is modified either by application of magnetic fields or by tilting with respect to gravity, BECs can be studied both undergoing pendular motion (terminated guide) when launched at small velocities, and undergoing circular motion (unterminated guide) at larger velocities.

Another appealing possibility is the study of BECs that fill the ring-shaped trap, rather than forming in just a portion of the ring. Such a system would allow for studies of quantized and persistent circulation [17–20], untruncated motion of solitons [14–16], and other aspects of nonlinear dynamics [33]. For this purpose, the azimuthal variation in the TORT potential must be reduced below the typical ~ 100 nK scale of the BEC chemical potential. From measurements of the kinetic energy of BEC's undergoing circular motion in our trap (data of Fig. 4), we estimate the TORT potential varied by as much as 5 μK . We believe this figure can be reduced greatly by using traps of smaller radius, by designing better electromagnets for both the static and the rotating fields used for the TORT, and also by controlling the orientation of the electromagnets with respect to gravity.

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Note added.—After the submission of this paper, related work was reported by Arnold and collaborators [34].

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